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Ising Model

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angle} \sigma_i \sigma_j$$
 with $\sigma_k = \pm 1$
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We are interested in the 2 dimensional case on a square lattice with periodic boundary conditions and J > 0:



J > 0 ferromagnetic $\leftrightarrow J < 0$ antiferromagnetic

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$$\Leftrightarrow T > T_{C} = \frac{2}{\ln 3}J \approx 1.8J$$

Qualitatively different behaviour at different temperatures



Phase Transition (Ehrenfest (1933)):

n-th order phase transition has a discontinuity at any *n*-th partial derivative of the free energy F.

Exact solution in the thermodynamic limit (Onsager (1944)):

$$-\beta f = \ln 2 + \frac{1}{8\pi} \int_0^{2\pi} \mathrm{d}\,\phi\,\mathrm{d}\,\theta\,\ln\left[\cosh^2 2\beta J - \sinh 2\beta J\left(\cos\phi + \cos\theta\right)\right]$$

Heat capacity $c=-\beta^2\partial_\beta^2 f$ has a log-divergence at the critical temperature

$$T_C = \frac{2}{\ln\left(1 + \sqrt{2}\right)} J \approx 2.27 J$$

Easy to determine from magnetization:

$$m = \begin{cases} \left[1 - (\sinh 2\beta J)^4\right]^{\frac{1}{6}}, \ T < T_C\\ 0, \ T > T_C \end{cases}$$

Phase Transition in the Ising Model



Source: http://farside.ph.utexas.edu/teaching/329/lectures/node110.html

Published on 13 February 2017 by Carrasquilla and Melko

Try to find critical temperature of a ferromagnetic (square) Ising model with supervised machine learning

Analytical toy model with free parameter ϵ :



Input layer is $\sigma_i = \pm 1$

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Hidden layer consists of perceptrons (Heaviside)

Analytical toy model with free parameter ϵ :



$$Wx + b = \frac{1}{(1+\epsilon)} \begin{pmatrix} m(x) - \epsilon \\ -m(x) - \epsilon \\ m(x) + \epsilon \end{pmatrix}$$

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Learning toy model:

Train same network with arbitrary weights far from T_C





Fully connected network with single hidden layer consisting of 100 sigmoid neurons



 $\rightarrow T_C = (2.266 \pm 0.002) J$

Literature: $T_C = 2.2692J$



Use same network without retraining on a triangular Ising model



 $\rightarrow T_C = (3.65 \pm 0.01)J$ Literature: $T_C = 3.641J$

Fermion Sign Problem

Heisenberg model for spin 1/2-particles:

$$\mathcal{H}=-J\sum_{\langle i,j
angle}ec{S}_iec{S}_j$$

Two particle Hamiltonian $H = -J\vec{S_1}\vec{S_2}$ has eigenbasis

$$\left|\uparrow\uparrow\right\rangle, \left|\downarrow\downarrow\right\rangle, \left|\frac{1}{\sqrt{2}}\left(\left|\uparrow\downarrow\right\rangle - \left|\downarrow\uparrow\right\rangle\right), \left|\frac{1}{\sqrt{2}}\left(\left|\uparrow\downarrow\right\rangle + \left|\downarrow\uparrow\right\rangle\right)$$

Entangled eigenstates!

1d-Chain with periodic boundary conditions:

$$\mathcal{H} = -J\sum_{i} \vec{S}_{i}\vec{S}_{i+1} = \sum_{i} H_{i} = \sum_{\text{even } i} H_{i} + \sum_{\text{odd } i} H_{i} = \mathcal{H}_{e} + \mathcal{H}_{o}$$

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Try to get the partition function numerically:

$$Z = \operatorname{Tr} e^{-\beta \mathcal{H}} = \operatorname{Tr} \left(e^{-\Delta \tau \mathcal{H}} \right)^{m} = \operatorname{Tr} \left[\left(e^{-\Delta \tau \mathcal{H}_{e}} e^{-\Delta \tau \mathcal{H}_{o}} \right)^{m} \right] + \mathcal{O}(\Delta \tau^{2})$$

$$= \sum_{\substack{|\sigma_{1}\rangle, \dots, |\sigma_{2m}\rangle \\ \{|\sigma_{i}\rangle\} \text{ Basis } \forall i}} \langle \sigma_{1} | e^{-\Delta \tau \mathcal{H}_{e}} | \sigma_{2m} \rangle \langle \sigma_{2m} | e^{-\Delta \tau \mathcal{H}_{o}} | \sigma_{2m-1} \rangle \dots \langle \sigma_{2} | e^{-\Delta \tau \mathcal{H}_{o}} | \sigma_{1} \rangle$$

$$+ \mathcal{O}(\Delta \tau^{2}) = \sum_{\omega} p(\omega)$$

ightarrow Is $p(\omega) \ge 0$?

$$Z = \sum_{\substack{|\sigma_1\rangle, \dots, |\sigma_{2m}\rangle \\ \{|\sigma_i\rangle\} \text{ Basis } \forall i}} \langle \sigma_1 | e^{-\Delta \tau \mathcal{H}_e} | \sigma_{2m} \rangle \langle \sigma_{2m} | e^{-\Delta \tau \mathcal{H}_o} | \sigma_{2m-1} \rangle \dots \langle \sigma_2 | e^{-\Delta \tau \mathcal{H}_o} | \sigma_1 \rangle$$

Look at "spin flip" term:

$$\left<\uparrow\downarrow\right|e^{-\Delta auec{S_1}ec{S_2}}\left|\downarrow\uparrow\right>=e^{-\Delta au J/4}\sinhrac{\Delta au J}{4}<0 ext{ if } J<0$$

$$Z = \sum_{\substack{|\sigma_1\rangle, \dots, |\sigma_{2m}\rangle \\ \{|\sigma_i\rangle\} \text{ Basis } \forall i}} \left\langle \sigma_1 \right| e^{-\Delta \tau \mathcal{H}_e} \left| \sigma_{2m} \right\rangle \left\langle \sigma_{2m} \right| e^{-\Delta \tau \mathcal{H}_o} \left| \sigma_{2m-1} \right\rangle \dots \left\langle \sigma_2 \right| e^{-\Delta \tau \mathcal{H}_o} \left| \sigma_1 \right\rangle$$

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angle = e^{-\Delta au J/4} \sinh rac{\Delta au J}{4} < 0 ext{ if } J < 0$$

Two particle system (m=1 for simplicity):

$$Z = \dots + \langle \uparrow \downarrow | e^{-\Delta \tau \vec{S}_1 \vec{S}_2} | \downarrow \uparrow \rangle \langle \downarrow \uparrow | e^{-\Delta \tau \vec{S}_2 \vec{S}_1} | \uparrow \downarrow \rangle + \dots$$

 $\rightarrow p \geq 0$

Three particles: $\mathcal{H} = -J\left(ec{S_1}ec{S_2} + ec{S_2}ec{S_3} + ec{S_3}ec{S_1}
ight)$:

$$Z = \sum \left\langle \sigma_1 \right| e^{-\Delta \tau \vec{S_1} \vec{S_2}} \left| \sigma_3 \right\rangle \left\langle \sigma_3 \right| e^{-\Delta \tau \vec{S_2} \vec{S_3}} \left| \sigma_2 \right\rangle \left\langle \sigma_2 \right| e^{-\Delta \tau \vec{S_3} \vec{S_1}} \left| \sigma_1 \right\rangle$$

One can create a configuration similar to before which generates negative probabilites:

 $\left|\downarrow\uparrow\uparrow\right\rangle,\left|\uparrow\downarrow\uparrow\right\rangle,\left|\uparrow\uparrow\downarrow\right\rangle$

 \rightarrow "frustrated antiferromagnet"

Ingredients for trouble:

- fermions
- negative couplings
- non bipartite lattice



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theoretical fix: always work in hamiltonian eigenbasis!

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but diagonalization scales exponentially \rightarrow NP-complete problems

Machine learning quantum phases of matter beyond the fermion sign problem (2017)

Also treat sign-problematic models with machine learning (Hubbard model)

(a)



(b)



Learning phase transitions by confusion (2017)

Published on 13 February 2017 by van Nieuwenburg, Liu and Huber Try to find an unknown critical point of a system with pseudo-supervised machine learning Suppose a model which depends on a parameter T and there exists some critical value T_C in the interval (T_a, T_b) which yields differently structured output above and below T_C .

General Strategy to find the critical point T_C :

- Propose some critical point T'_C
- Generate configurations far from T'_C and train the network with it
- Generate more configurations and record the performance of the trained network

Do the above for all $T'_C \in (T_a, T_b)$.

The performance $P(T'_C)$ should be maximal for $T'_C = T_a$, T_C , T_b as it is easier to distinguish between differently structured configurations which should be separated by T_C :

$$P(T_{C}') \propto 1 - \frac{\min\{|T_{C}' - T_{a}|, |T_{C}' - T_{C}|, |T_{C}' - T_{b}|\}}{T_{b} - T_{a}}$$

Learning phase transitions by confusion (2017)



Ising model with T_C being the on at the thermodynamic limit

Learning phase transitions by confusion (2017)



Heisenberg model with external field; $T_C = 3J$ being the literature value

Neural networks can be trained to distinguish different phases



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Neural networks can find unkown phase transitions on their own



Neural networks can be trained to distinguish different phases but choosing the correct approach and network may be difficult.

Neural networks can find unkown phase transitions on their own but that doesn't have to be a phase transition.



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lsing model and statistical physics
Schwarz: Statistical Physics
https://www.thphys.uni-heidelberg.de/~biophys/PDF/Skripte/
StatPhys.pdf
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Quantum monte carlo Assaad and Evertz: World line and determinantal Quantum Monte Carlo methods for spins, phonons, and electrons. https://pawn.physik.uni-wuerzburg.de/~assaad/Reprints/ assaad_evertz.pdf

Paper

Machine learning phases of matter DOI:10.1038/nphys4035 Machine learning quantum phases of matter beyond the fermion sign problem DOI:10.1038/s41598-017-09098-0 Learning phase transitions by confusion DOI: 10.1038/nphys4037