

High-energy tests of QED: e^+e^-

Complementarity w/r/t low-energy precision measurements

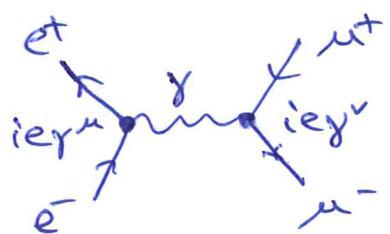
- test at high q^2
- equivalently short distance scales

Major task: radiative effects at higher orders

* theory but also Monte Carlo for realistic detectors and analysis studies \rightarrow needed, to make meaningful experimental validations

- test QED
- disentangle from weak/QCD
- limits on possible substructure of elementary fermions

Example: $e^+e^- \rightarrow \mu^+\mu^-$



$$\frac{d\sigma}{d\Omega} = \frac{1}{64\pi^2 s} \frac{|\vec{p}_f|}{|\vec{p}_i|} \langle |M_{fi}|^2 \rangle$$

} spin sums
(helicity vs. chirality)

leading order: 2 vertices $\Rightarrow |M|^2 \sim e^4 \sim \alpha^2$

next-to-leading order (NLO): 4 vertices $\Rightarrow |M|^2 \sim \alpha^4$

diagrams like , , etc

need to sum amplitudes before squaring to get observables like decay rates or cross sections:

$$M_{fi} = M_{fi}^* + \sum M_{NLO} + \dots$$

$$|M_{fi}|^2 = |M_{fi}^*|^2 + \sum (M_{fi}^* M_{NLO}^* + c.c.) + |\sum M_{NLO}|^2 + \dots$$

$\sim \alpha^2$

$\sim \alpha^3$

$\sim \alpha^4$ etc

Each term in $|M_{fi}|^2$ is suppressed by $\alpha \sim \frac{1}{137}$ w/r/t the previous one. (But there are more diagrams!)
 \Rightarrow naively, LO should be good to $\sim 1\%$ for QED

Spin sums and helicity vs. chirality

LO diagram for $e^+e^- \rightarrow \mu^+\mu^-$: 2 possible spin states

\Rightarrow 16 orthogonal helicity combinations

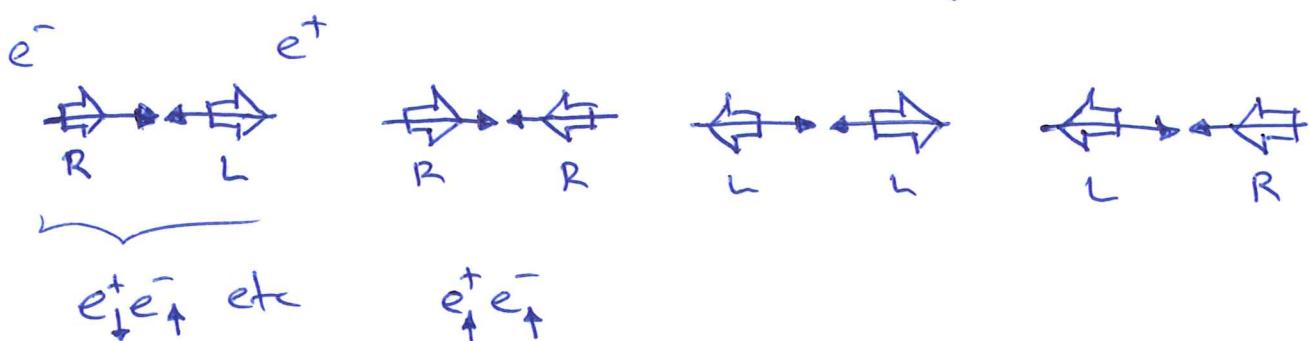
Need to calculate 16 matrix elements?

- trace techniques \rightarrow see muon decay also
- direct calculation for each helicity (more laborious)

Total annihilation rate for one e^+e^- helicity state

$$= \sum_{\mu^+\mu^- \text{ states}} (\text{rate for each } \mu^+\mu^- \text{ state})$$

Unpolarized beams \Rightarrow equal amounts of all helicities are initially present for e^+e^-

$$\Rightarrow \langle \rangle \sim \frac{1}{4} \sum_{\text{spins}}$$


$$\langle |M_{fi}|^2 \rangle = \underbrace{\frac{1}{(2s_e + 1)^2} \sum_{e^+e^-} \sum_{\mu^+\mu^-}}_{\text{avg. over initial spins}} \sum_{\mu^+\mu^-} |M_{fi}|^2$$

sum over final spins

helicity: $\frac{\vec{e} \cdot \vec{p}}{2p} \sim \frac{\vec{S} \cdot \vec{p}}{p}$

- sph projection along the momentum direction
- not Lorentz-invariant, $m > 0$
- but is a constant of motion

chirality: $P_R = \frac{1}{2}(1 + \gamma^5)$
 $P_L = \frac{1}{2}(1 - \gamma^5)$

- transformation in LH/RH representation of Poincaré
- is Lorentz-invariant
- not a constant of the motion

chiral structure of QED:

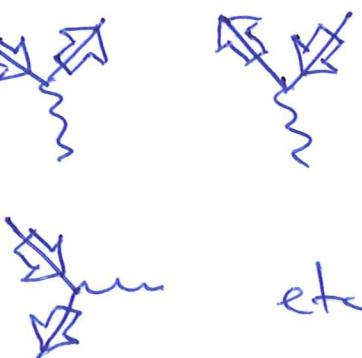
$E \gg m$, only 4 of the 16 helicity combinations give $u_f \neq 0$

helicity and chirality eigenstates are the same for massless particles
 { ultrarelativistic limit

helicity "conserved" (so can also speak of LH/RH helicity) for high-energy QED interactions

e.g. $\bar{u}_f \gamma^\mu u_f = 0$

nonzero:



Returning to $e^+e^- \rightarrow e^+e^-$ as a concrete example:

$$M_{fi} = -\frac{e^2}{q^2} \underbrace{(\bar{e}^+ \gamma^\mu e^-)}_{\text{propagator } ie} \underbrace{(\bar{e}^+ \gamma^\mu e^-)}_{\text{propagator } ie} \quad \text{Junk}$$

coupling

The helicity spinors can be used to work out the various possibilities explicitly, e.g.

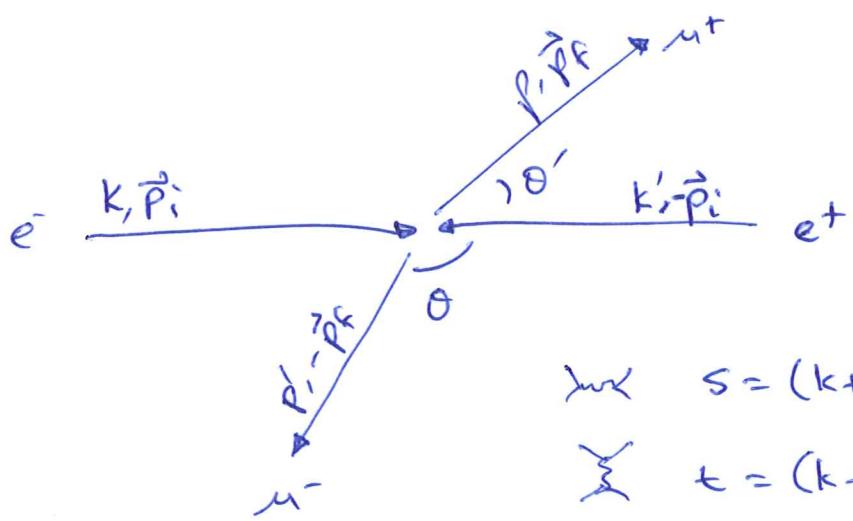
$$N_{RL \rightarrow RL} = e^2(1 + \cos\theta) = 4\pi\alpha(1 + \cos\theta)$$

$$N_{RL \rightarrow LR} = e^2(1 - \cos\theta)$$

$$\langle |M|^2 \rangle = e^4 \left[(1 + \cos\theta)^2 + (1 - \cos\theta)^2 \right] = e^4 (1 + \cos^2\theta)$$

$$\Rightarrow \frac{d\sigma}{ds} = \underbrace{\frac{e^4}{64\pi^2 s}}_{\alpha^2/4s} (1 + \cos^2\theta) = \frac{e^4}{32\pi^2 s} \frac{t^2 + u^2}{s^2}$$

Integrating, $\sigma_{\text{tot}} = \frac{4\pi\alpha^2}{3s} \approx \frac{87 \text{ nb}}{s} \text{ GeV}^2$



When is $\theta' \neq \pi - \theta$?

$$\text{Ansatz } s = (k+k')^2 = 4E_i^2$$

$$\begin{aligned} t &= (k-p)^2 = -2E_i^2(1+\cos\theta) \\ &= -\frac{s}{2}(1+\cos\theta) \end{aligned}$$

$$\begin{aligned} u &= (k-p')^2 = -2E_i^2(1-\cos\theta) \\ &= -\frac{s}{2}(1-\cos\theta) \end{aligned}$$

When is a u-channel diagram relevant?

Tests of cut-off scales in QED

- are fundamental fermions point-like?
- could there be a heavy photon w/ modified propagator?

Let $\Delta \sim$ mass of heavy photon w/ coupling strength still α

⇒ modified propagator and vertices

$$\frac{1}{q^2} \rightarrow \frac{1}{q^2} - \frac{1}{q^2 - \Delta^2} = \frac{1}{q^2} \left(1 - \frac{\Delta^2}{q^2 - \Delta^2} \right)$$

⇒ change of EM potential: $\frac{1}{r} \rightarrow \frac{1}{r} (1 - e^{-\Delta r})$ Yukawa
modifies point-like nature of the em interaction

⇒ modified $e^+e^- \rightarrow \gamma\gamma$ form factor:

$$\frac{d\sigma}{ds} = \frac{\alpha^2}{s} \frac{1 + \cos^2 \theta}{\sin^2 \theta} \left(1 \pm \frac{s^2}{2\Delta_\pm^4} \frac{\sin^4 \theta}{1 + \cos^2 \theta} \right)$$

- lower sign (Δ_-) allows a possibly lower cross-section
- note $\frac{s^2}{\Delta^4}$ scaling and form at $\theta = \frac{\pi}{2}$ (maximal sensitivity at 90°)
- QED allows coupling via magnetic moments - need to conserve the electromagnetic current

if switched to "-"
then this is a heavy (excited?) electron w/ coupling via the same charge e

↓
set limits on its mass via Δ

Sensitivity typically rises with COM energy
⇒ limited by luminosity measurement

Second example: $e^+e^- \rightarrow e^+e^-$ (Bhabha)

total cross-section: $\sigma = \frac{4\pi\alpha^2}{3s} \left(1 \mp \frac{s}{s-\Delta_{\pm}} \right)^2$

$$\frac{d\sigma}{dt} = \frac{\alpha^2}{2s} \left[\left(\frac{s}{t} \right)^2 F(t)^2 + \left(\frac{t}{s} \right)^2 F(s)^2 + \left(1 + \frac{s}{t} \right)^2 \left(1 + \frac{t}{s} \right)^2 \left(\frac{s}{t} F(t) + F(s) \right)^2 \right]$$

with $F(q^2) = 1 \mp \frac{q^2}{q^2 - \Delta_{\pm}^2}$

timelike: $q^2 = s$

spacelike: $q^2 = t$
 $= -\frac{s}{2}(1 + \cos\theta)$

(Note different scaling w/r/t $e^+e^- \rightarrow \gamma\gamma$)

Bhabha: $\sim \frac{q^2}{\Delta^2}$