

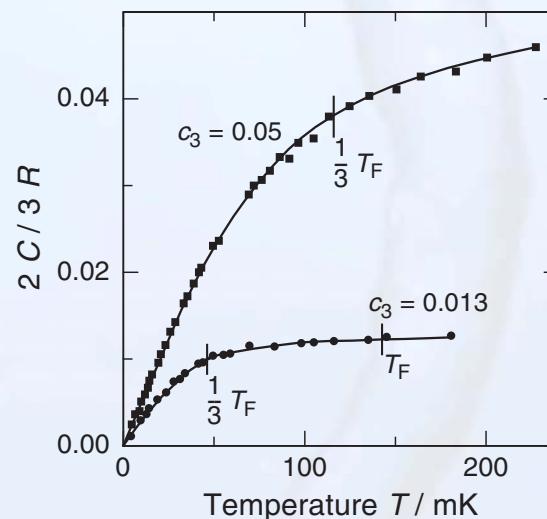
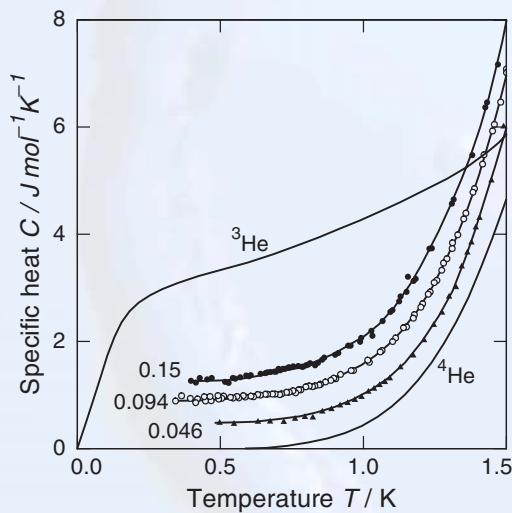


5.1 Specific heat and phase diagram

dilute solutions of ^3He in He-II ($c_3 < 0.15$, $T < 0.5$ K)→ ^4He : passive background fluid→ ^3He : “free” atoms in a **quasi vacuum** and effective mass $m_3^* = 2.4 m_3$

Fermi gas $T_F = \frac{\hbar^2}{2m_3^* k_B} (3\pi^2 n_3)^{2/3} \propto c_3^{2/3}$

$$\begin{aligned} C &\rightarrow T > T_F, \quad C \propto c_3 T^0 \quad (\hat{=} \frac{3}{2} R) && \text{high } T \\ &\rightarrow T < \frac{1}{3} T_F, \quad C \propto T && \text{low } T \end{aligned}$$



- ▶ T_λ depends on c_3
- ▶ pure ^3He : transition Fermi gas → Fermi liquid
- ▶ high T , dilute solution: classical gas with m^*
- ▶ low T : transition classical gas → Fermi gas
- ▶ lines correspond to theory



Finite solubility of ^3He in liquid ^4He at $T = 0$

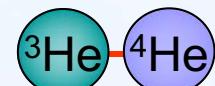
reason: difference in zero-point motion of ${}^3\text{He}$ and ${}^4\text{He}$

v. Waals interaction **identical** for ${}^3\text{He}-{}^3\text{He}$ and ${}^3\text{He}-{}^4\text{He}$

but: larger zero-point motion of ${}^3\text{He}$ weakens the bonding

stronger effective binding force

compared to



in equilibrium on finds

► $T = 0 \longrightarrow c_3 = 1$ for concentrated phase (pure ${}^3\text{He}$)

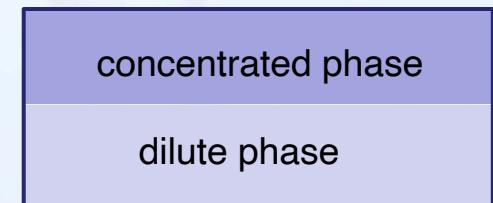
- necessary energy to bring one ${}^3\text{He}$ atom into “vacuum” $L_3(T = 0)$

$$\longrightarrow \mu_{3,c}(0,1) = \mu_3(0) = -L_3(0) = -2.473 \text{ K} \quad \text{latent heat}$$

► dilute phase: $E_3 = -\mu_{3,d}(0, 0)$  binding energy

 $c_{3,d} \rightarrow 0$

- with increasing concentration, the effective binding energy for ${}^3\text{He}$ is reduced because of the Pauli principle \longrightarrow Fermi gas: $E_F = k_B T_F(c_3)$

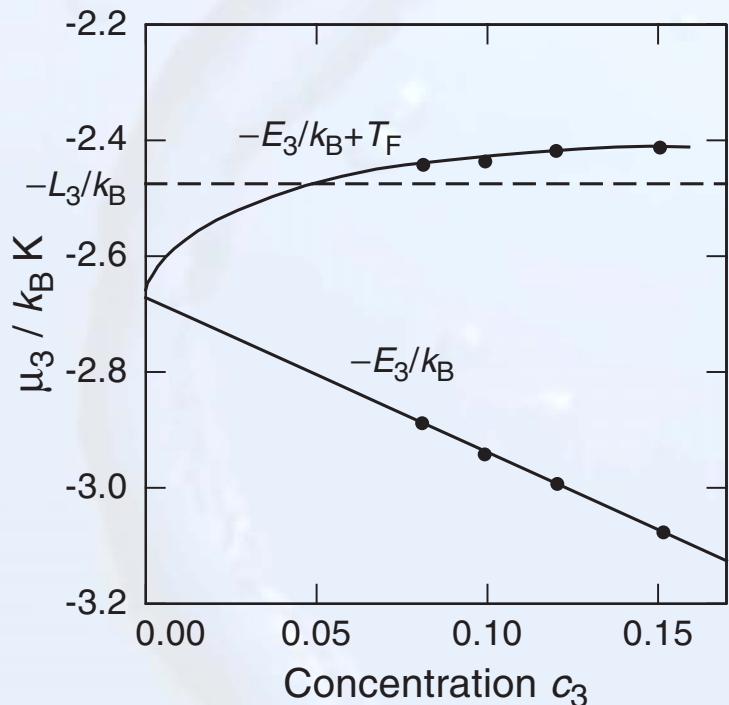


5.1 Specific heat and phase diagram

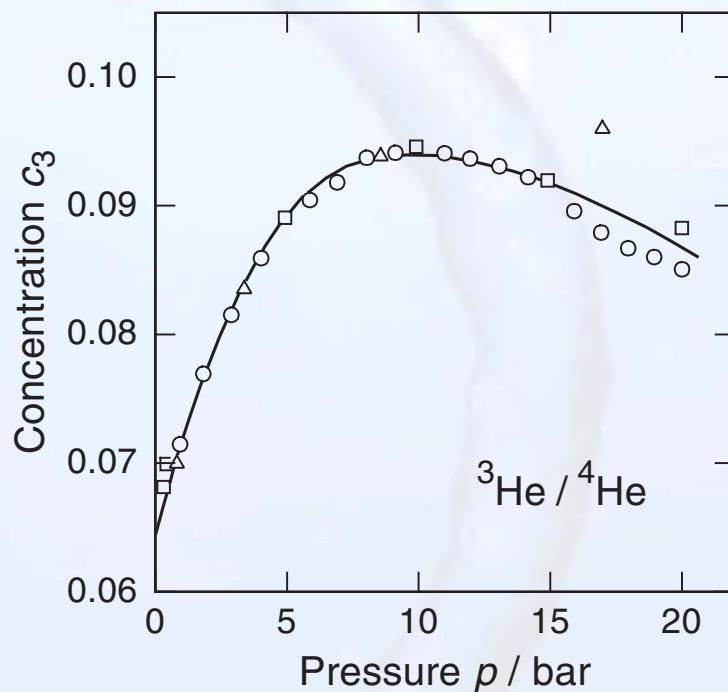


equilibrium concentration at $T = 0$

$$-L_3(0) = -E_3(0, c_3) + k_B T_F(c_3)$$



pressure dependence



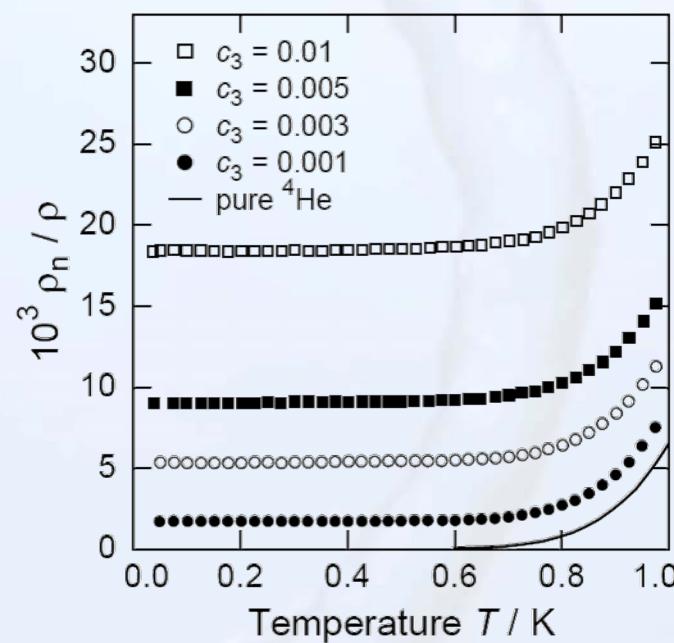
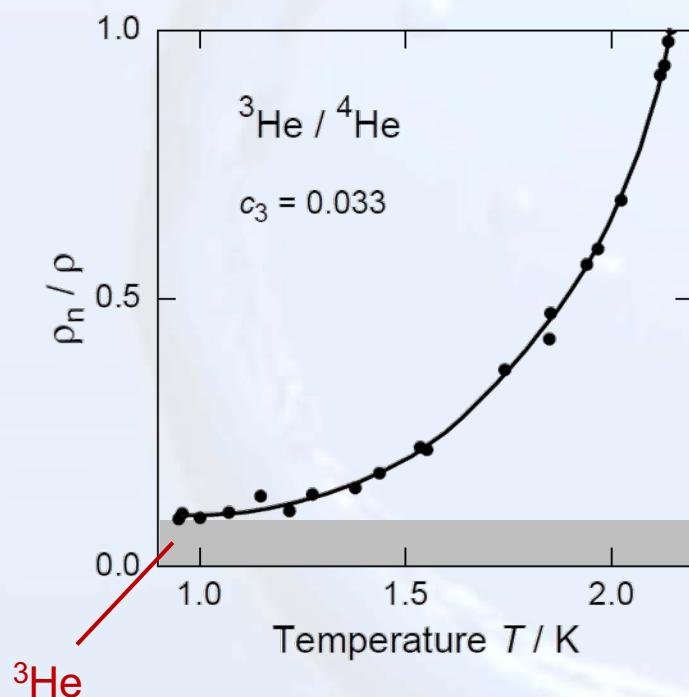
- ▶ calculation of $E_3(0, c_3)$ is not trivial
- Bardeen, Baym, Pines model

- ▶ maximum at 8.7 bar
- ▶ concentration $c_3 = 0.096$



determination of ϱ_n —→ Andronikasvili-type experiment

15 mica sheets
4 cm diameter
190 μm spacing



$$\rightarrow \varrho_n(T \rightarrow 0) = \text{const} \propto c_3$$



Osmotic pressure

- ▶ ${}^4\text{He}$ flows to solution to thin the ${}^3\text{He}$ concentration
- ▶ ${}^3\text{He}$ is blocked
- osmotic pressure

van't Hoff law ($T \gg T_F$, classical regime)

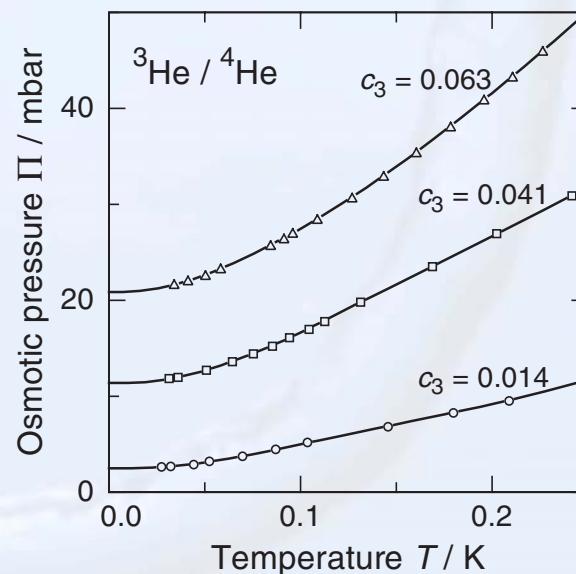
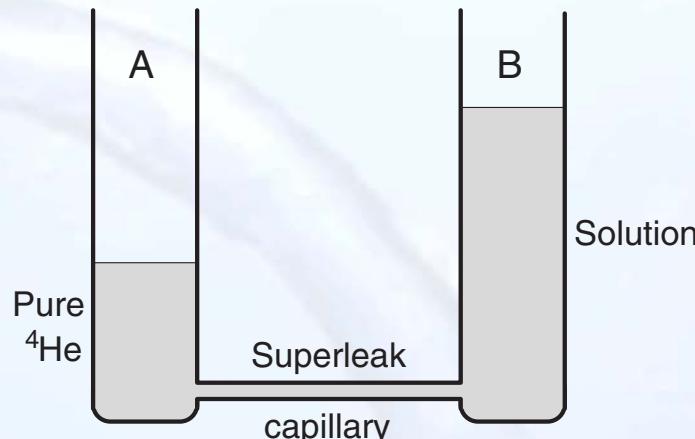
$$\Pi = n_3 k_B T \propto c_3 T$$

$T \ll T_F$, degenerate Fermi gas

$$\Pi = \frac{2}{5} n_3 k_B T_F \propto c_3^{5/3} = \text{const}$$

depends on c_3

→ transition from FG to classical gas

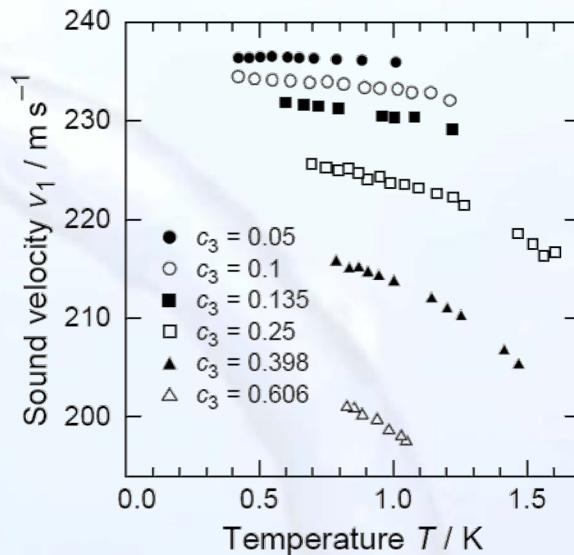




5.3 Sound Propagation

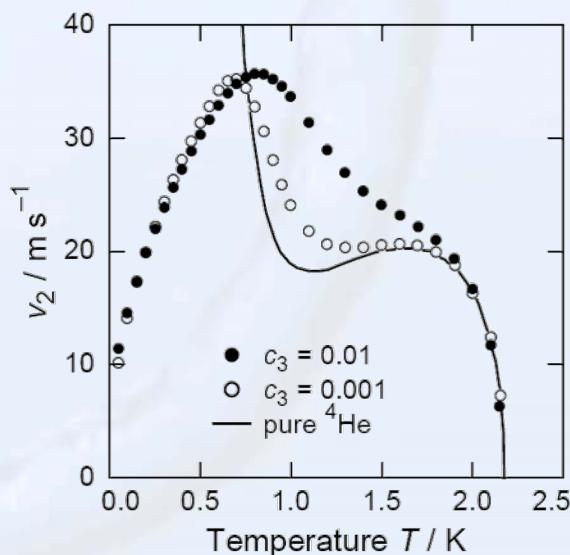
First sound

$$v_1^2 = \left(\frac{\partial p}{\partial \varrho} \right)_{S, c_3} \left[1 + \frac{\varrho_s}{\varrho_n} \left(\frac{\partial \varrho}{\partial c_3} \frac{c_3}{\varrho} \right)^2 \right]$$



Second sound

$$v_2^2 = \frac{\varrho_s}{\varrho_n} \left[\bar{S} \left(\frac{\partial T}{\partial S} \right)_{\varrho, c_3} + c_3^2 \frac{\partial(\mu_3 - \mu_4)}{\partial c_3} \right] \left[1 + \frac{\varrho_s}{\varrho_n} \left(\frac{\partial \varrho}{\partial c_3} \frac{c_3}{\varrho} \right)^2 \right]^{-1}$$
$$\bar{S} = S_{4,0} - \frac{k_B}{m_4} [c_3 + \ln(1 - c_3)] + \frac{k_B}{m_3} c_3$$





5.4 Transport Properties

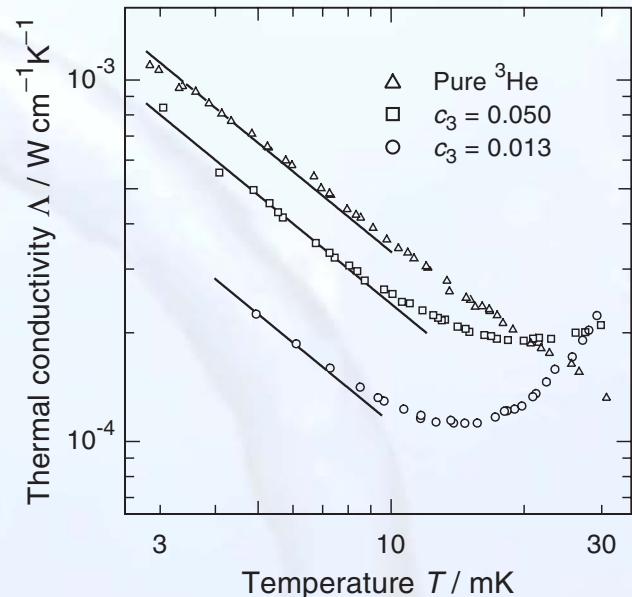
Thermal transport (rather complex)

high Temp.: ϱ_n flow leads to ${}^3\text{He}$ concentration gradient

- ${}^3\text{He}$ atoms diffuse back
- ${}^3\text{He}$ form scattering centers for ϱ_n
- reduction of heat transport

low Temp.: ${}^3\text{He}$ atoms form FG

$$\Lambda = \frac{1}{3} C v_F \ell \propto \frac{c_3}{T}$$
$$\ell = v_F \tau \propto (T_F/T)^2$$
$$v_F = (\hbar/m_3^*)(3\pi^2 n_3)^{1/3},$$
$$C \propto T/T_F$$
$$m_3^* = (1 + F_1/3)m_3$$



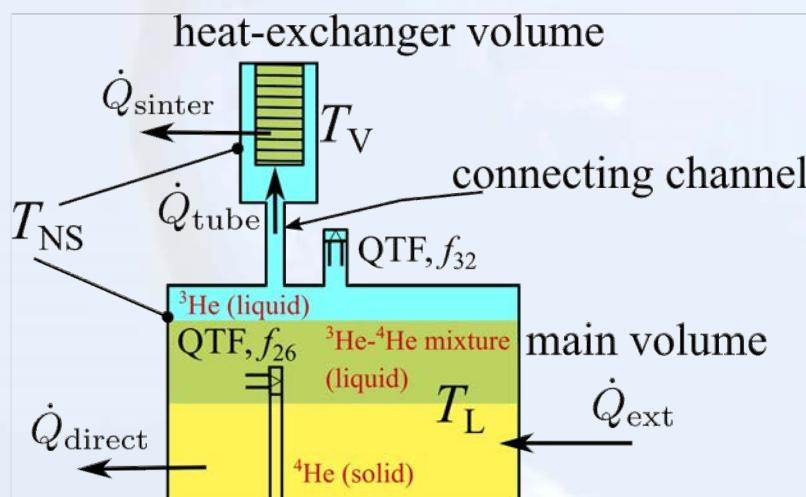


very interesting: 3 superfluid phases in the same container → ${}^4\text{He}$, ${}^3\text{He}$, and dilute ${}^3\text{He}$

Problem: ${}^3\text{He}/{}^4\text{He}$ mixtures are hard to cool to below 200 μK because of Kapitza resistance

new initiative:

- cooling by melting of ${}^4\text{He}$ crystal
- lowest temperature so far 90 μK



↓
acoustic mismatch hinders cooling

