# Reconstructing quantum states with generative models<sup>[1]</sup>

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## Structure

- Quantum computing / Quantum state tomography
- IC-POVMs
- General idea / Methods
- Results
- Summary

## A general introduction to quantum computing

- Ongoing field of research since the 80th
- Qubit = two state system
  - Use  $|0\rangle \, {\rm and} \, |1\rangle$  instead of classical bits 0 and 1
- Research by: Google, Microsoft, Intel, IBM and Alibaba
- Application:
  - Cryptography
  - Material science
  - Stock market
  - ...





Select investors betting on quantum computing startups [2] 2010 - 2019 YTD (1/7/2019)



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## The Bloch sphere

• A general single qubit state can be visualized on the Bloch sphere by the Bloch vector

$$\vec{v} = \left(\begin{array}{c} \langle \sigma_x \rangle \\ \langle \sigma_y \rangle \\ \langle \sigma_z \rangle \end{array}\right)$$

For a pure state the angles θ and φ can be obtained by the parametrisation (up to a gobal phase):

 $|\psi\rangle = \cos(\frac{\theta}{2})|0\rangle + e^{i\phi}\sin(\frac{\theta}{2})|1\rangle$ 



The representation of a state  $\psi$  on the Bloch sphere [3]

#### Quantum state tomography

how good does a quantum computer work?

- Quantum state tomography is the method to benchmark the system
- Many measurements on a well-known state are used to reconstruct the state
- The reconstructed state is compared to the original one to quantify the quality of the quantum computer



#### Scheme of quantum state tomography

## EXAMPLE: Single Qubit state tomography

• An arbitrary single qubit density matrix can be written as:

$$\hat{\rho} = \frac{1}{2} \sum_{i=0}^{3} S_i \hat{\sigma}_i \quad \hat{\sigma}_0 = \hat{\mathbb{1}} \quad S_i = \text{Tr}(\hat{\sigma}_i \hat{\rho})$$

- Only the expectation values are required
- IBM provides an online platform that provides access to a few qubit quantum computer, so we can try it out! (use 100 measurements for every S<sub>i</sub>)





Philipp Schultzen



Quality criterium:  $F(\rho_{state}, \rho_{model}) = \text{Tr}[\sqrt{\sqrt{\rho_{state}}\rho_{model}\sqrt{\rho_{state}}}] = 0.982$ 

**Problem:** In general there are 4<sup>N</sup> -1 free parameter which have to be obtained from measurements:

• Makes exact state tomography impractical for large N: *curse of dimensionality* 

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Information Complete Positive Operator Valued Measure (IC-POVM)

#### **Definition POVM:**

- Set of Measurements  $\mathbf{M} = \{\mathbf{M}^{(\alpha)}\}$
- Each  $M^{(\alpha)}$  is a positive semidefinite hermitian operator
- $\sum_{\alpha} M^{(\alpha)} = \mathbb{1}$

$$1 = Tr(\rho) = \sum_{i} \langle i | \rho | i \rangle$$
  
=  $\sum_{i} \langle i | \sum_{\alpha} \rho M^{(\alpha)} | i \rangle$   
=  $\sum_{\alpha} Tr(\rho M^{(\alpha)})$   
=  $\sum_{\alpha} P_{\alpha}$  with  $P_{\alpha} \coloneqq Tr(\rho M^{(\alpha)}) \ge 0$ 

#### **Definition IC-POVM:**

• the POVM spans the whole space of bounded norm operators on the observed hilbert space

A probability can be assigned to each measurement  $M^{(\alpha)}$  !

## Example: Thetrahedral single qubit IC-POVM

• 4 rank-1 projectors are used (see picture)

$$\begin{split} \boldsymbol{M}_{\text{tetra}} &= \left\{ M^{(a)} = \frac{1}{4} (\mathbb{1} + \boldsymbol{s}^{(a)} \cdot \boldsymbol{\sigma}) \right\}_{a \in \{0, 1, 2, 3\}} \\ s^{0} &= (0, 0, 1) \\ s^{1} &= \left(\frac{2\sqrt{2}}{3}, 0, -\frac{1}{3}\right) \\ s^{2} &= \left(-\frac{\sqrt{2}}{3}, \sqrt{\frac{2}{3}}, -\frac{1}{3}\right) \\ s^{3} &= \left(-\frac{\sqrt{2}}{3}, -\sqrt{\frac{2}{3}}, -\frac{1}{3}\right) \\ P_{\alpha} &= \operatorname{Tr}(\rho M^{(\alpha)}) \end{split}$$



- Couple to an ancillary qubit •
- Apply a unitary operator Uin a way, that the new state shows  $P_{\alpha}$



 $10\rangle$ 

 $|11\rangle$ 

 $|1\rangle$ 

 $|0\rangle$ 

Single-qubit  $M_{thetra} \implies Multi-qubit M_{Thetra}$ 

Natural extension:

- Apply M<sub>thetra</sub> on each single qubit
- $M \to M_1 \otimes M_2 \otimes M_3 \dots$

Requires only single-qubit measurements:

• Easy implementation in real experiment



#### IC-POVMs: some mathtools

#### How can we model ρ with an IC-POVM?

$$\rho \stackrel{IC}{=} \sum_{\alpha'} c_{\alpha'} M^{(\alpha')} = \vec{c} \cdot \vec{M}$$
$$\vec{P} = \{P_{\alpha}\} = \operatorname{Tr}(\rho \vec{M}) = \vec{c} \cdot \mathbf{T}$$
with  $\mathbf{T}_{\alpha,\beta} \coloneqq \operatorname{Tr}(M^{(\alpha)}M^{(\beta)})$ 
$$\vec{c} = \vec{P} \cdot \mathbf{T}^{-1}$$
$$\rho = \vec{P} \cdot \mathbf{T}^{-1} \vec{M} = \mathbb{E}_{P}[\mathbf{T}^{-1} \vec{M}]$$

Multi-qubit M<sub>thetra</sub>:

$$\rho = \mathbb{E}_P \left[ \mathbf{T}^{-1} \vec{M_1} \otimes \mathbf{T}^{-1} \vec{M_2} \otimes \dots \right]$$

How can we estimate arbitrary operators O?

$$O \stackrel{IC}{=} \sum_{\alpha'} d_{\alpha'} M^{(\alpha')} = \vec{d} \cdot \vec{M}$$
$$\vec{d} = \operatorname{Tr}(O\vec{M}) \cdot \mathbf{T^{-1}}$$

- determines each d univocally

$$\langle O \rangle = \operatorname{Tr}(\rho O) = \vec{d} \cdot \vec{P} = \mathbb{E}_P[\vec{d}]$$

$$c, d \in \mathbb{R};$$
  $M^{(\alpha')}, O = \text{operators}$   
 $\vec{M} = \text{vector of operators}$   
 $\vec{c}, \vec{d} = \text{vector of coefficients}$ 



## Restricted Boltzmann machine (RBM)

Energy:  $E(v,h) = -\sum_{i,j,k,l} W_{ij}^{kl} v_i^k h_j^l - \sum_{i,k} b_i^k v_i^k - \sum_{j,l} a_j^l h_j^l$ 

• *W* = Weights

upper indices = dimension of v, h

- *v*, *h* = visible and hidden units
- *b, a* = *biases for v and h respectively*

#### Training:

Gibbs sampling (1 step): v is the measurement P(h|v) is calculated, h is drawn from this distribution P(v|h) is calculated, v<sub>new</sub> is drawn from this distribution

Update general parameter  $\theta \in (W, b, a)$  by contrastive divergence

 $\mathcal{L} = \langle \log(p_{\theta}(x)) \rangle_{data}$ 

$$\frac{\partial \mathcal{L}}{\partial \theta_i} = \mathbb{E} \Big[ \frac{\partial E(\mathbf{v}, \theta)}{\partial \theta_i} \Big] - \mathbb{E} \Big[ \frac{\partial E(\mathbf{v}_{\text{new}}, \theta)}{\partial \theta_i} \Big]$$

#### Sampling:

Gibbs sampling after training (with optimized parameter)



Visible units v<sub>i</sub> of dimension M, where M is the number of measurements in the POVM

#### Recurrent neural network (RNN)



Variables:

- Inputs *x*
- Input weights U
- Outputs o
- Hidden state *s* (*memory*)
- Hidden weights W

RNN:

- Optimized to process sequential data with *memory*
- Each decision depends on the decision before:  $s_t = f(Ws_{t-1} + Ux_t)$ 
  - e.g. f is a ReLU or tanh

#### For this paper:

- Sequence step = one qubit
- Output = single qubit prob. dist.
- Hidden state = contains all relevant information for the output
- 1. P<sub>1</sub> is calculated (first qubit)
- P<sub>2</sub> is calculated with the memory of P<sub>1</sub>: P<sub>2</sub>|P<sub>1</sub> is obtained
  ...

This is done for every qubit, then we get

$$P_{model} = P_1 x (P_2 | P_1) x ... x (P_N | P_1, P_2 ... P_{N-1})$$

How do we measure the quality of the reconstruction?

• Kullback-Leibler divergence

 $D_{KL}(P_{state}||P_{model}) = \mathbb{E}_{P_{state}} \left[ \log \frac{P_{state}}{P_{model}} \right]$ 

- Measures how much P<sub>model</sub> diverges from P<sub>state</sub>
- Classical Fidelity

$$F_C(P_{state}, P_{model}) = \mathbb{E}_{\alpha \sim P_{state}} \left[ \sqrt{\frac{P_{model}(\alpha)}{P_{state}(\alpha)}} \right]$$

• Quantum Fidelity

$$F(\rho_{state}, \rho_{model}) = \text{Tr}[\sqrt{\sqrt{\rho_{state}}\rho_{model}\sqrt{\rho_{state}}}]$$

In the case of pure states this simplifies to

$$F(\psi_{state}, \psi_{model}) = \langle \psi_{state} | \psi_{model} \rangle$$





KL divergence for different gaussians. [6]

#### Which reconstruction criterium should be used?

- Measure the different reconstruction quality criteria using an RBM with the N=2 GHZ state
- As POVM, the multi-qubit thetrahedral POVM is used
- Add single-qubit depolarisation probability *p*

• Greenberger-Horne-Zeilinger (GHZ) state:  $|\Psi\rangle_{GHZ} = \frac{1}{\sqrt{2}}(|0\rangle^N + |1\rangle^N)$ 

• Depolarisation of a single qubit:  $\rho \rightarrow \mathbf{p} \cdot \left(\frac{1}{2} \cdot \mathbb{1}\right) + (1-p) \cdot \rho$ 

## Which reconstruction criterium should be used?



Epoch: 1 step of improving the weights and biases after Gibbs sampling

F<sub>c</sub> and F approach unity in a similar behaviour

## F<sub>c</sub> for different number of training samples N<sub>s</sub>

- F<sub>c</sub> is measured for different number of qubits N and for the depolarization probability p = 0.0 or 0.4
- Use the General GHZ state  $|\Psi\rangle_{GHZ} = \frac{1}{\sqrt{2}}(|0\rangle^N + |1\rangle^N)$
- Advantage of F<sub>C</sub>: It only requires P<sub>model</sub> and can be efficiently estimated for large N by simply sampling and averaging from P<sub>model</sub>
- RBM is difficult to train; RNN is used

$$F_C(P_{state}, P_{model}) = \mathbb{E}_{\alpha \sim P_{state}} \left[ \sqrt{\frac{P_{model}(\alpha)}{P_{state}(\alpha)}} \right]$$



How does the number of samples N<sub>s</sub> scale with the number of qubits for a given  $F_c$ ?

b • A Fidelity value is fixed [1]  $10 imes 10^4$  $N_s^* \coloneqq F_C(N_s^*) = 0.99$  $8 \times 10^{\circ}$ p=0.0 • The scaling is approximately linear!  $6 imes 10^4$  $N_s^*$ • Pure luck? Is it just a pecularity of M<sub>thetra</sub>?  $4 imes 10^4$  $2 \times 10^4$ 30 5060 102040 7080 N

## Pauli M<sub>6</sub> POVM

- Set of 6 measurements
- Projections on the 6 eigenstates of the three pauli matrices





 Even though 6 different measurements per single-qubit are used, a linear behaviour can still be observed!

#### Antiferromagnetic 1D transverse field ising model

$$\begin{aligned} \mathcal{H} &= J \sum_{\langle ij \rangle} \sigma_i^z \sigma_j^z + h \sum_i \sigma_i^x \\ J \gg h: \text{ The ground state is } |\uparrow\downarrow\uparrow\downarrow\uparrow\downarrow\dots\rangle \\ J \ll h: \text{ The ground state is } |\to\to\to\to\to\to\dots\rangle \end{aligned}$$

- The phase transition at J = h is observed
- The problem can be solved analytically and is therefore the testbed par excellence for quantum simulators.

- $F_c = 0.998$
- Also investigate 1-body correlation function



Antiferromagnetic Heisenberg model triangular lattice

 $\mathcal{H} = \sum_{\langle ij \rangle} \vec{\sigma}_i \vec{\sigma}_j$ 

- No analytical solution
- A 8 x 8 lattice is used (64 Qubits)

•  $F_c = 0.98$ 

• Remarkable agreement in the correlation function





Model for the path used by the RNN ( 4 x 4 lattice)

$$\langle O \rangle = \operatorname{Tr}(\rho O) = \vec{d} \cdot \vec{P} = \mathbb{E}_P[\vec{d}]$$

• Can be easily and efficiently computed!

## Conclusion

- Method: Quantum state tomography based on probability distribution estimation using unsupervised machine learning
- Could be certified on pure and mixed states, as well as on ground states of local hamiltonians
- A linear scaling between N and  $N_s^*$  is observed!
- Can efficiently estimate  $F_c$  and expectation values of local operators, e. g.  $\langle \sigma_x \rangle$

## References

- [1]: arXiv:1810.10584: Reconstructing quantum states with generative models
- [2]: https://www.cbinsights.com/research/report/quantumcomputing/
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#### Limitation:

 Expectation values of non-local operators can not always be well estimated by sampling, e.g. F