

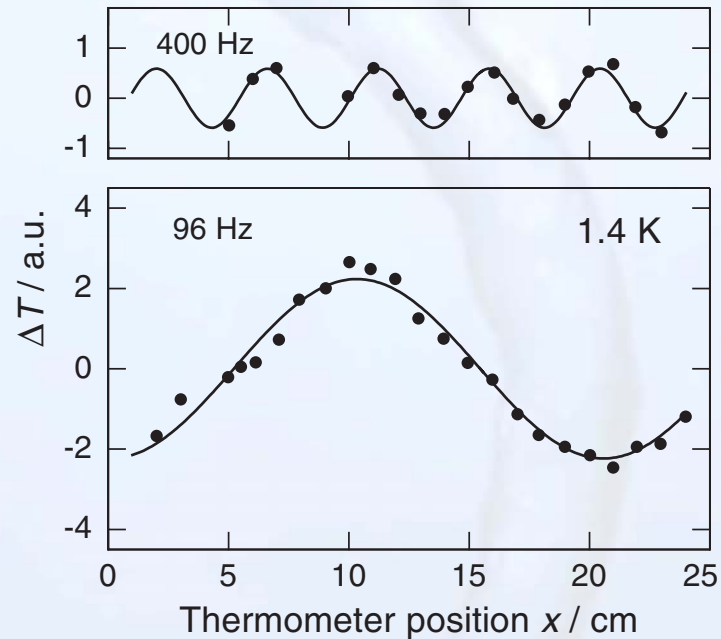
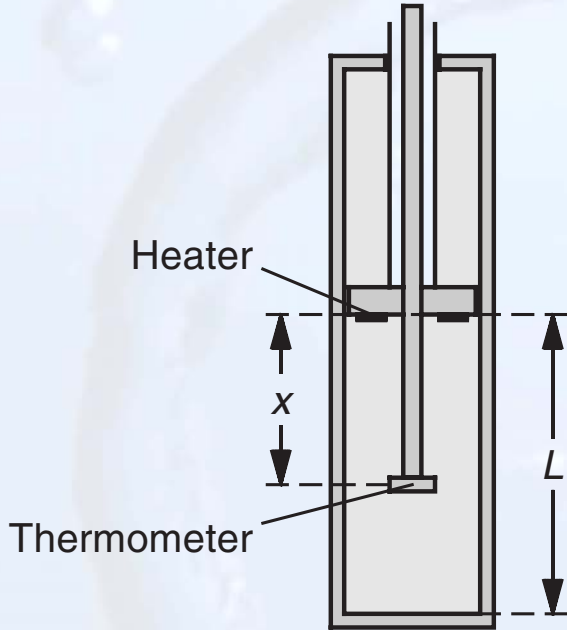


f) Second Sound



Propagation of temperature waves similar to sound waves

suggested by Kapitza
first seen by Peshkov 1944



resonance condition $v_2 = 2L\nu/n$

- ▶ Seen up to 100 kHz (experimental limit)
- ▶ v_2 independent of frequency



2.2 Two-Fluid Model



Basic idea: He-II has two components



normalfluid

superfluid

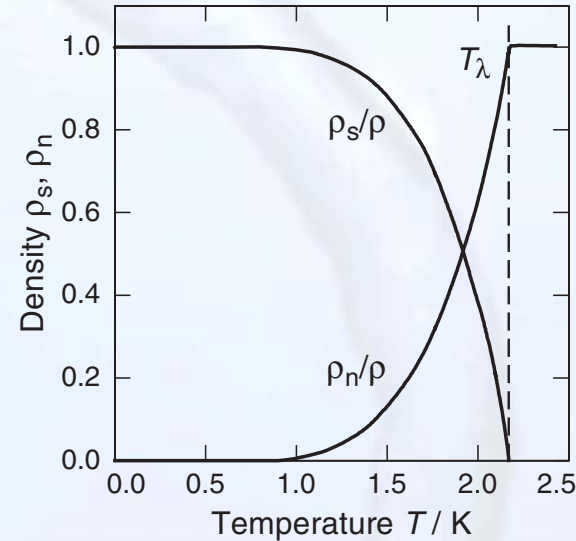
Tisza 1938
London 1938
Landau 1941, 1947
Feynman 1953

Assumptions and Properties:

$$\rho = \rho_n + \rho_s \quad (1)$$

$$T = T_\lambda : \rho_s = 0 \text{ and } \rho_n = \rho$$

$$T = 0 : \rho_s = \rho \text{ and } \rho_n = 0$$



| | density | viscosity | entropy |
|------------------------|----------|-----------------|-----------|
| normal-fluid component | ρ_n | $\eta_n = \eta$ | $S_n = S$ |
| superfluid component | ρ_s | $\eta_s = 0$ | $S_s = 0$ |

In addition: no turbulence associated with $\rho_s \rightarrow \text{rot } \mathbf{v}_s = 0$



density $\rho = \rho_n + \rho_s$ (1)

mass flow $\mathbf{j} = \rho_n \mathbf{v}_n + \rho_s \mathbf{v}_s$ (2)

continuity eqn.
(mass conservation) $\frac{\partial \rho}{\partial t} = -\text{div } \mathbf{j}$ (3)

He-II is **ideal fluid** $\eta_n < 10^{-5} \text{ P} \sim 0$

➔ **Euler eqn.** (Newton's 2nd law of motion for continua)

$$\frac{\partial \mathbf{j}}{\partial t} + \underbrace{\rho \mathbf{v} \text{ div } \mathbf{v}}_{\approx 0} = -\text{grad } p$$

➔

for **small velocities** since quadratic in v
(approximation for **linear regime**)

$$\frac{\partial \mathbf{j}}{\partial t} = -\text{grad } p$$
 (4)



entropy conservation

motion is reversible since **no dissipative** processes \rightarrow **He-II is an ideal fluid**
(in first approximation)

$$\frac{\partial(\overset{\text{entropy / mass}}{\rho S})}{\partial t} = -\text{div}(\underbrace{\rho S \mathbf{v}_n}_{\text{entropy density}}) \quad (5)$$

only ρ_n contributes

One more equation is needed \rightarrow an equation of motion for ρ_s (or ρ_n)

this is difficult to derive \rightarrow see R.B. Dingle, Proc. Phys. Soc. A62, 648 (1949) (40 pages)

here: [Gedankenexperiment](#) according to Landau

idea: **Superfluid component** is **added** at **“constant” volume** in the system



Consider change of internal energy

$$dU = T dS - p dV + G dm$$

$$dS = 0$$

reversible

$$dV = 0$$

$$V = \text{constant}$$

Gibbs free energy per unit mass

$$dU = G dm$$

Gibbs free energy is **potential energy** of **superfluid component/mass**

→ $-\text{grad } G$ is corresponding **force**

$$\frac{d\mathbf{v}_s}{dt} = -\text{grad } \mu \quad \text{and} \quad d\mu = -S dT + \frac{1}{\rho} dp$$

G/m
Chemical potential

$$\rightarrow \frac{\partial \mathbf{v}_s}{\partial t} = S \text{ grad } T - \frac{1}{\rho} \text{ grad } p \quad (6)$$



density $\rho = \rho_n + \rho_s$ (1)

mass flow $\mathbf{j} = \rho_n \mathbf{v}_n + \rho_s \mathbf{v}_s$ (2)

mass conservation
continuity eqn. $\frac{\partial \rho}{\partial t} = -\text{div } \mathbf{j}$ (3)

ideal fluid $\frac{\partial \mathbf{j}}{\partial t} = -\text{grad } p$ (4)

entropy conservation $\frac{\partial(\rho S)}{\partial t} = -\text{div}(\rho S \mathbf{v}_n)$ (5)

an equation of motion for
superfluid component $\frac{\partial \mathbf{v}_s}{\partial t} = S \text{ grad } T - \frac{1}{\rho} \text{ grad } p$ (6)



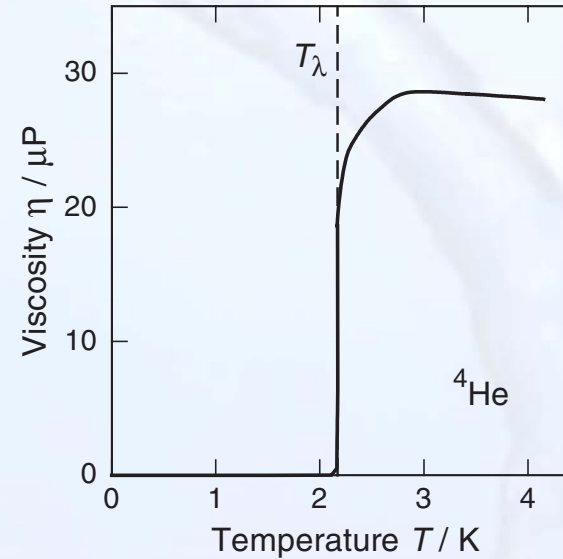
a) Viscosity

(i) capillaries (extremely thin)

Interpretation: $v_n \approx 0$

→ only superfluid phase is observed

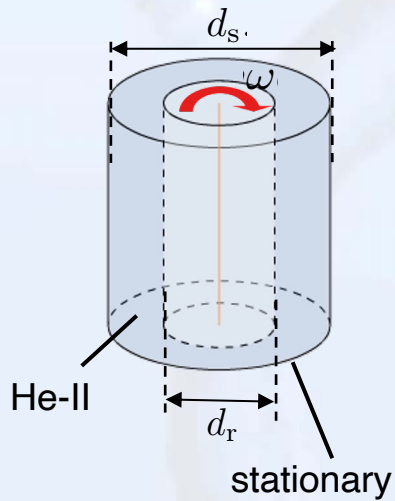
$$\eta = \eta_s = 0$$





a) Viscosity

(ii) rotary viscosimeter



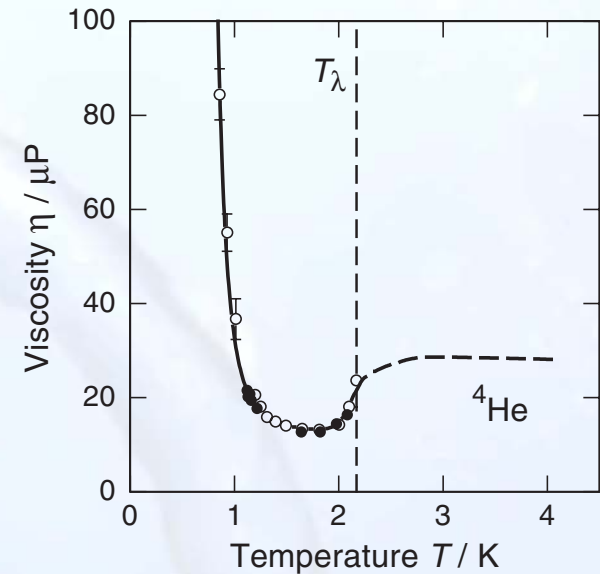
Torque acting on stationary cylinder is measured

$$M_r = \pi \eta \omega d_r^2 d_s^2 / (d_s^2 - d_r^2)$$

since $\eta_s = 0$ no torque resulting from ρ_s

$$\rightarrow M_r \propto \eta = \eta_n$$

↑
two-fluid model



Temperature dependence

$\eta_n(T)$ at very low temperatures $T < 1.8 \text{ K}$?

$\eta_n \propto \ell_n \rightarrow$ mean free path increases with decreasing temperature because thermal excitations disappear

Viscosity

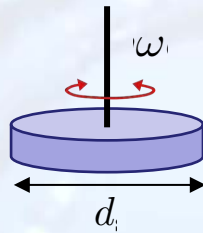
$$\eta = \frac{1}{3} \rho v \ell$$

Landau-Chaladnikov Theory



a) Viscosity

(iii) oscillating disc



Torque acting on the disc:

$$M_d = \pi \sqrt{\rho \eta} \omega^{3/2} r^4 \Theta(\omega)$$

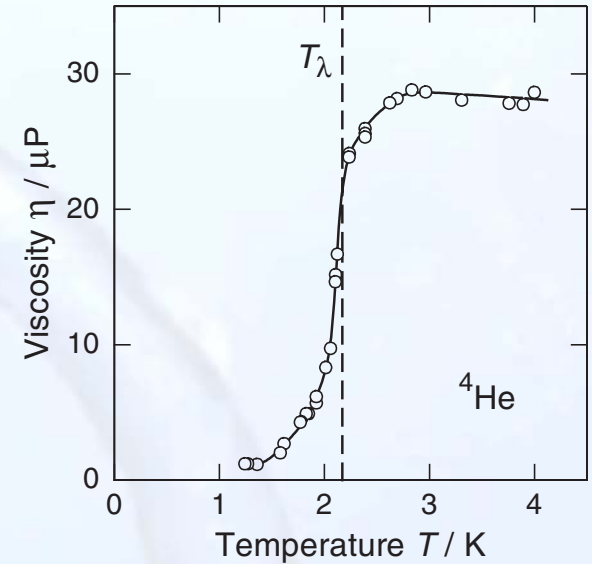
$$\Theta(\omega) = \Theta_0 \cos(\omega t - \pi/4)$$

$$M_d \propto \sqrt{\rho \eta}$$

product is important for M_d

$$T < T_\lambda \quad \longrightarrow \quad \eta_s = 0 \quad \longrightarrow \quad \eta_n \rho_n \text{ is measured}$$

$$\text{for } T \rightarrow 0 \quad \longrightarrow \quad \rho_n \rightarrow 0 \quad \longrightarrow \quad \rho_n \eta_n \rightarrow 0$$





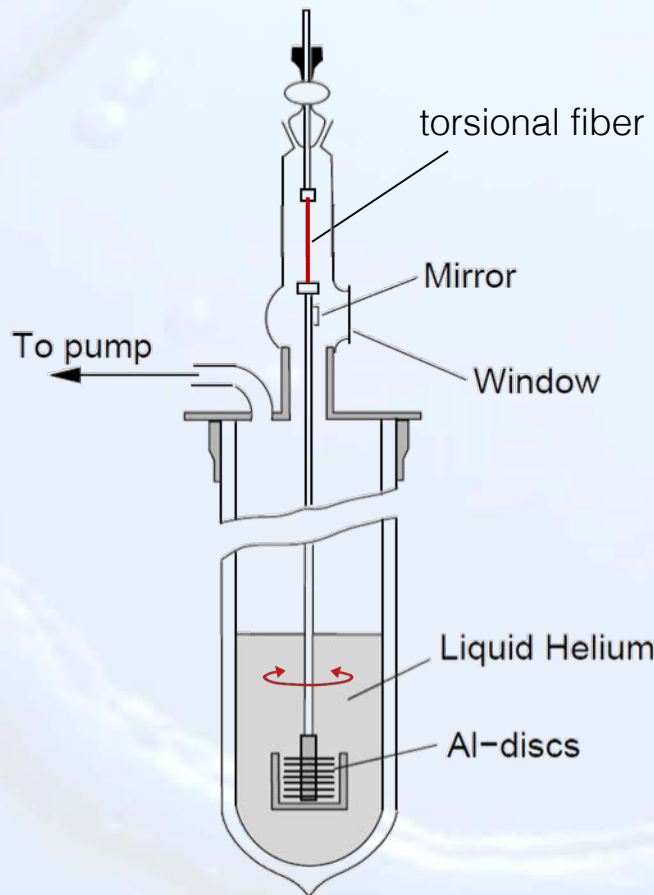
Determination of ρ_n

Experiment of Andronikasvili (1948)

First **direct** observation of ρ_n



Elepter Luarsabovich
Andronikashvili (1910-1989)



50 aluminum discs

thickness 13 μm

diameter 3.5 cm

spacing 210 μm

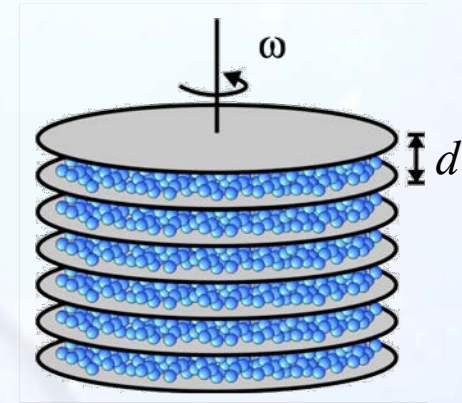


observation \rightarrow slow resonant oscillations (mass and torsion fiber)

Important parameter is the viscos penetration depth for wave with frequency ω

$$\delta = \sqrt{2\eta_n / \rho_n \omega}$$

- $d < \delta$:
- \blacktriangleright ρ_n is dragged along with torsion oscillator above and below T_λ
 - \blacktriangleright ρ_s remains stationary
 - \blacktriangleright period of oscillation determined by mass of torsion oscillator (and spring constant)
- \rightarrow ρ_n can be determined



temperature dependence (empirical relation)

$$\rho_n = \rho_\lambda \left(\frac{T}{T_\lambda} \right)^{5.6}$$

comparison with 2nd Sound \rightarrow fits well

