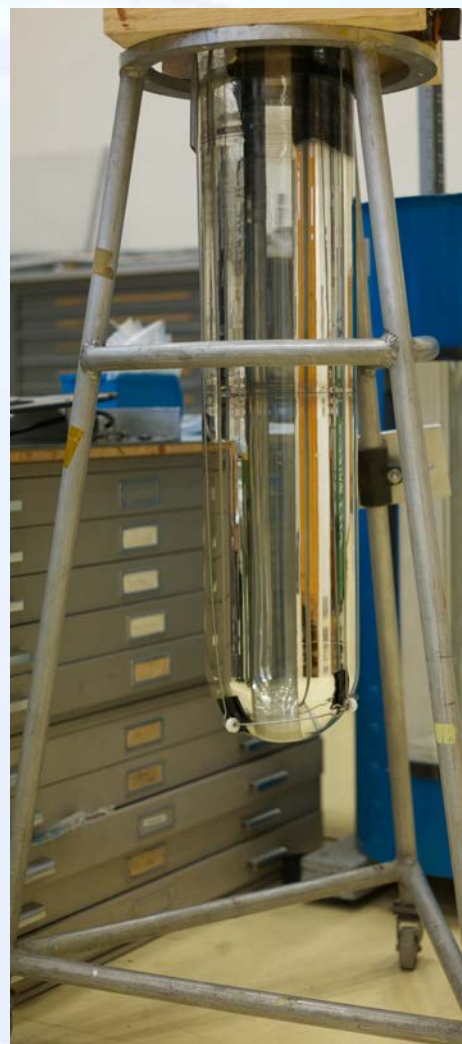
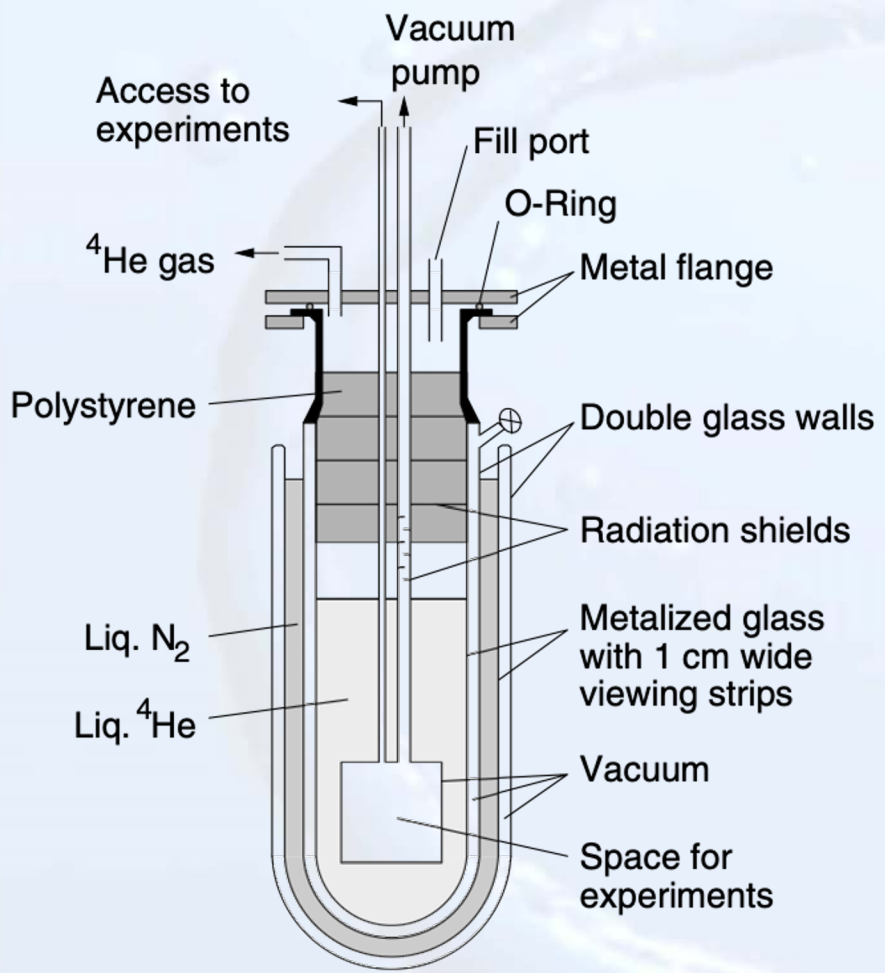




11. Cooling Techniques

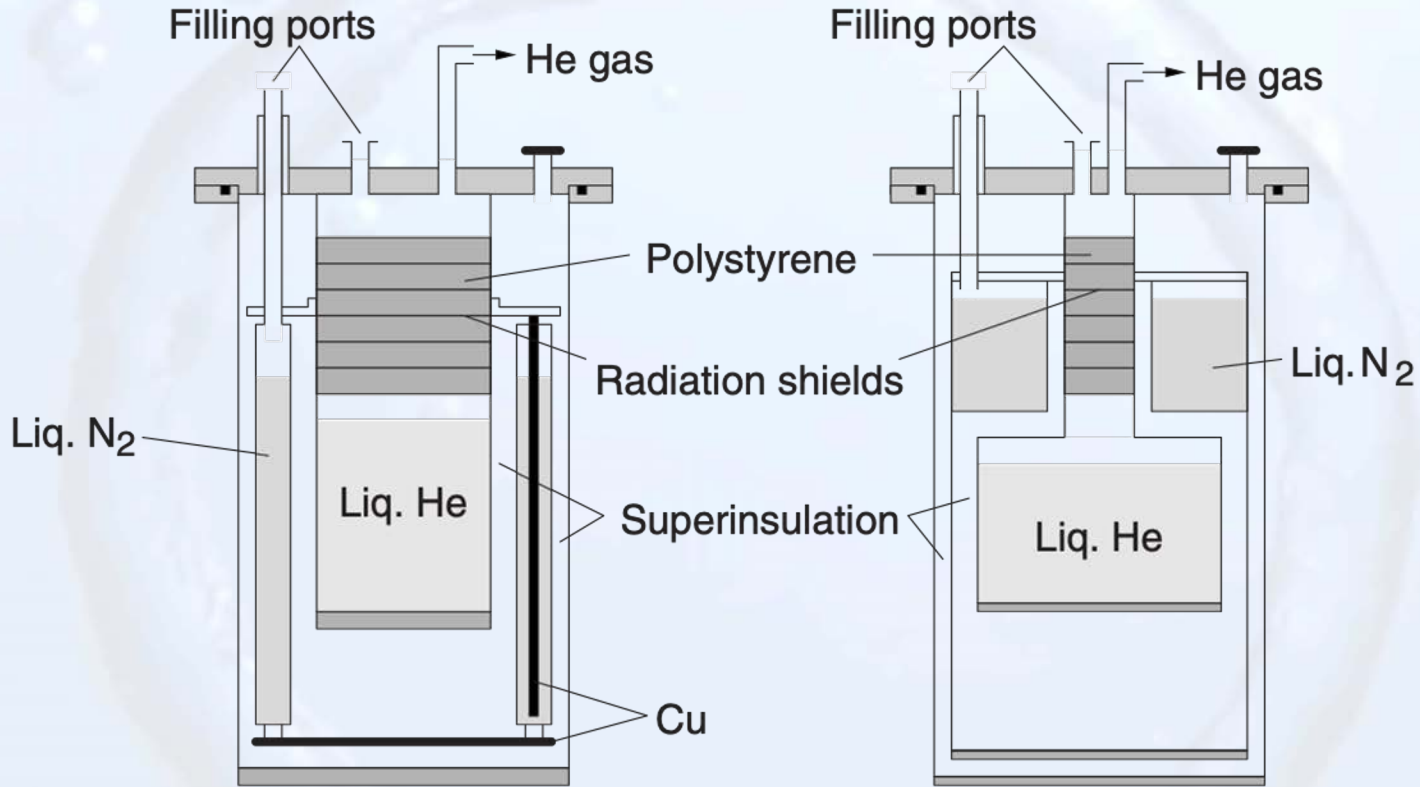


^4He bath cryostat: glass dewar





^4He Bath cryostat: metal dewar



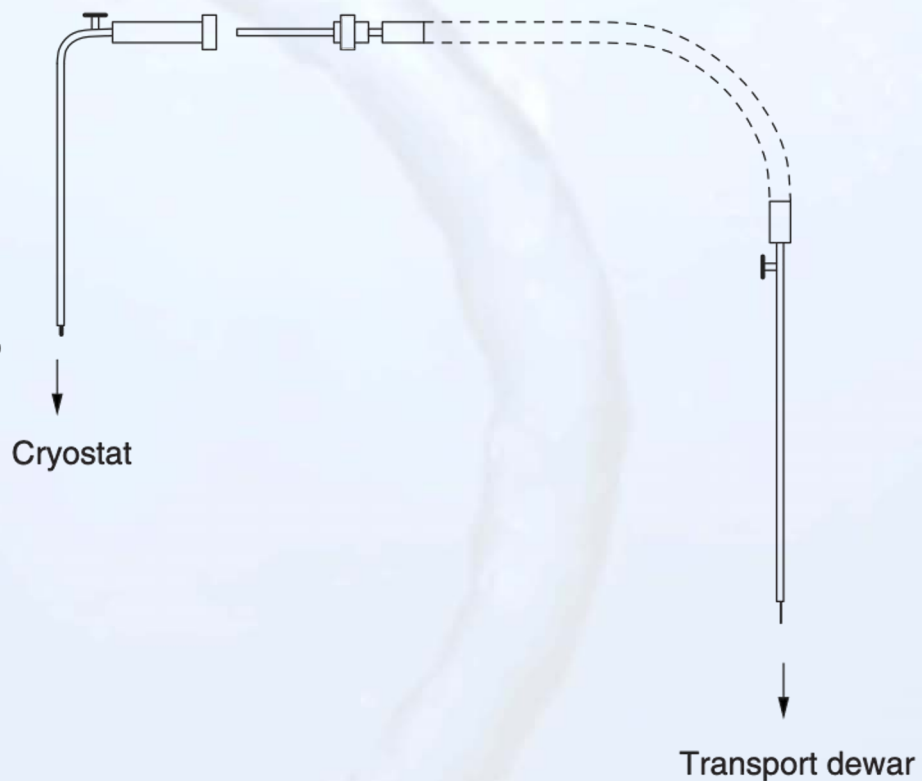
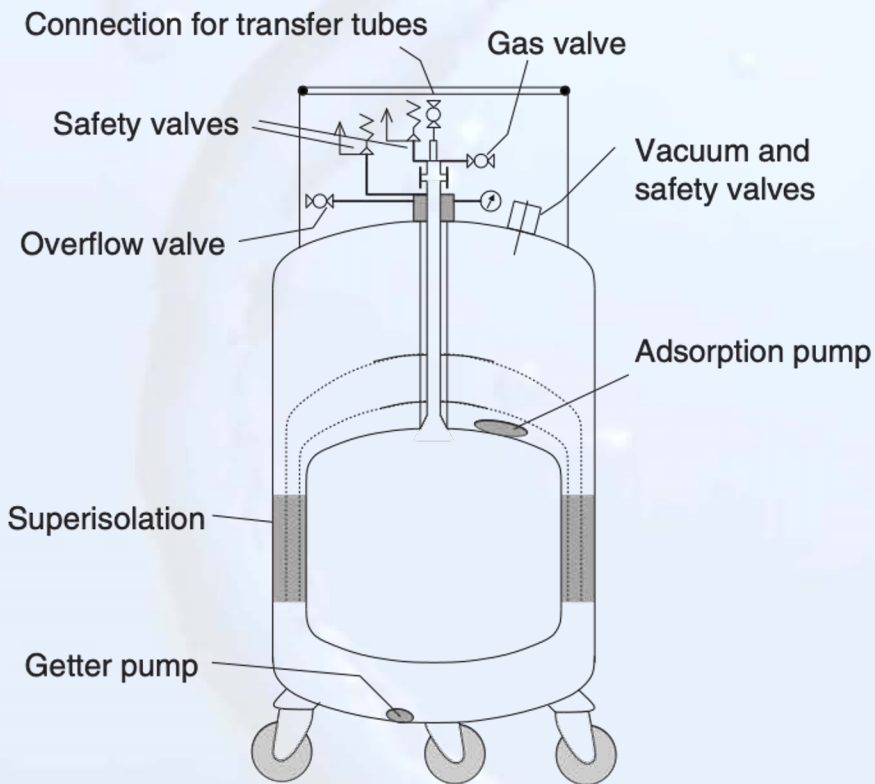


11.1 Bath Cryostats



helium transport vessel

helium transfer tube





Radiation shields – super insulation



multiple radiation shields → smaller steps → reduction of heat flow

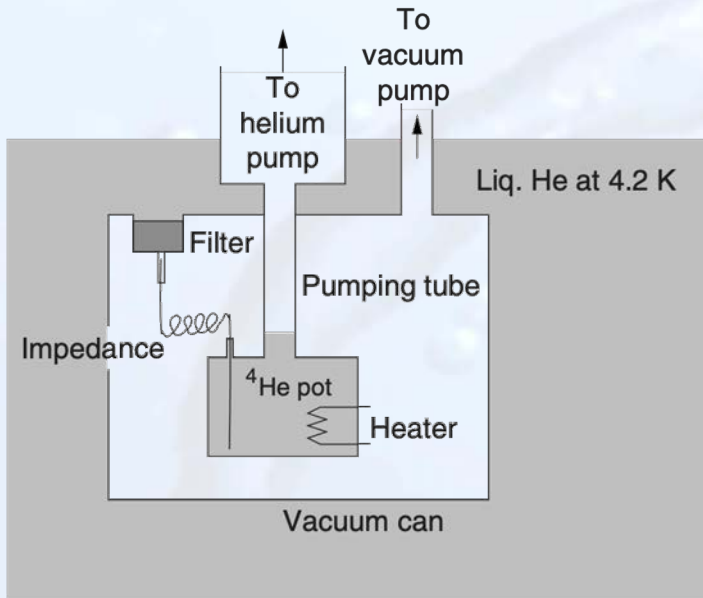
30 to 80 layers of low conductivity
high reflection material → **aluminized Mylar**

apparent thermal conductivity
 $\sim 10^{-4}$ to 10^{-5} W/(m K)





Cryostats with 1-K-Pot



Vapor pressure curve of various cryogenic liquids

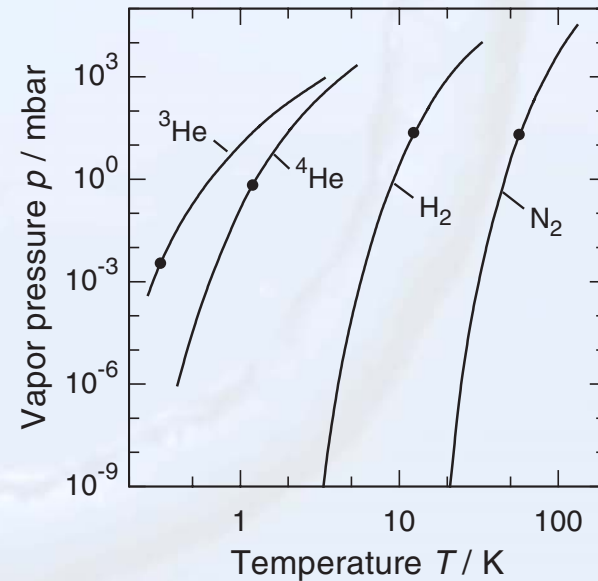
Clausius-Clapeyron equation

$$\frac{dp}{dT} = \frac{L}{\Delta V T} \quad pV_g = RT$$

$\Delta V = V_g - V_l \approx V_g$

$$\frac{dp}{dT} = \frac{L}{RT^2} p \quad \rightarrow \quad p(T) = p_0 e^{-L/RT}$$

vapor pressure curve



$${}^4\text{He} \quad L = 90 \text{ J mol}^{-1}$$

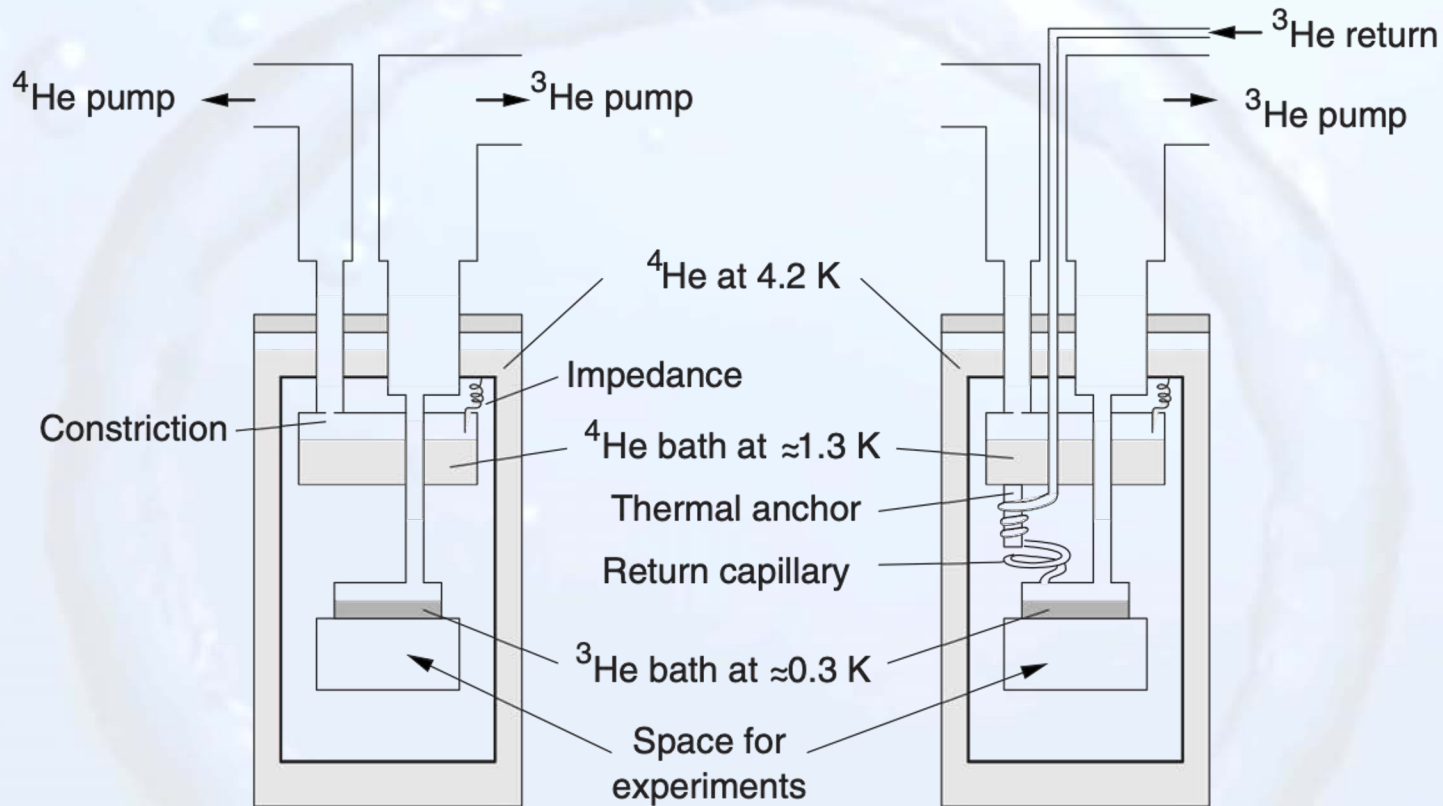
$${}^3\text{He} \quad L = 40 \text{ J mol}^{-1}$$



11.1 Bath Cryostats



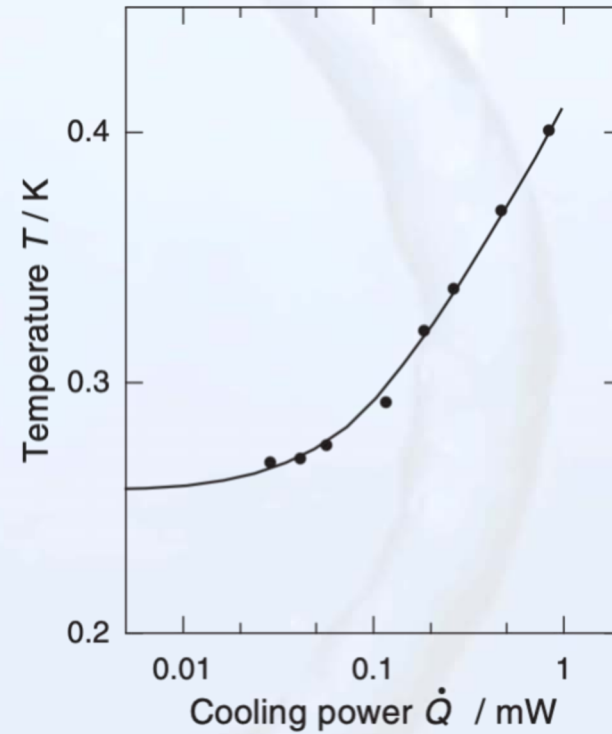
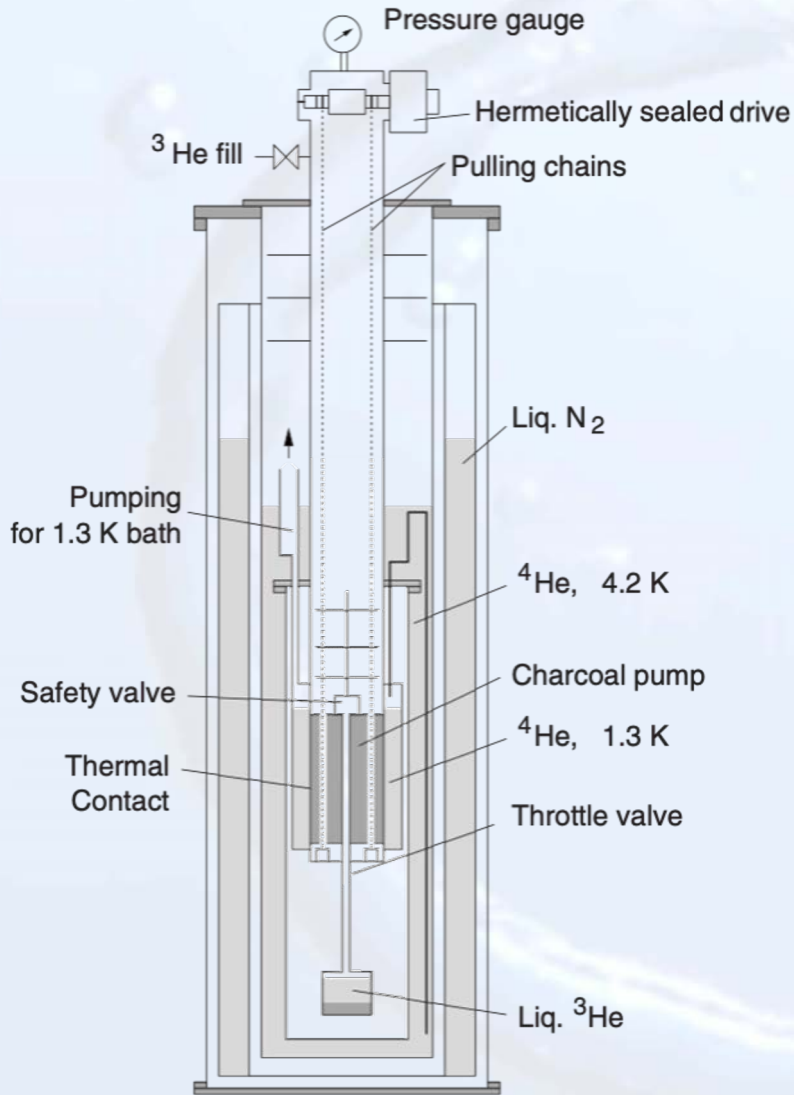
³He cryostats



cooling power $\dot{Q} = \dot{n}_g L \propto p \propto e^{-L/RT}$



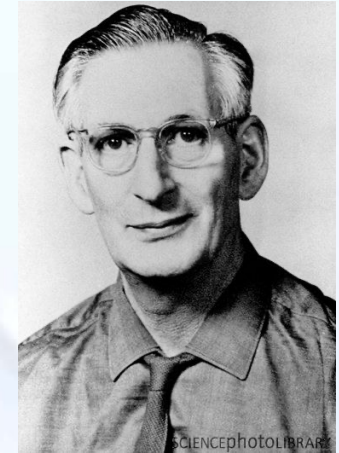
Cooling power of a ^3He cryostat with charcoal absorption pump



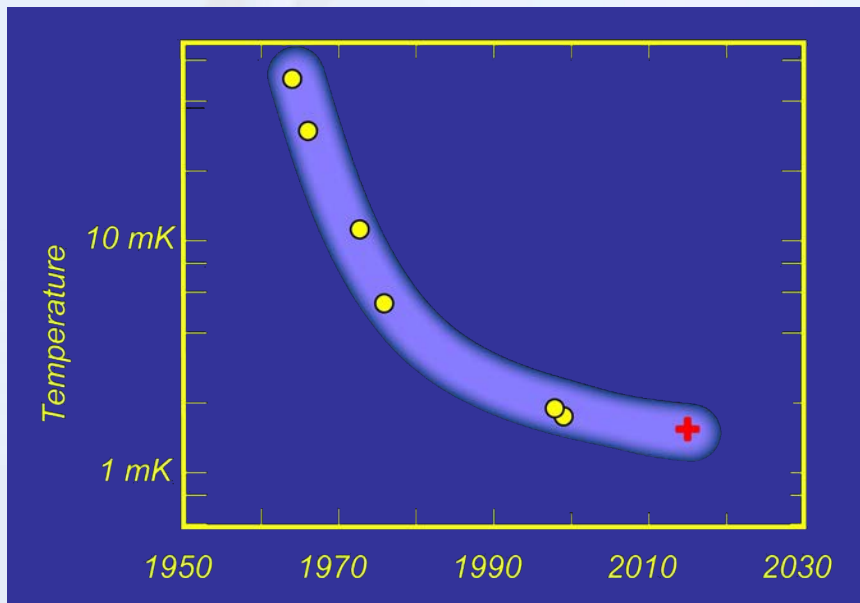


History

- 1951 **basic idea** suggested by Heinz London
- 1962 **detailed concept** worked out by London, Clark, Mendoza
- 1965 **first realization** Das, De Bruyn Ouboter, Taconis $T_{\min} = 220 \text{ mK}$
- 1999 **lowest temperature** obtained, J.C. Cousins *et al.* $T_{\min} = 1.75 \text{ mK}$

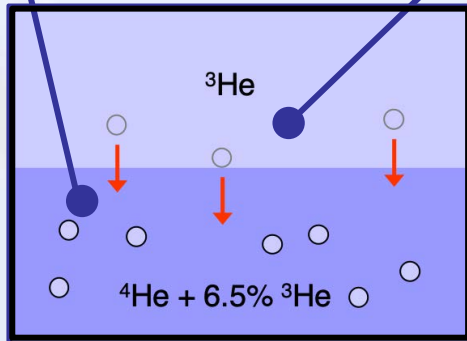
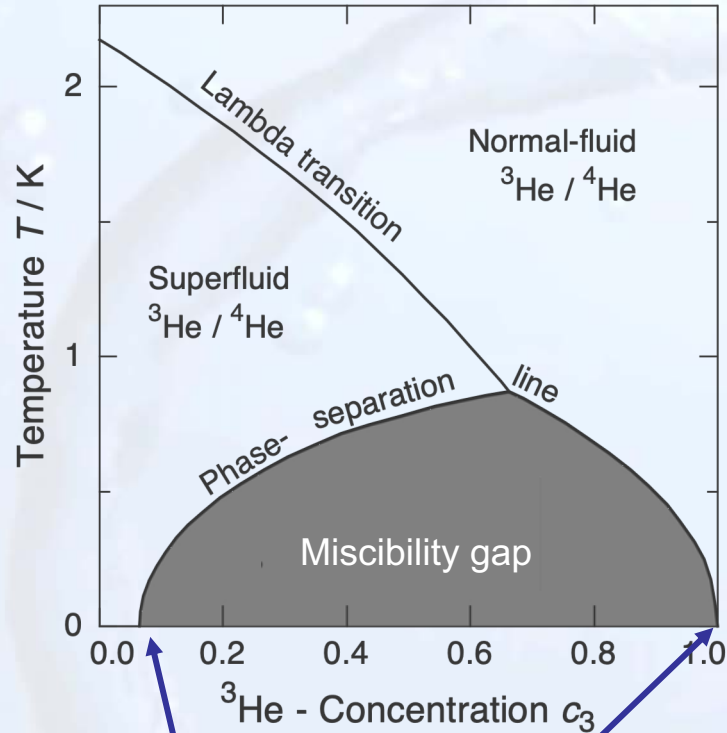


Heinz London





11.2 Dilution Refrigerators

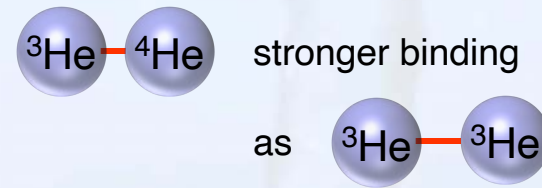


occurrence of miscibility gap

but 6.5% ^3He in ^4He at $T = 0\text{ K}$

reason:

zero-point motion weakens binding



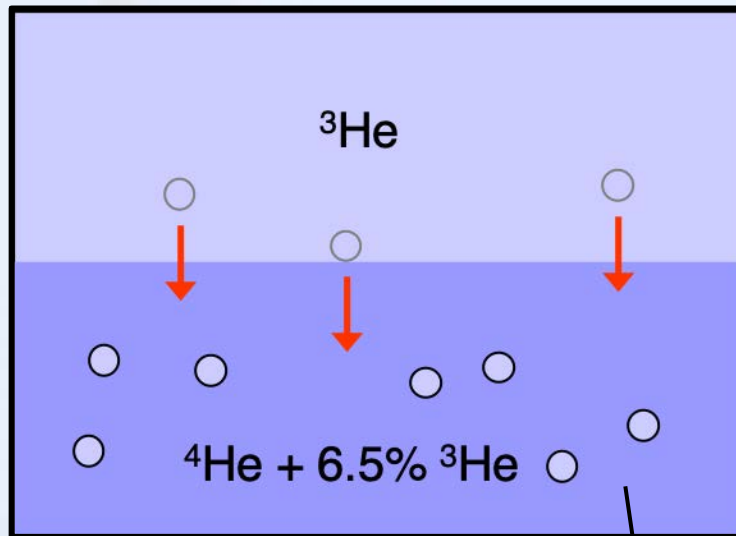
but: Fermi energy

max. 6.5% ^3He in ^4He at $T = 0\text{ K}$



principal of cooling by mixing $^3\text{He}/^4\text{He}$

- ▶ **transition** of ^3He into the ^4He rich phase
- ▶ cooling by „evaporation“ of ^3He into ^4He quasi vacuum



heat of solubility per Mol:

$$\Delta Q = T\Delta S = aT^2$$

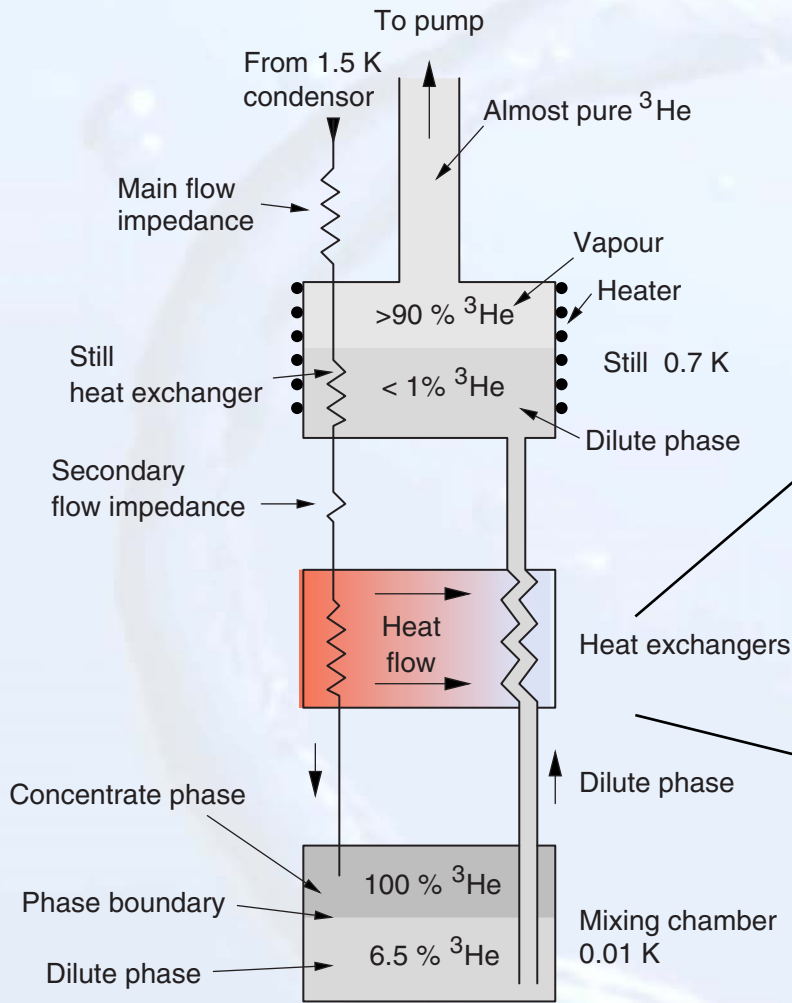
$$a = -84 \text{ J/K}^2$$



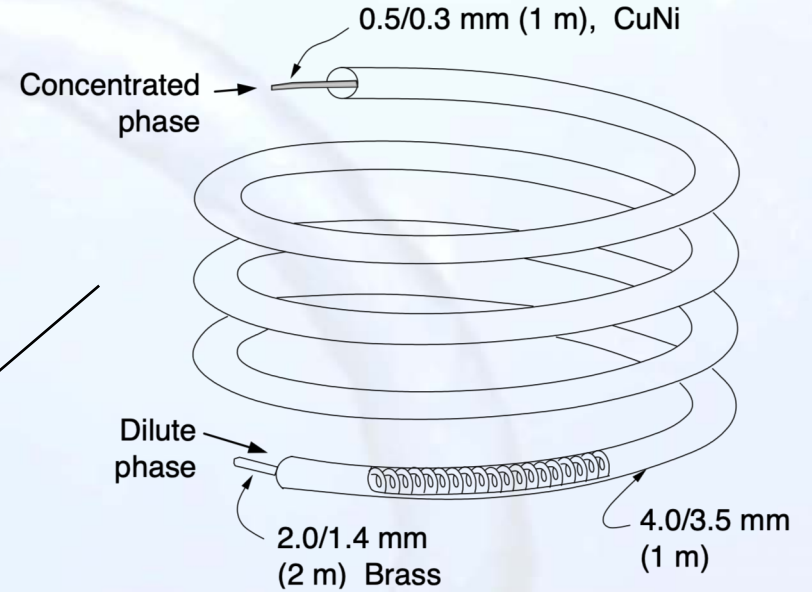
11.2 Dilution Refrigerators



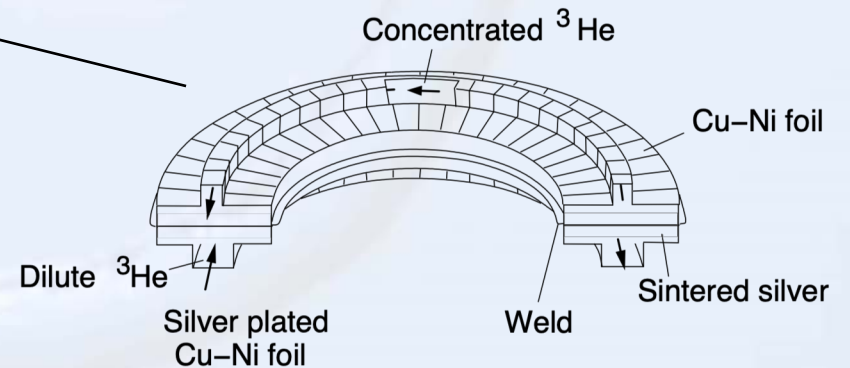
Realisation of $^3\text{He}/^4\text{He}$ cooling cycle



continuous heat exchanger



step heat exchanger





Kapitza Resistance – thermal boundary resistance

Snell's law of refraction $\frac{\sin \alpha_l}{\sin \alpha_s} = \frac{v_l}{v_s}$

critical angle of **total reflection** $\alpha_l^c = \arcsin \left(\frac{v_l}{v_s} \right)$

for liquid helium and copper $\alpha_l^c \approx 4^\circ$

fraction of phonons incident **within critical angle**

$$f = \frac{1}{2} \sin^2 \alpha_l^c = \frac{1}{2} \left(\frac{v_l}{v_s} \right)^2 < 10^{-2}$$

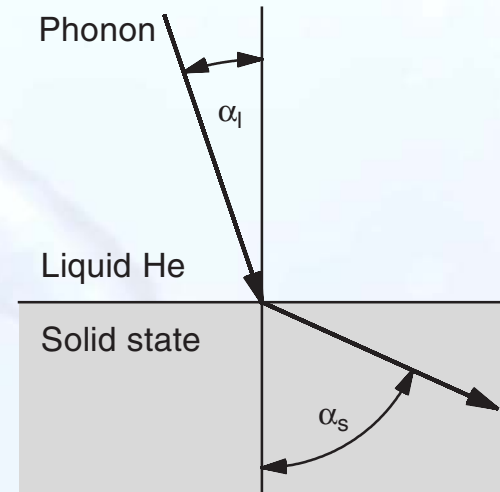
transmission coefficient

$$t = \frac{4Z_l Z_s}{(Z_l + Z_s)^2} \approx \frac{4Z_l}{Z_s} = \frac{4\rho_l v_l}{\rho_s v_s}$$

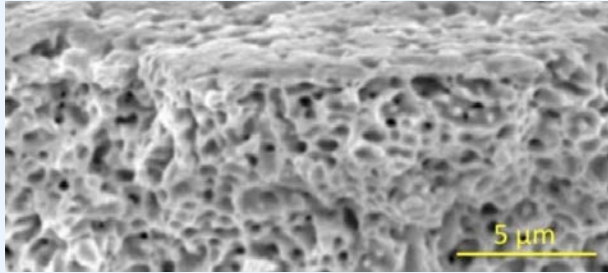
$$Z_l = \rho_l v_l \quad Z_s = \rho_s v_s \quad \text{acoustic impedances}$$

fraction of phonons crossing the interface

$$ft = \frac{2\rho_l v_l^3}{\rho_s v_s^3}$$



- ▶ Kapitza resistance occurs at **any solid-solid, liquid-solid** interface
- ▶ particular problematic for **liquid helium** because of the **low** sound velocity
- ▶ helium-copper $ft < 10^{-5}$



silver sinter
SEM image

heat flow from liquid to solid (using Debye model)

$$\dot{Q} = \frac{1}{2} f t u v_\ell A = \frac{\pi^2 k_B^4 \rho_\ell v_\ell}{30 \hbar^3 \rho_s v_s^3} AT^4$$

$u = U/V = \pi^2 k_B^4 T^4 / (30 \hbar^3 v_\ell^3)$

in equilibrium **identical** heat flow from solid to liquid

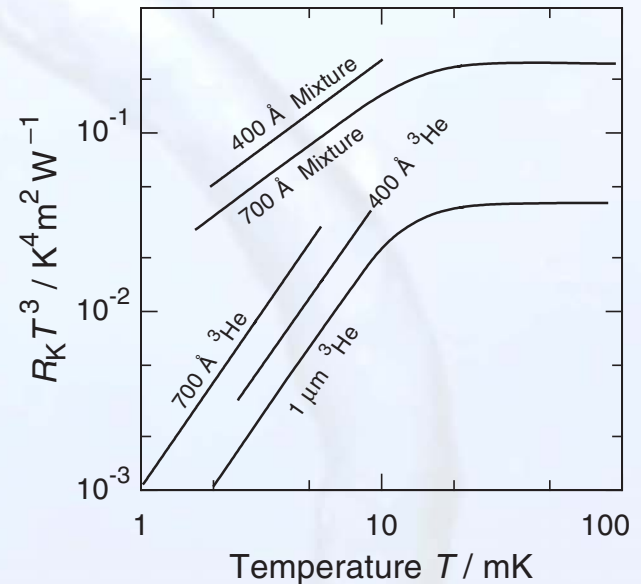
net flow in non-equilibrium (ΔT)

$$\dot{Q} = \frac{d\dot{Q}}{dT} \Delta T = \frac{2\pi^2 k_B^4 \rho_\ell v_\ell}{15 \hbar^3 \rho_s v_s^3} AT^3 \Delta T$$

Kapitza resistance

$$R_K = \frac{A \Delta T}{\dot{Q}} = \frac{15 \hbar^3 \rho_s v_s^3}{2 \pi^2 k_B^4 \rho_\ell v_\ell} \frac{1}{T^3}$$

Kapitza resistance between **pure ^3He** and **$^3\text{He}/^4\text{He}$ mixtures** and **silver sinters** of **different grain sizes**



- ▶ $20 \text{ mK} < T < 100 \text{ mK}$ good agreement with Debye model $R_K \propto T^{-3}$
- ▶ below 20 mK $R_K \propto T^{-2}$ or $R_K \propto T^{-1}$

→ **anomalous** Kapitza resistance

origin: **TLS**, coupling to **zero and second sound** modes, phonon **wavelength** larger than sinter grains



Cooling power

assuming **100% ³He circulation** one finds in equilibrium:

$$\dot{Q}_{mc} + \dot{N}_3 [H_3(T_{ex}) - H_3(T_{mc})] = \dot{N}_3 [H_{3,d}(T_{mc}) - H_3(T_{mc})]$$

\dot{Q}_{mc} : heat leak and/or available cooling power
 \dot{N}_3 : circulation rate
 $H_3(T_{ex})$: enthalpy of ³He-rich phase
 $H_3(T_{mc})$: enthalpy of ³He-dilute phase
 $H_{3,d}(T_{mc})$: enthalpy of ³He-dilute phase
 T_{mc} : mixing chamber temperature
 T_{ex} : temperature after last heat exchanger

enthalpy
 $H = U + pV$

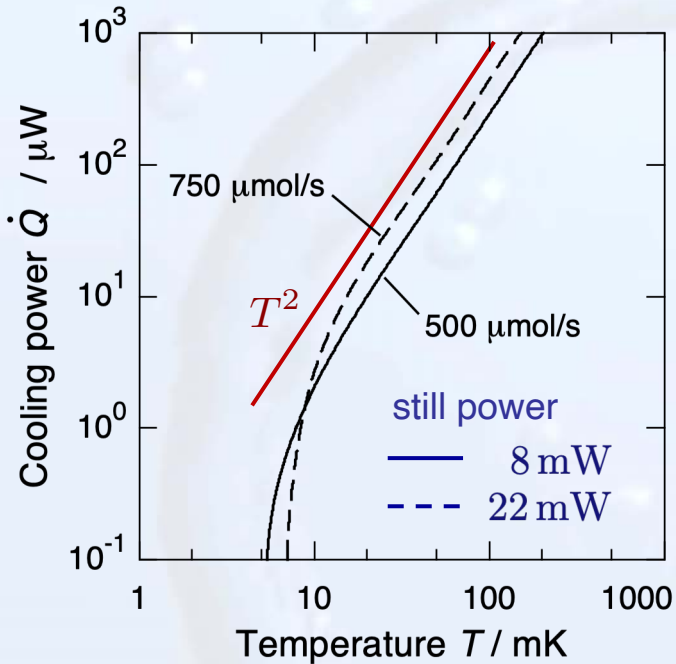
inserting the enthalpies

$$\dot{Q}_{mc} = \dot{N}_3 [H_{3,d}(T_{mc}) - H_3(T_{ex})]$$

$$= \dot{N}_3 (95 T_{mc}^2 - 11 T_{ex}^2) \left(\frac{\text{J}}{\text{mol K}^2} \right)$$



Temperature and circulation rate dependence of the cooling power



limiting case of vanishing cooling power: $\dot{Q}_{mc} = 0$

$$95 T_{mc}^2 - 11 T_{ex}^2 = 0$$

$$\frac{T_{ex}}{T_{mc}} = 2.8$$

→ this underlines the **importance** of the **heat exchanger** quality

▶ for $\dot{Q} \gg \dot{Q}_{heat\ leak}$ → $\dot{Q} \propto T^2$, $\dot{Q} \propto \dot{N}_3$

▶ heat leak determines lowest temperature

└─┬─┘
circulation rate



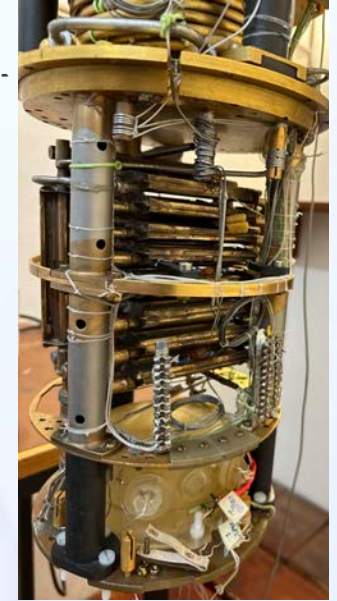
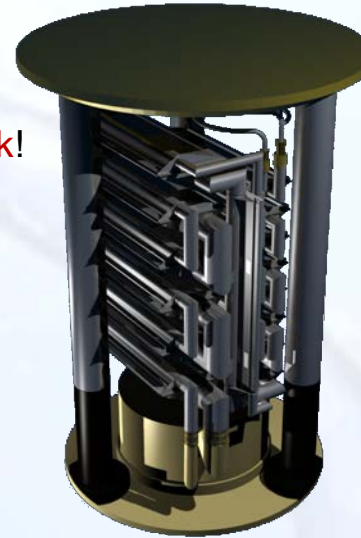
Minimum temperature

- ▶ there is **no principle limit** ... it is determined by the **heat leak!**
- ▶ unavoidable heat leak: viscous friction of ^3He

pressure difference along the heat exchanger:

$$\Delta p = G\eta\dot{V} \quad \text{Hagen-Poiseuille law}$$

$$G = 8L/(\pi r^4)$$

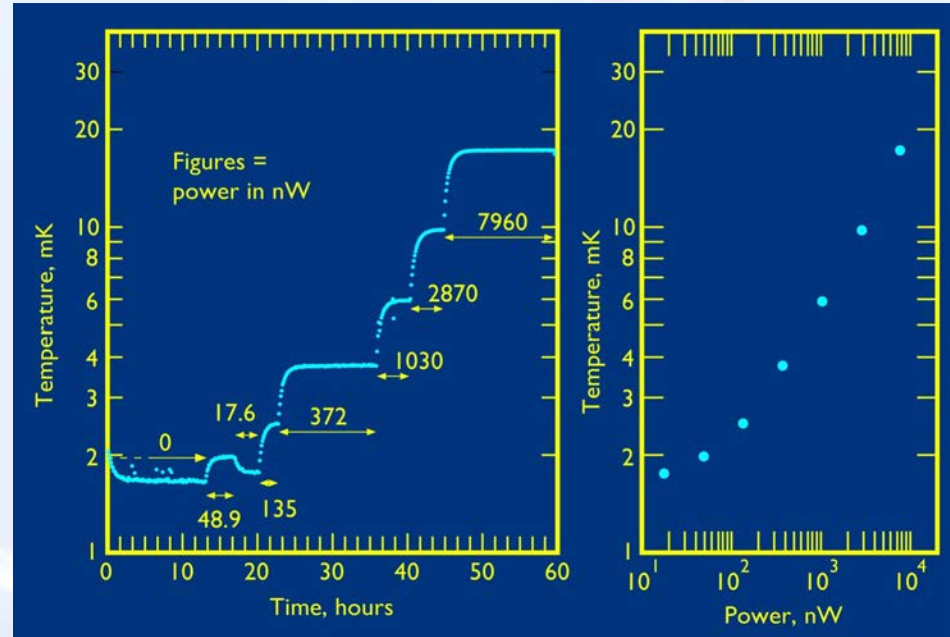


heat leak due to **viscous friction**

$$\dot{Q}_{\text{visc}} = \dot{V}\Delta p = G\eta\dot{V}^2$$

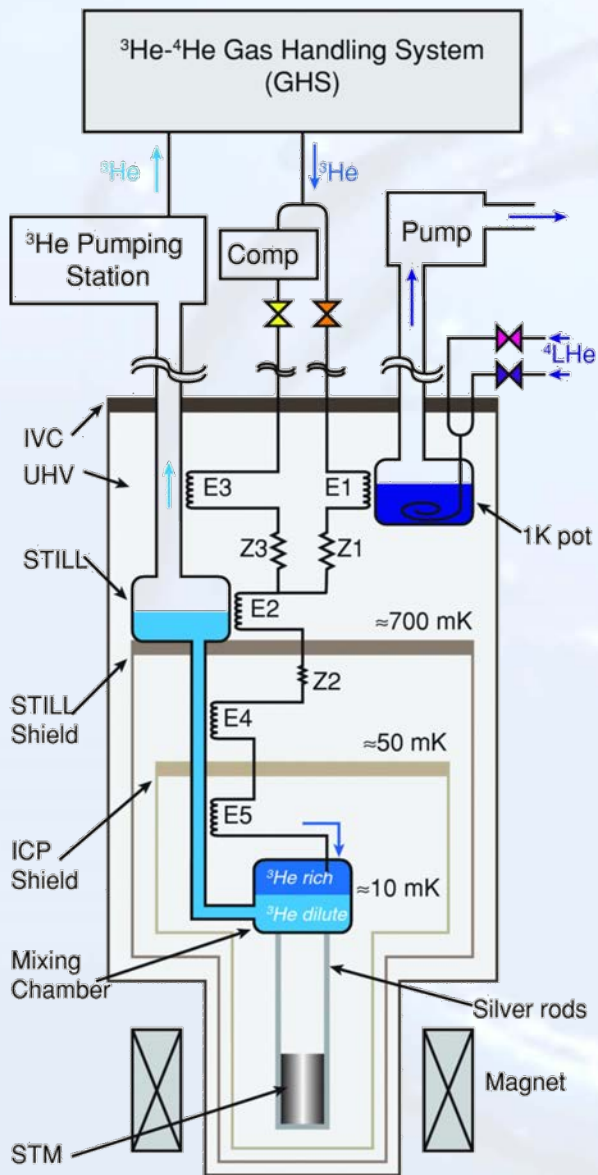
single shot minimum temperature

$$T_{\text{min.}} = \frac{4}{\sqrt[3]{d}} \text{mK}(\text{mm})^{1/3}$$



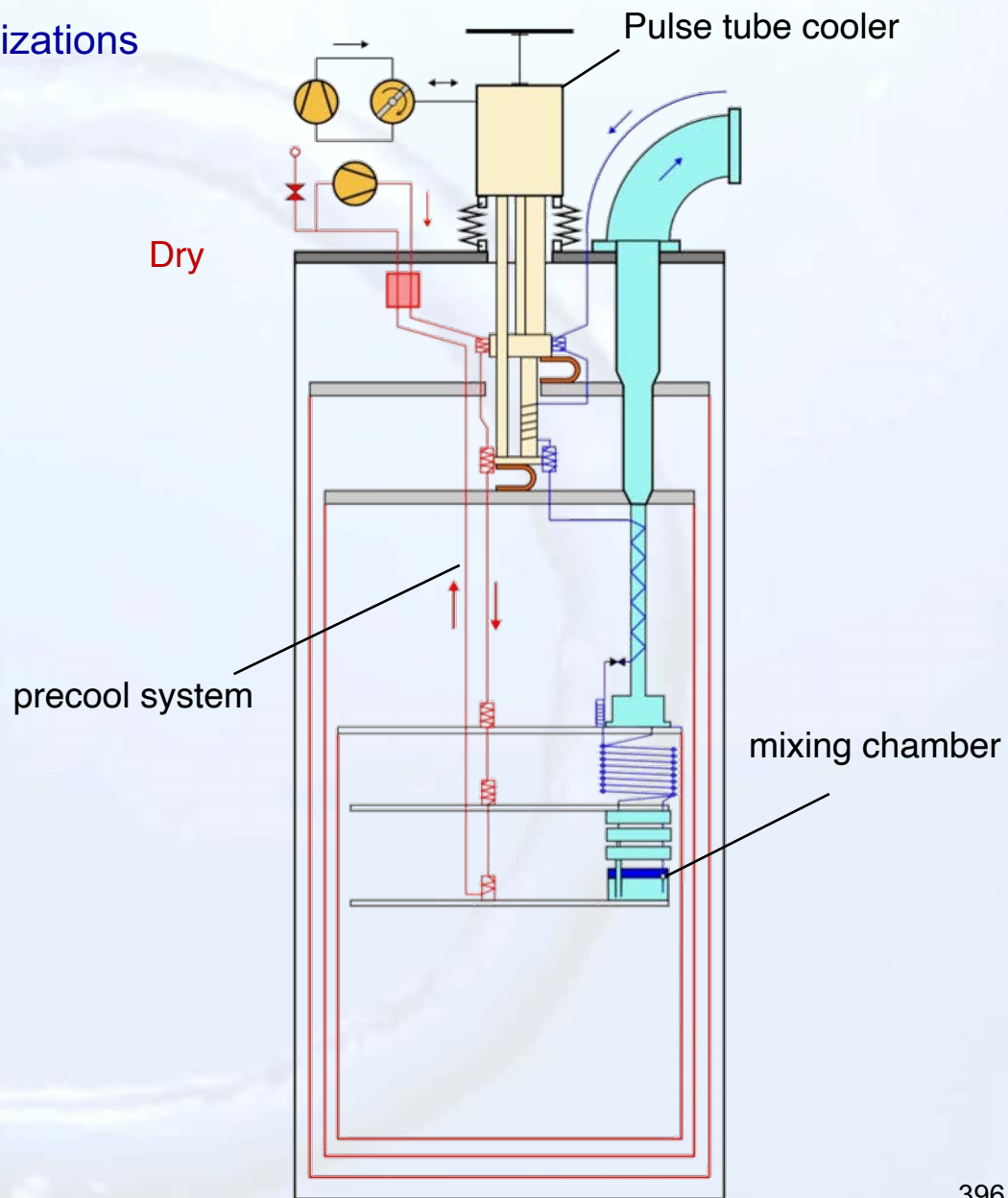


Realizations



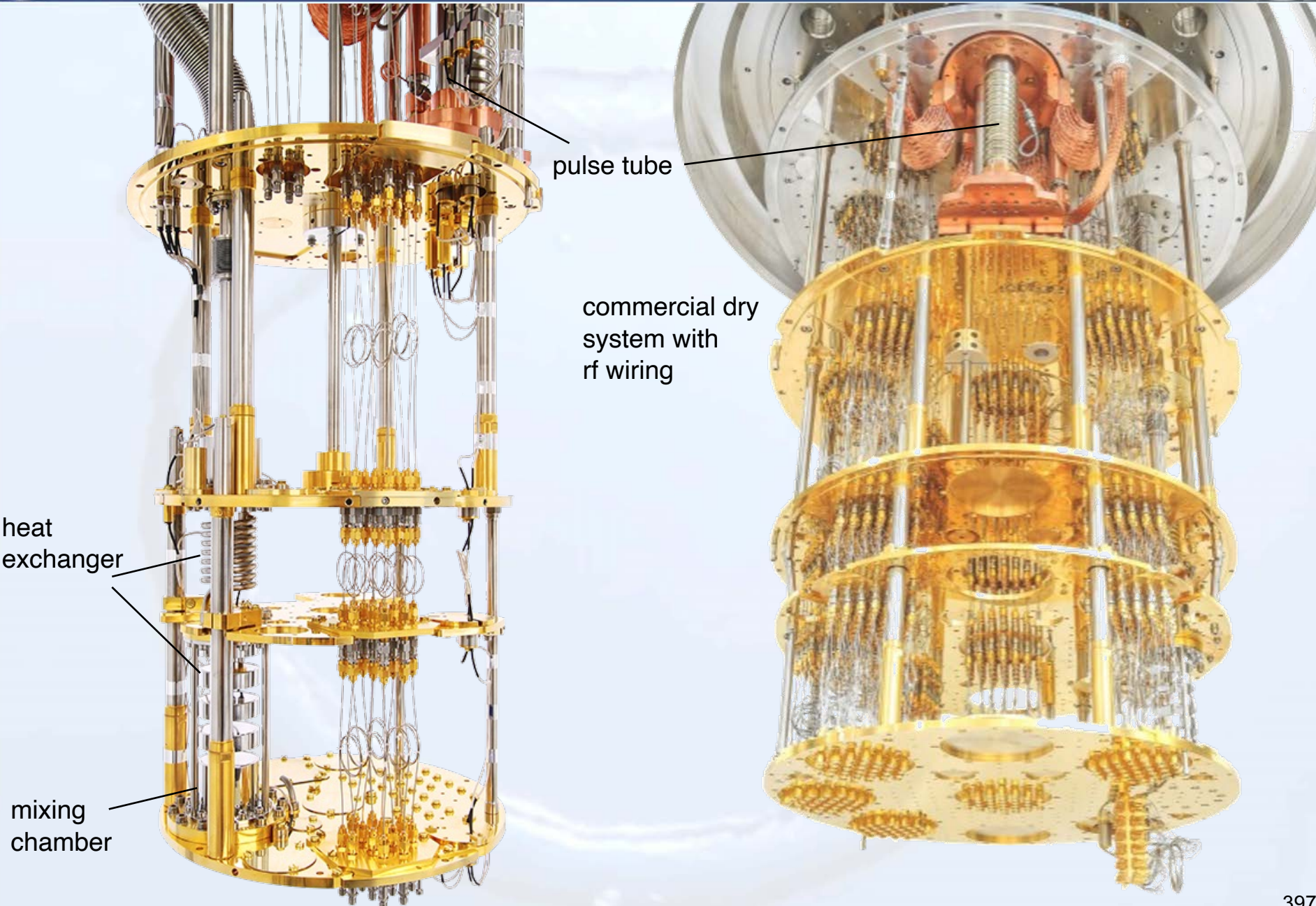
Wet

Dry





11.2 Dilution Refrigerators



pulse tube

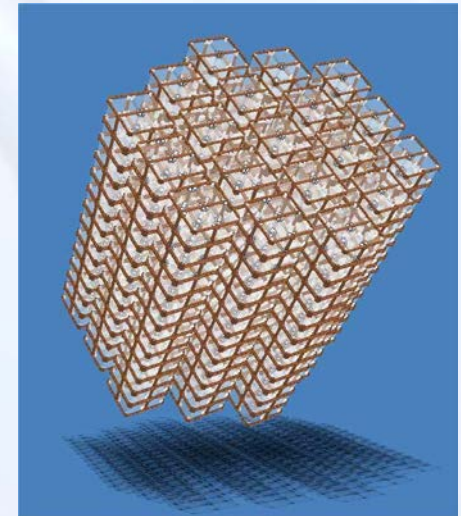
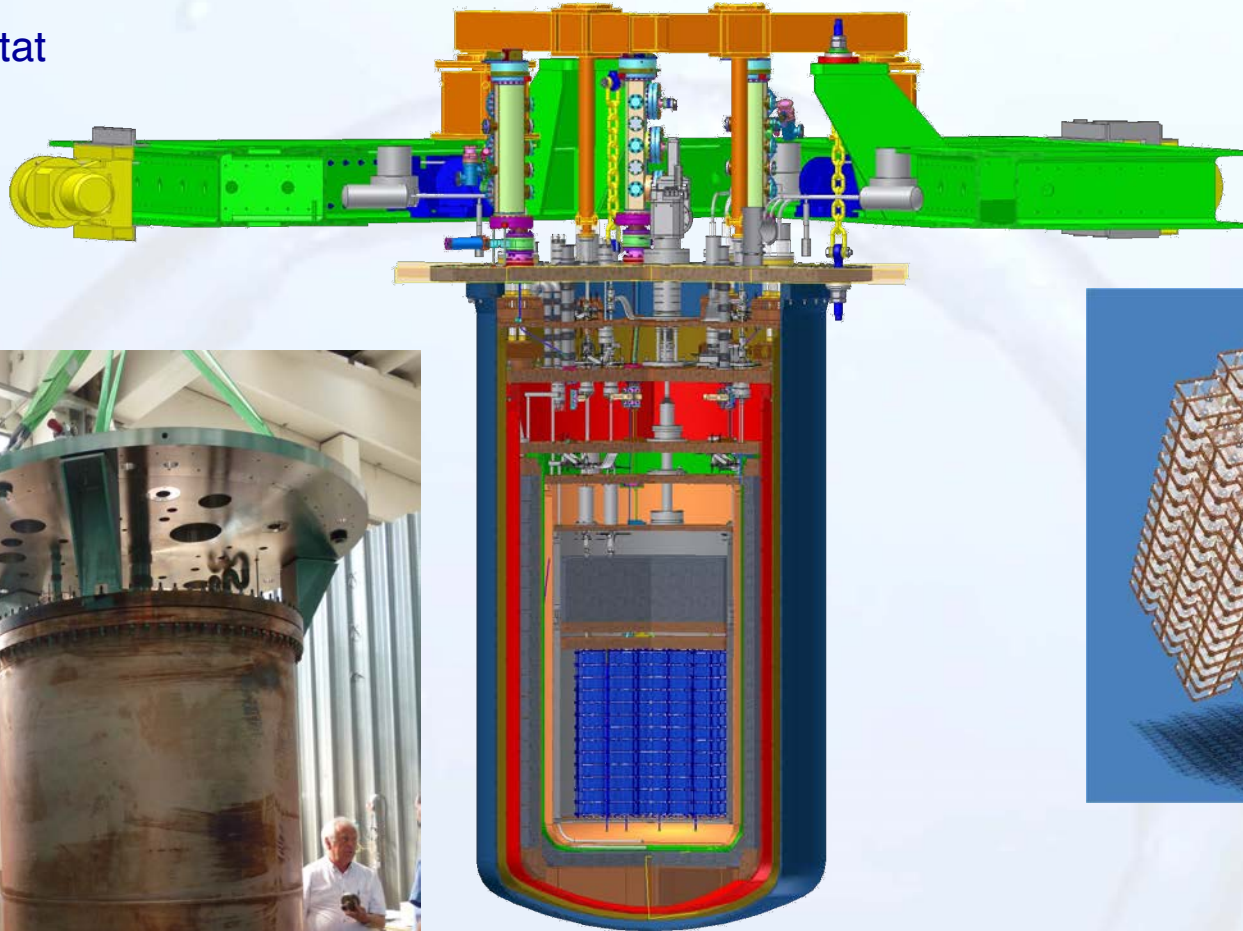
commercial dry
system with
rf wiring

heat
exchanger

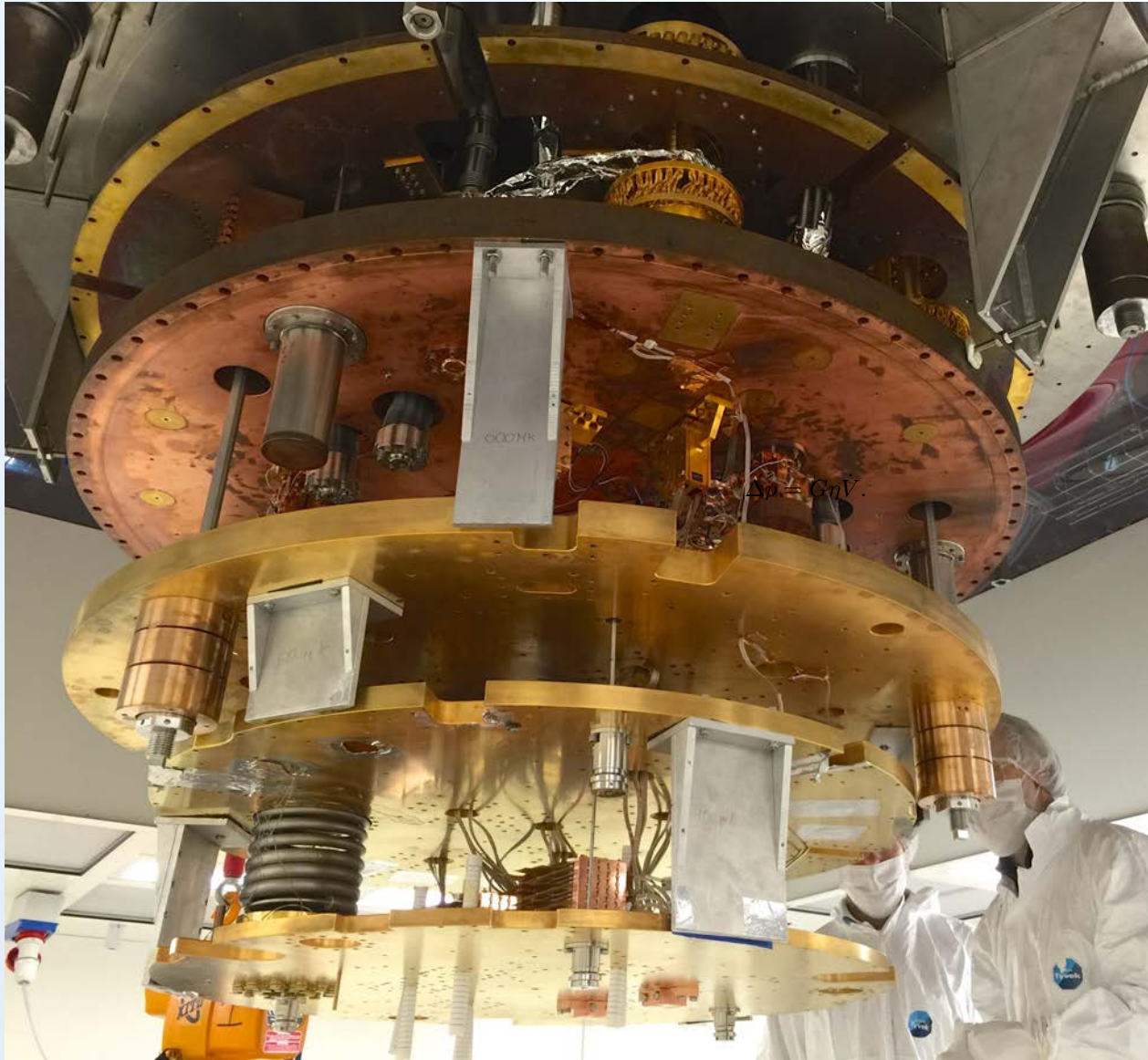
mixing
chamber



Cuore Cryostat



- ▶ 750 kg of TeO_2 cooled by the mixing chamber to 10 mK
- ▶ made of materials with low level radioactivity materials
- ▶ two cold shields made of ancient roman lead 10 cm thick
- ▶ 5 Pulse Tube Coolers



Cuore Cryostat

$$\Delta\rho = C_n \dot{V}$$



11.3 Adiabatic Demagnetization Refrigerators



1926 basic idea suggested by Debye, 1927 Giauque

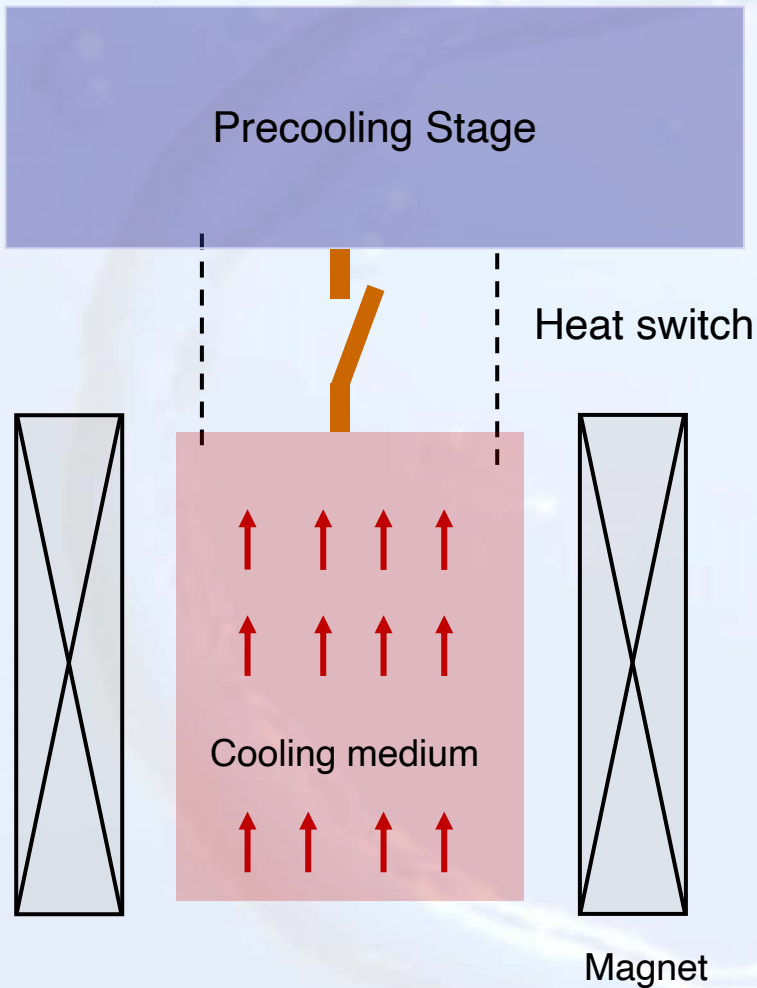
1933 first realization by two groups Leiden, Berkeley

↙ **electronic** spins

↘ **nuclear** spin, Gorter 1934, Kurti and Simon 1935



General cooling principle



▶ precooling

▶ isothermal magnetisation

$$\Delta Q_{\text{mag}} = -T_i [S(B_i, T_i) - S(0, T_i)]$$

▶ thermal isolation

▶ heat switch opened

▶ adiabatic demagnetisation

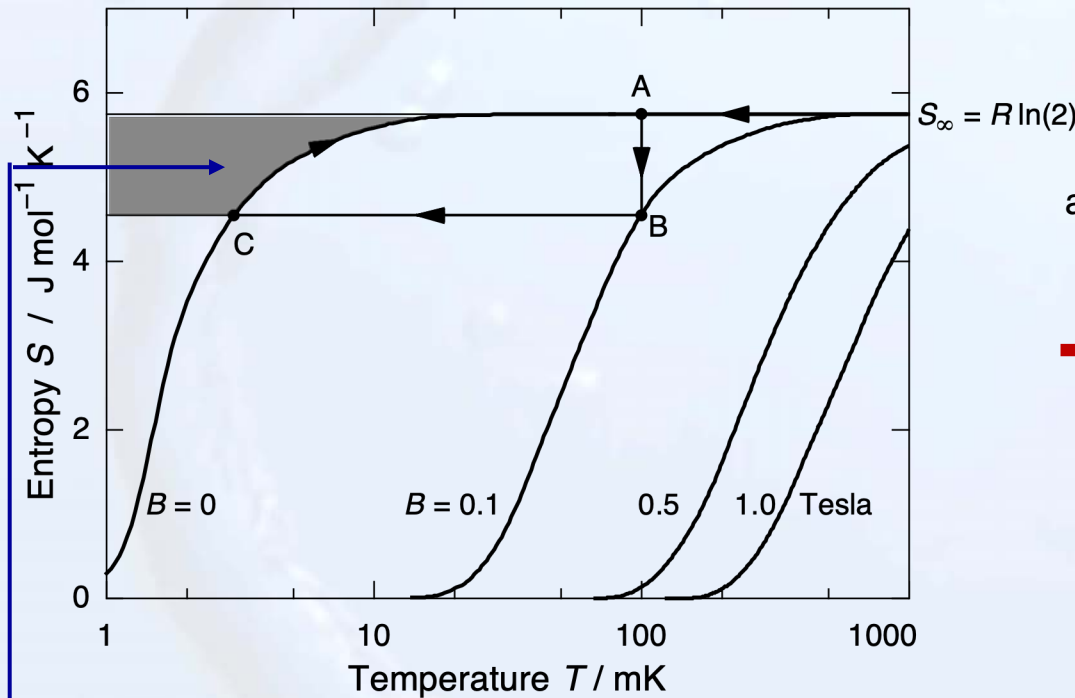
$$S = S\left(\frac{B}{T}\right) = \text{const.}$$



Entropy: ideal system, no heat leaks

internal field caused by spin-spin interaction

$$S = Nk_B \left\{ \ln(2J + 1) - \frac{g^2 J(J + 1) \mu_B^2}{6 k_B^2} \frac{B^2 + B_{\text{int}}^2}{T^2} \right\}$$



adiabatic demagnetization $S_B = S_C$

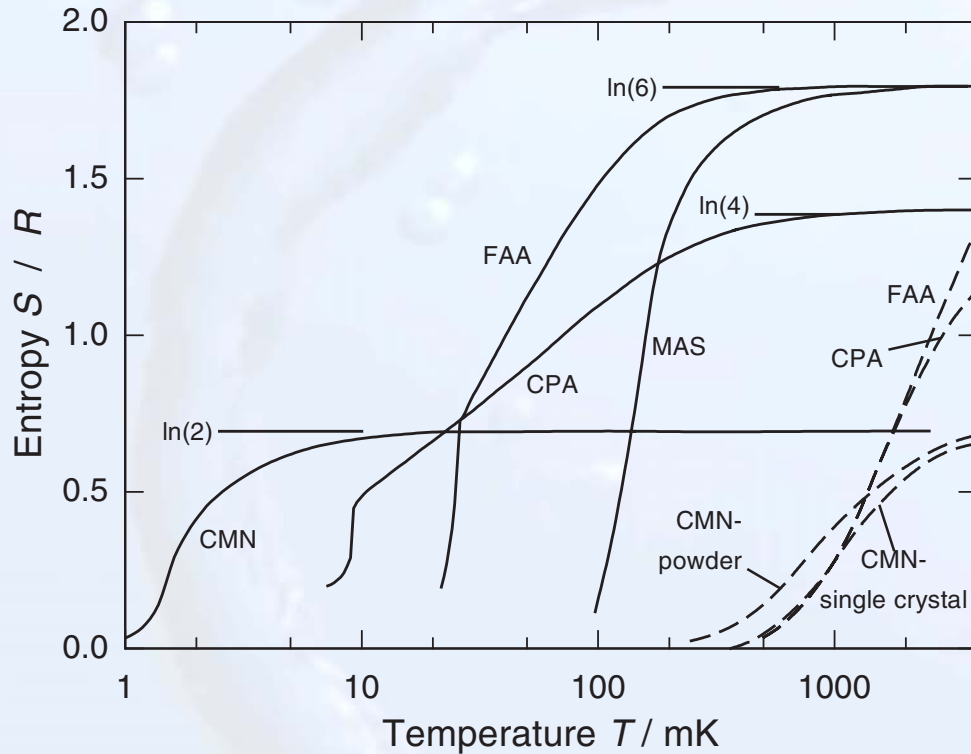
$$\rightarrow T_f = T_i \sqrt{\left(\frac{B_f^2 + B_{\text{int}}^2}{B_i^2 + B_{\text{int}}^2} \right)}$$

$$\Delta Q_{\text{spin}}(B_C) = \int_{T_C}^{T_A} C_{\text{spin}} dT = \int_{T_C}^{T_A} T \left(\frac{\partial S}{\partial T} \right)_{B_C} dT$$



a) Electronic spins

entropy of different paramagnetic salts



MAS for $\text{MnSO}_4 \cdot (\text{NH}_4)_2\text{SO}_4 \cdot 6\text{H}_2\text{O}$

FAA for $\text{Fe}_2(\text{SO}_4)_3 \cdot (\text{NH}_4)_2\text{SO}_4 \cdot 24\text{H}_2\text{O}$

CPA for $\text{Cr}_2(\text{SO}_4)_3 \cdot \text{K}_2\text{SO}_4 \cdot 24\text{H}_2\text{O}$

CMN for $2\text{Ce}(\text{NO}_3)_3 \cdot 3\text{Mg}(\text{NO}_3)_2 \cdot 24\text{H}_2\text{O}$

problems with paramagnetic salts

- T_c relatively high
- low thermal conductivity



high conductive wires to improve
low thermal conductivity of salt pills

NASA GSFC

- ▶ FAA salt pill for space application
- ▶ 15.000 gold wires
- ▶ salt pill grown around the wires





b) Nuclear spins

- ▶ metals with **fast relaxation time**
- ▶ nuclei with **large magnetic moment**
- ▶ isotopes with **large natural abundance**
- ▶ **cubic structure** to avoid quadropole contributions
- ▶ **no** superconductor
- ▶ **pure** material, **easy to machine**

$$\frac{dT_n^{-1}}{dt} = - \frac{(T_n^{-1} - T_e^{-1})}{\tau_1}$$

$\tau = \kappa/T_e$ Korringa relation

	Structure	I	μ/μ_N	κ (K s)	Abundance (%)
^{63}Cu	fcc	3/2	2.22	1.27	69.1
^{65}Cu	fcc	3/2	2.38	1.09	30.9
^{195}Pt	fcc	1/2	0.597	0.03	33.8
$^{141}\text{PrNi}_5$	fcc	5/2	4.28	<0.001	100

van Vleck paramagnet



Gas gap heat switch



exchange gas

→ pumping to open switch

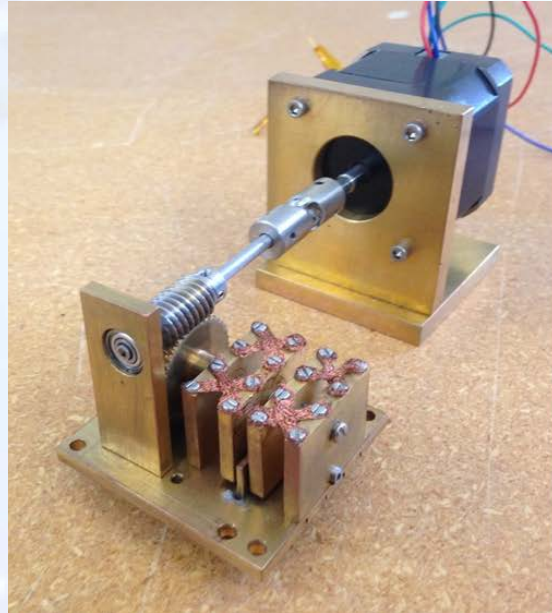
^4He : superfluid layer → creep

H_2 : ortho-para conversion

^3He : no exothermic reaction
no creep
high vapor pressure

→ ideal exchange gas

Mechanical heat switch

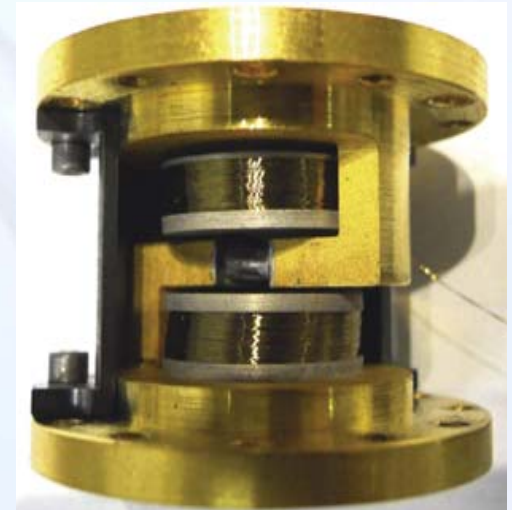


▶ large force needed ~ 100 N

▶ closed: mW/K ... 1 W/K @ 15K

▶ problem: heating on opening

Superconducting heat switch



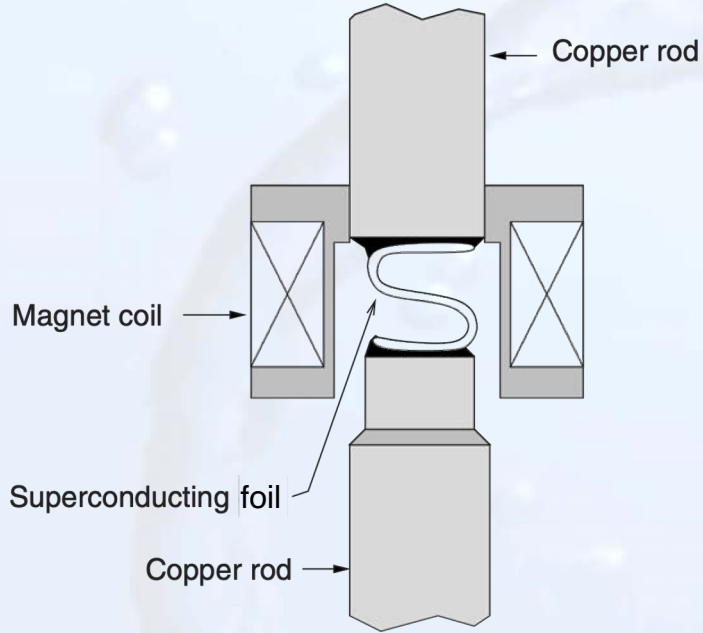
▶ only good well below T_c

▶ open means low conductivity

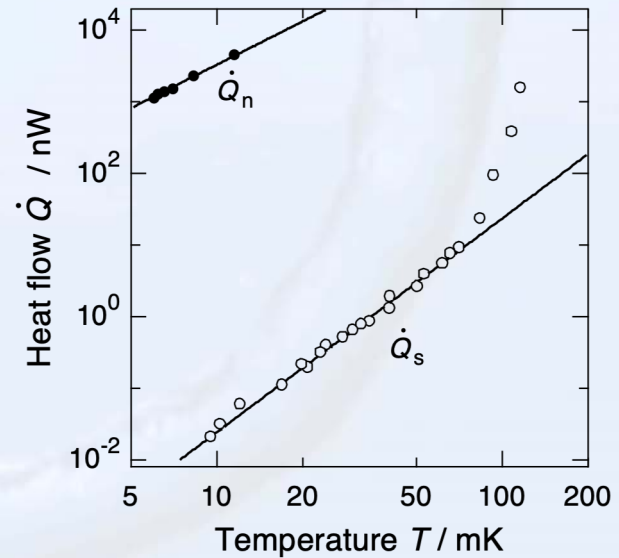
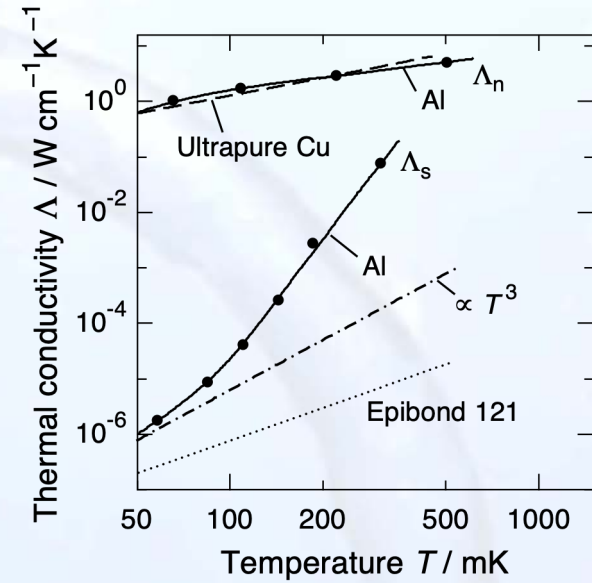
▶ problems: eddy currents
flux trapping



Performance of superconducting heat switch



- ▶ switching ratio 10^6 at 10 mK
- ▶ heat leak in open state 10 pW



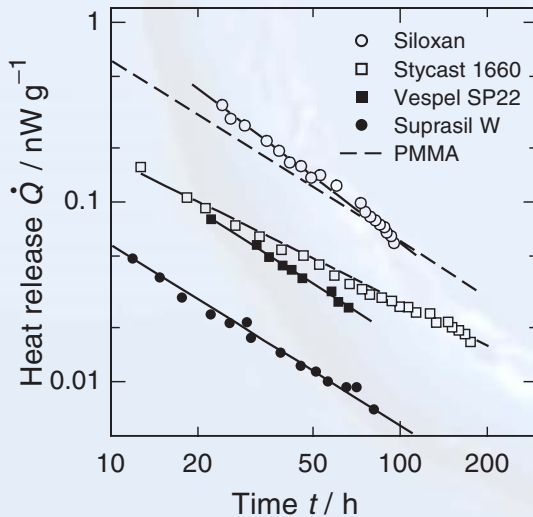


Heat leaks

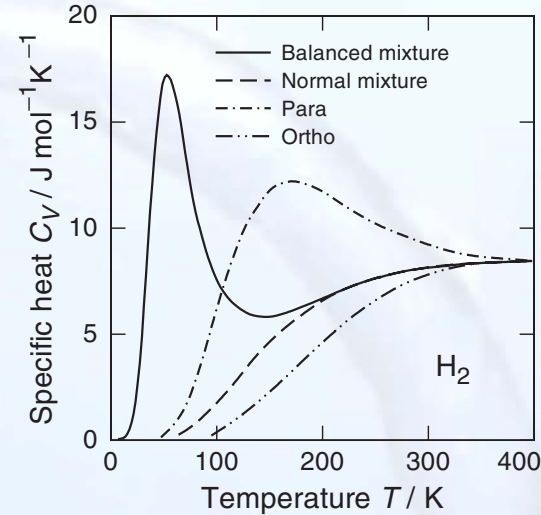
- ▶ eddy current heating $\dot{Q}_{\text{eddy}} = f \frac{V \dot{B}^2}{\rho}$
 - ▶ em fields and vibrations
 - ▶ ortho-para conversion
 - ▶ radioactive impurities
 - ▶ tunneling systems
- } time dependent heat leaks

atomic tunneling systems

$$\dot{Q} = \frac{\pi^2 k_B^2}{24} P_0 (T_1^2 - T_0^2) \frac{1}{t}$$



specific heat of H₂



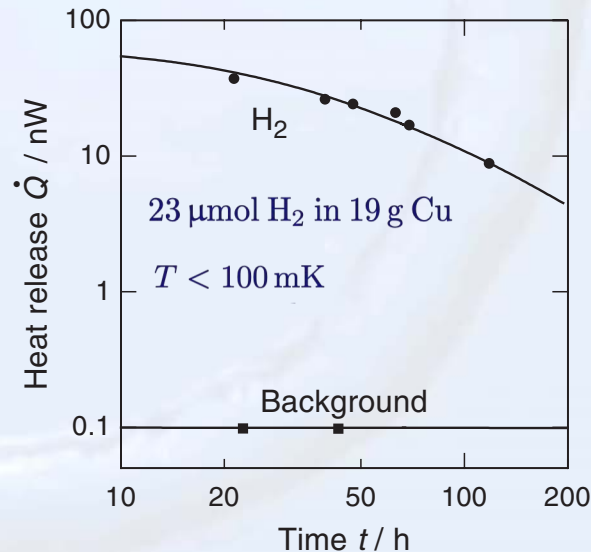
singlet state

para H₂ $j = 0$

triplet state

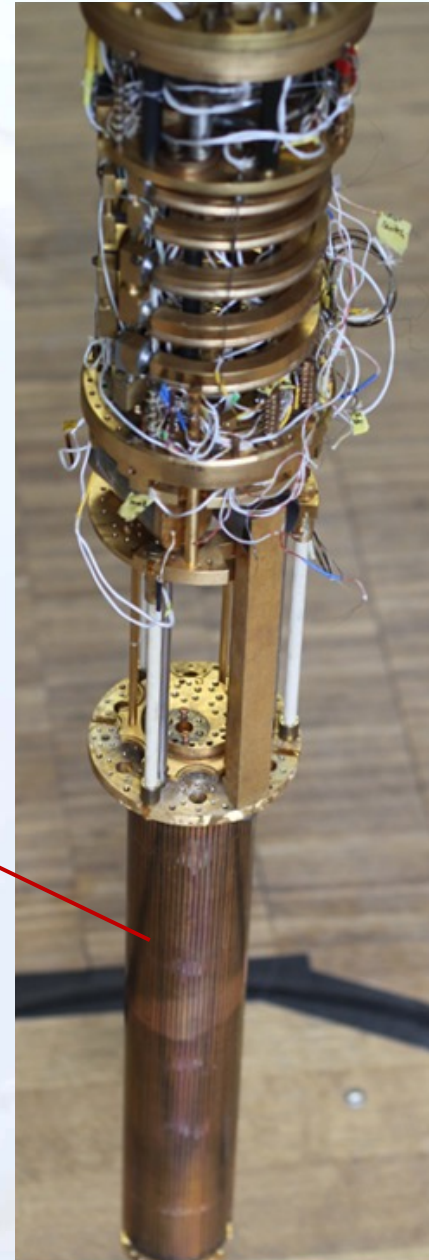
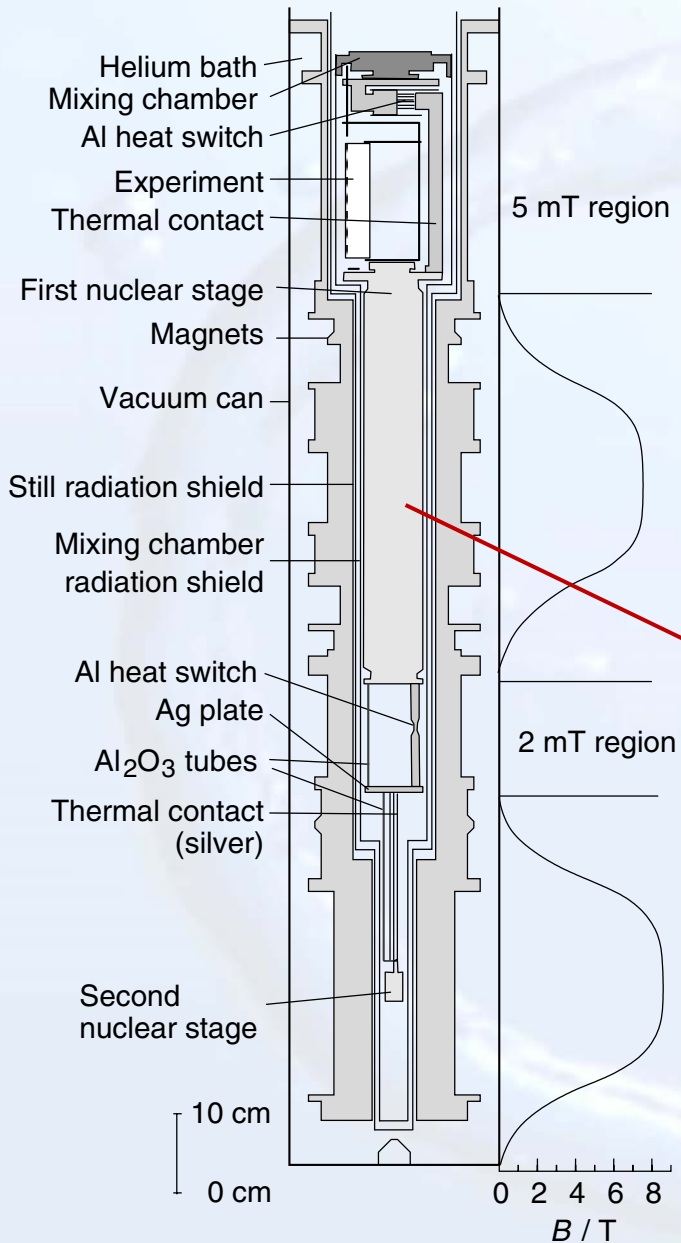
ortho H₂ $j = 1$

H₂ ortho-para conversion





11.3 Adiabatic Demagnetization Refrigerators





Cooling process

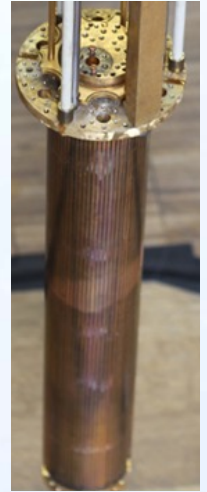
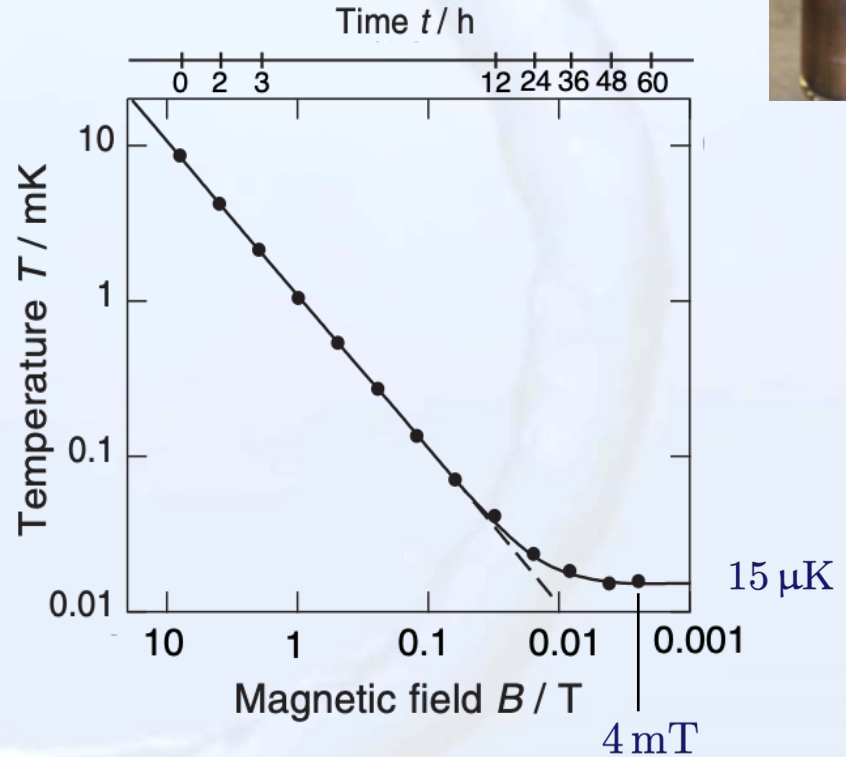
- ▶ precooling to T_A and isothermal magnetization

nuclear Curie constant $\lambda_n = \frac{nI(I+1)\mu_0\mu_n^2g_n^2}{3k_B}$

→ $Q = nT_A\Delta S = -\frac{\lambda_n B_i^2}{2\mu_0 T_A}$

- ▶ reducing B in steps to optimal final field

$B_{f,opt} = \sqrt{\frac{3k_B\kappa\dot{Q}}{ng_n^2I(I+1)\mu_n^2}}$ heat leak

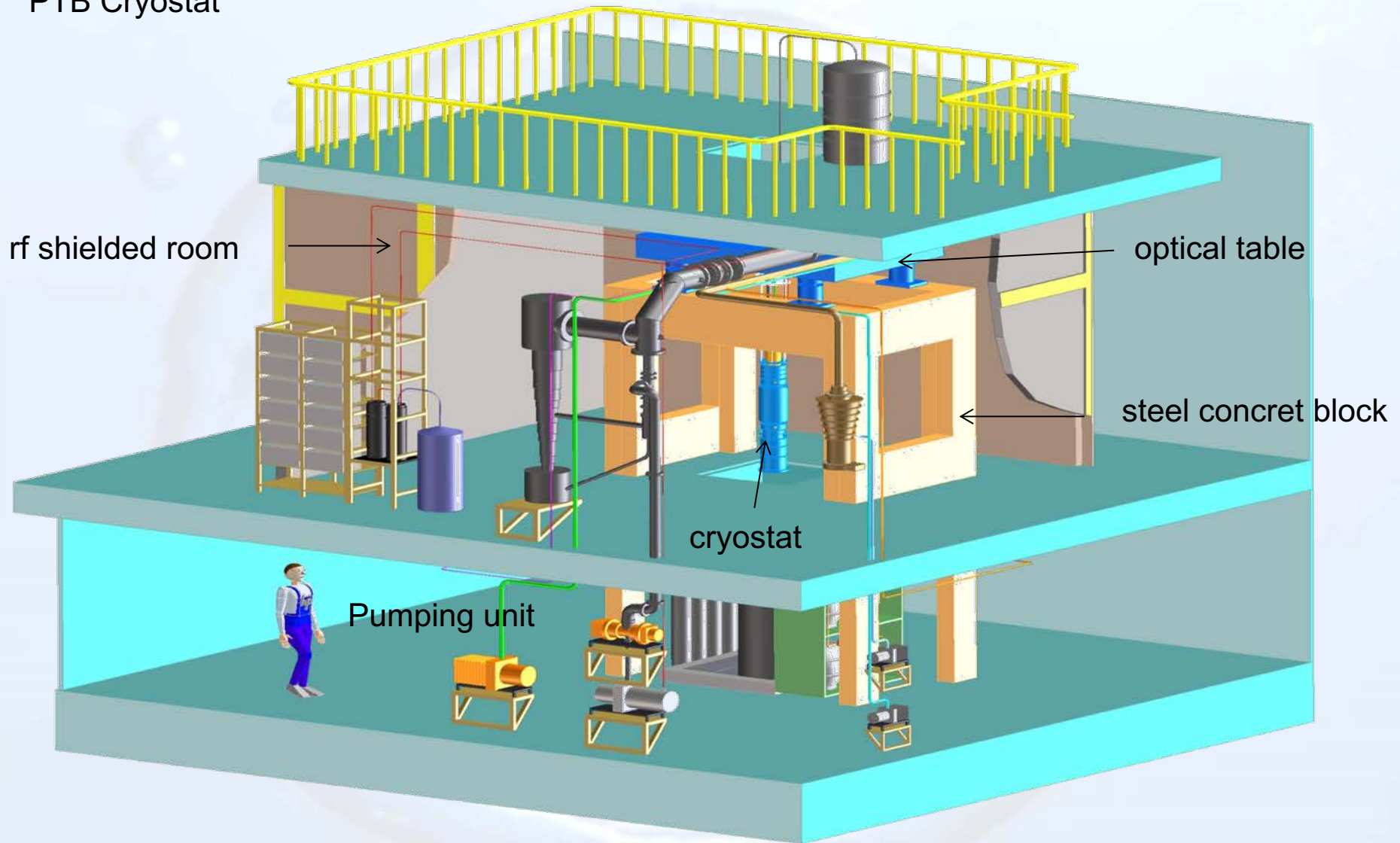




11.3 Adiabatic Demagnetization Refrigerators



PTB Cryostat

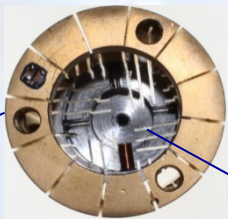




heat switch

Cu stage

heat switch

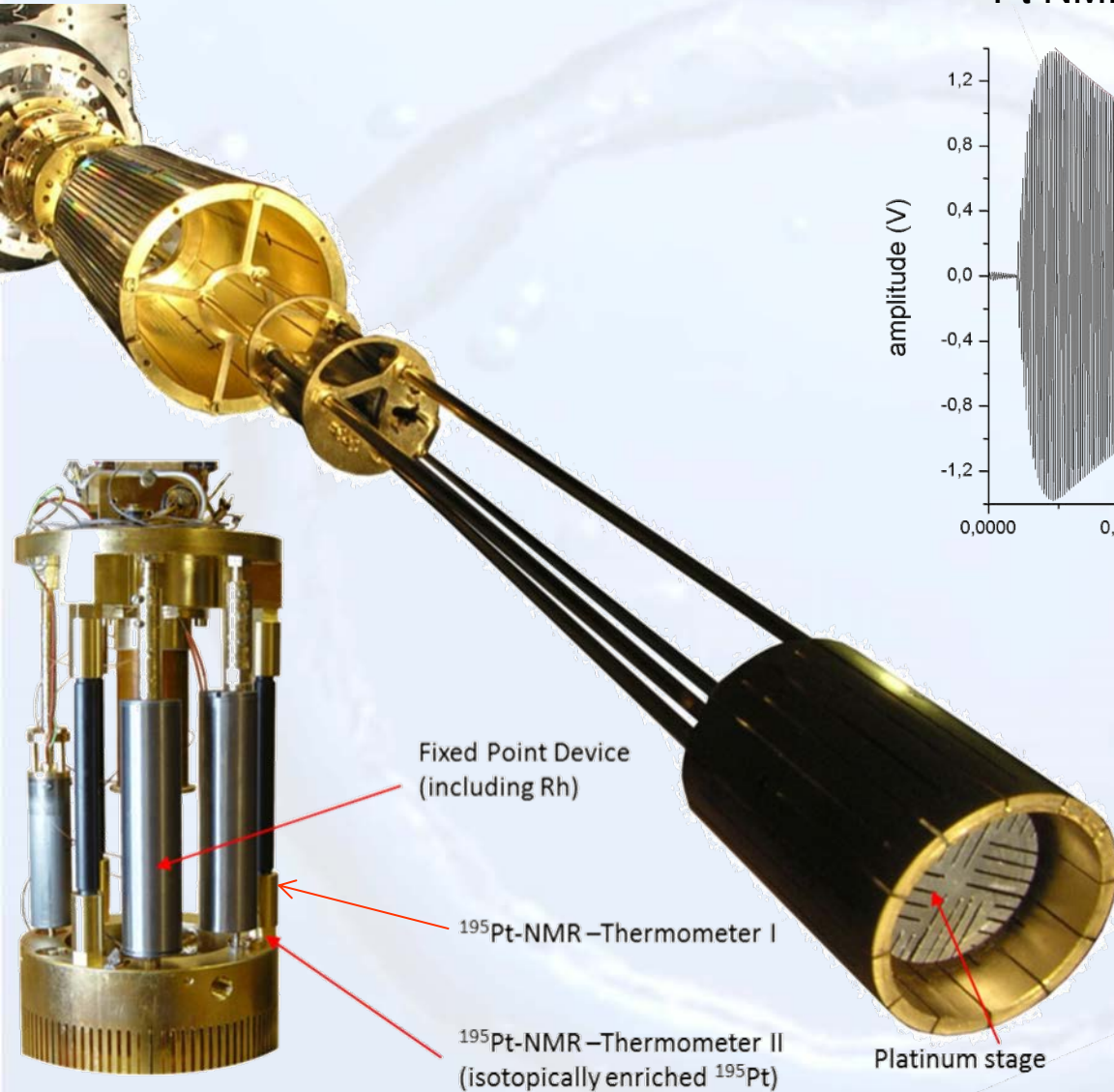


Pt stage

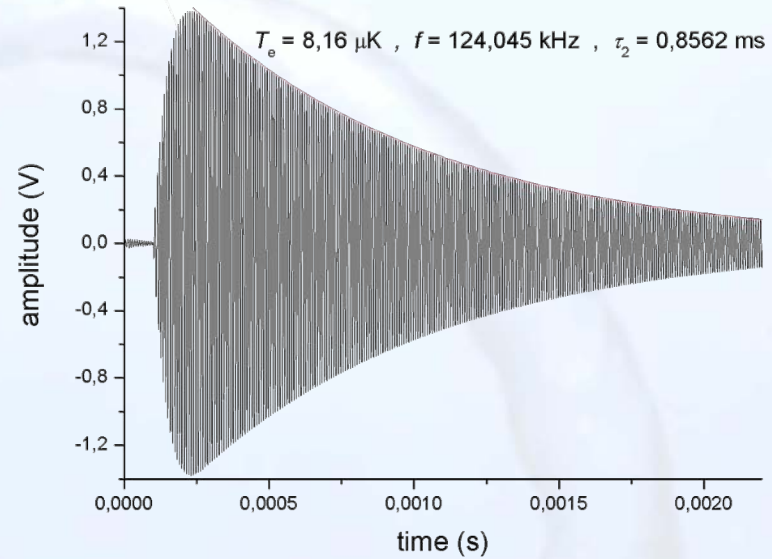




11.3 Adiabatic Demagnetization Refrigerators



¹⁹⁵Pt NMR thermometer signal



Lowest temperature at Pt stage

$$T_{\min} = 800 \text{ nK}$$



Primary thermometers

Superconducting fixpoints

Current/flux noise

^{195}Pt NMR

Coulomb blockade

Nuclear orientation

^3He melting curve

....

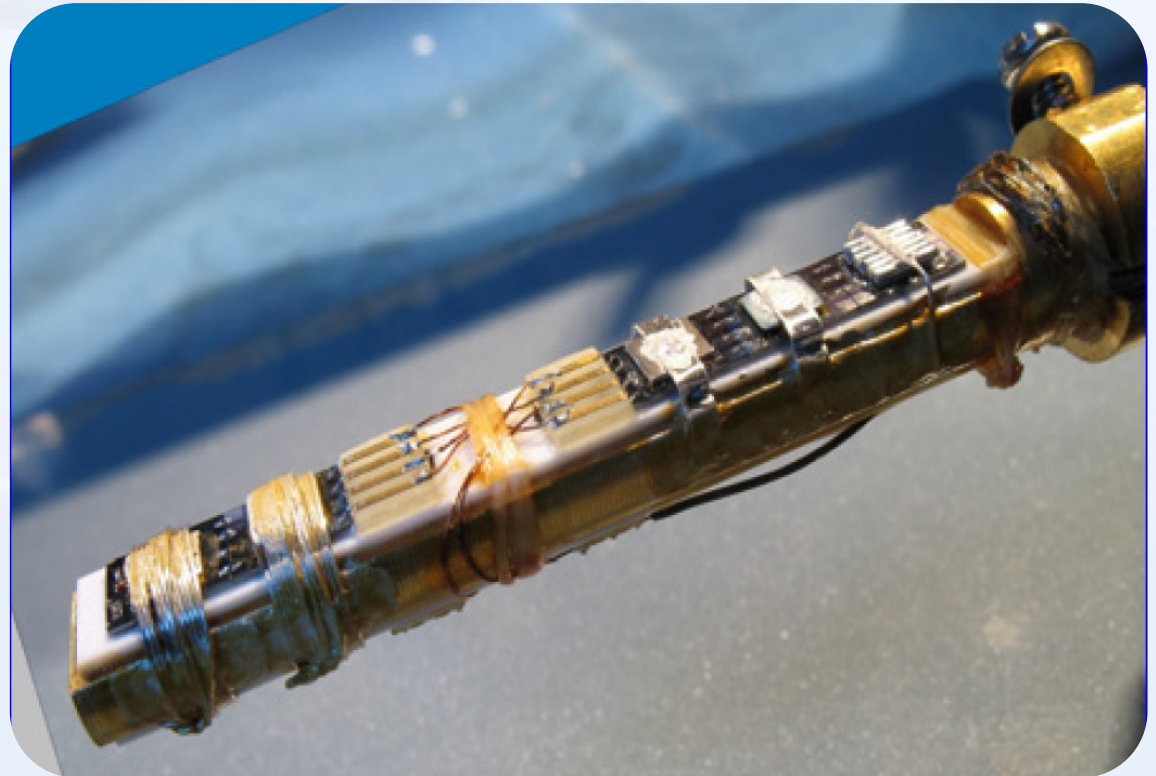
Secondary thermometers

Resistance

Capacitance

Magnetic susceptibility

.....





12. Thermometry at Low Temperature



Temperature is a **thermodynamic property of state**

It can be defined by a **reversible cycle**, like a **carnot cycle** $\oint T^{-1}dQ = 0$

not practical

primary thermometers: can be **used without** any **prior calibration**

secondary thermometers: must be **calibrated against** an **other thermometer**

distinction is often somewhat arbitrary ...

Temperature scales

defined by *Comité International des Poids et Mesures*

based on **fixpoints** like the triple point of water and **interpolation** like Pt-100 resistance thermometry or gas thermometry

ITS-90 0.65 K to 1358 K

PLTS-2000 0.9 mK to 1358 K

