# MVCMP-1 Bardeen – Josephson Debate

## The Nobel Laureate Versus the Graduate Student

In a recent note, Josephson uses a somewhat similar formulation to discuss the possibility of superfluid flow across the tunneling region, in which no quasi-particles are created. However, as pointed out by the author [Bardeen, in a previous publication], pairing does not extend into the barrier, so that there can be no such superfluid flow.



Physics Today 54, 46-51 (2001)



## **10.4 Unconventional Superconductors**

LaH<sub>10</sub> @ 188 GPa <sup>O</sup>.





### Quantum states

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		S	L
<i>s</i> -wave	conventional superconductors	0	0
<i>p</i> -wave	(Sr <sub>2</sub> RuO <sub>4</sub> )	1	1
<i>d</i> -wave	cuprate high- $T_{\rm c}$ superconductors	0	2
<i>f</i> -wave	(UPt <sub>3</sub> )	1	3

energy gap $\Delta_{m{k}} = \Delta_0(T) f(\widehat{m{k}})$		nodes
conventional superconductors	isotropic	
cuprate high-T <sub>c</sub> superconductors	2D	
Sr <sub>2</sub> RuO <sub>4</sub>	axial 2D	
UPt <sub>3</sub>	3D	Ó
UBe <sub>13</sub>	axial 3D	



# **10.4 Unconventional Superconductors**

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#### cuprate high- $T_c$ superconductors

discovered June 1986

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#### Possible High $T_c$ Superconductivity in the Ba – La – Cu – O System

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Georg Bednorz

Karl Alexander Müller





APS March Meeting of 1987 The "Woodstock of Physics" Hilton Hotel New York

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## best investigated system: YBCO Yttrium barium copper oxide

important are the Cu-O planes — different phases depending on hole concentration introduced by oxygen surplus



structure

•	
Ва	
Cu	

$0 < \delta < 0.4$	insulator
$\delta \approx 0.4$	insulator-metal transition
$\delta > 0.4$	superconductor
$\delta = 0.92$	superconductor with $T_{\rm c}$ = 95 K

AF anti-ferromagnetic phase
PG pseudo gap
CDW charge density wave
SM strange metal
FL Fermi liquid phase



# **10.4 Unconventional Superconductors**

## superconducting transition

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a) resistance

- resistivity depends on crystal direction
- in *c*-direction (perpendicular to CuO<sub>2</sub> planes) resistivity is much higher

b) specific heat

- second order phase transition
- rounded onset of transition caused by large fluctuation
  - indicates short coherence length



# **10.4 Unconventional Superconductors**

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## Critical fields

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$$B_{
m c1}=rac{\Phi_0}{4\pi\lambda_{
m L}^2}$$

- $\blacktriangleright$   $B_{c1}$  depends on crystal direction
- $B_{c1}$  is very small at T = 0
  - $\rightarrow$   $\lambda_{L} = 150 \text{ nm} \longrightarrow \text{ factor 10 larger as for Al}$
  - $ightarrow \lambda_{
    m L}^2 \propto 1/n_{
    m s} ~~$  factor 100 less Cooper pairs

$$B_{\rm c2} = \frac{\Phi_0}{2\pi\xi_{\rm GL}^2}$$

- $\triangleright$  B<sub>c2</sub> depends on crystal direction
- ►  $B_{c2}$  is very large at  $T = 0 \longrightarrow B_{c2} > 800$  T for  $B \parallel ab$ 
  - $\rightarrow$  very short coherence length  $\xi_{GL} = 1.5$  nm
- $\longrightarrow$   $\lambda_{\rm L} \gg \xi_{
  m GL}$   $\longrightarrow$  extreme type II superconductor







Experimental determination of *d*-wave nature of Cooper pairs inYBCO

flux quantization Josephson effect

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a) flux quantization

normal geometry -- YBCO ring

- flux quantization measured with SQUID
- result:  $\Phi_0 = h/2e$





## Experimental determination of *d*-wave nature of Cooper pairs in YBCO

unconventional flux quantization

## tri-crystalline rings

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- ► YBCO film is grown epitaxially on SrTiO<sub>3</sub> substrate
- ► SrTiO<sub>3</sub> is suitably tailored having 3 crystal orientations
- YBCO is patterned by ion milling to produces rings
- rings are positioned that they contain grain boundaries
- macroscopic quantum states in different parts of the rings have different orientations
- the grain boundaries produce  $\pi$  junctions

**0** or even number of  $\pi$  junctions  $\Phi = n \Phi_0$ 

odd number of  $\pi$  junctions

 $\Phi = \left( n + \frac{1}{2} \right) \Phi_0 \quad \longrightarrow \quad \text{half-flux quantum}$ 

## Experimental result

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## scanning SQUID microscopy





- rings with even number of  $\pi$  junctions show no flux
- ring in the middle with 3  $\pi$  junctions shows spontaneous formation of half-flux quantum



**11. Cooling Techniques** 



## <sup>4</sup>He bath cryostat: glass dewar











## <sup>4</sup>He Bath cryostat: metal dewar







## helium transport vessel

#### helium transfer tube







## Radiation shields – super insulation



multiple radiation shields  $\rightarrow$  smaller steps  $\rightarrow$  reduction of heat flow

30 to 80 layers of low conductivity high reflection material  $\rightarrow$  aluminized Mylar

apparent thermal conductivity  $\sim 10^{-4}$  to  $10^{-5}$  W/(m K)





#### Cryostats with 1-K-Pot

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<sup>4</sup>He  $L = 90 \text{ J mol}^{-1}$ <sup>3</sup>He  $L = 40 \text{ J mol}^{-1}$ 

#### Vapor pressure curve of various cryogenic liquids

**Clausius-Clapeyron equation** 



vapor pressure curve



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<sup>3</sup>He cryostats



cooling power  $\dot{Q} = \dot{n}_{
m g} \, L \propto p \propto {
m e}^{-L/RT}$ 





Cooling power of a <sup>3</sup>He cryostat with charcoal absorption pump







#### History

- 1951 basic idea suggested by Heinz London
- 1962 detailed concept worked out by London, Clark, Mendoza
- 1965 first realization Das, De Bruyn Ouboter, Taconis  $T_{min} = 220 \text{ mK}$
- 1999 lowest temperature obtained , J.C. Cousins *et al.*  $T_{min} = 1.75$  mK



Heinz London











occurrence of miscibility gap

but 6.5 % <sup>3</sup>He in <sup>4</sup>He at T = 0 K

reason:

zero-point motion weakens binding



but: Fermi energy

max. 6.5% <sup>3</sup>He in <sup>4</sup>He at T = 0 K



principal of cooling by mixing <sup>3</sup>He/<sup>4</sup>He

- transition of <sup>3</sup>He into the <sup>4</sup>He rich phase
- cooling by "evaporation" of <sup>3</sup>He into <sup>4</sup>He quasi vacuum



heat of solubility per Mol:

$$\Delta Q = T\Delta S = aT^2$$

$$a = -84 \,\mathrm{J/K^2}$$





#### Realisation of <sup>3</sup>He/<sup>4</sup>He cooling cycle

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# **11.2 Dilution Refrigerators**



#### Kapitza Resistance – thermal boundary resistance

Snell's law of refraction

$$\frac{\sin \alpha_\ell}{\sin \alpha_{\rm s}} = \frac{v_\ell}{v_{\rm s}}$$

critical angle of total reflection  $\alpha_{\ell}^{c} = \arcsin\left(\frac{v_{\ell}}{v_{s}}\right)$ 

for liquid helium and copper  $~~ lpha_\ell^{
m c} pprox 4^\circ$ 

fraction of phonons incident within critical angle

$$f = \frac{1}{2}\sin^2 \alpha_{\ell}^{c} = \frac{1}{2} \left(\frac{v_{\ell}}{v_{s}}\right)^2 < 10^{-2}$$

transmission coefficient

 $t = \frac{4Z_{\ell}Z_{\rm s}}{\left(Z_{\ell} + Z_{\rm s}\right)^2} \approx \frac{4Z_{\ell}}{Z_{\rm s}} = \frac{4\varrho_{\ell}v_{\ell}}{\varrho_{\rm s}v_{\rm s}}$  $\overbrace{Z_{\ell} = \varrho_{\ell}v_{\ell}} \sum Z_{\rm s} = \varrho_{\rm s}v_{\rm s} \quad \text{acoustic impedances}$ 

fraction of phonons crossing the interface

$$ft=rac{2arrho_\ell v_\ell^3}{arrho_{
m s} v_{
m s}^3}$$



- Kaptiza resistance occurs at any solidsolid, liquid-solid interface
- particular problematic for liquid helium because of the low sound velocity
- helium-copper  $ft < 10^{-5}$



# **11.2 Dilution Refrigerators**





silver sinter SEM image

Kapitza resistance between pure <sup>3</sup>He and <sup>3</sup>He/<sup>4</sup>He mixtures and silver sinters of different grain sizes



- ▶  $20 \,\mathrm{mK} < T < 100 \,\mathrm{mK}$  good agreement with Debye model  $R_{\mathrm{K}} \propto T^{-3}$
- $\blacktriangleright$  below 20 mK  $R_{
  m K} \propto T^{-2}$  or  $R_{
  m K} \propto T^{-1}$ 
  - → anomalous Kapitza resistance
  - origin: TLS, coupling to zero and second sound modes, phonon wavelength larger than sinter grains

heat flow from liquid to solid (using Debye model)

$$\dot{\mathcal{Q}} = \frac{1}{2} ftuv_{\ell} A = \frac{\pi^2 k_{\rm B}^4 \varrho_{\ell} v_{\ell}}{30\hbar^3 \varrho_{\rm s} v_{\rm s}^3} AT^4$$
$$\bigvee_{u = U/V = \pi^2 k_{\rm B}^4 T^4 / (30\hbar^3 v_{\ell}^3)$$

in equilibrium identical heat flow from solid to liquid

net flow in non-equilibrium ( $\Delta T$ )

$$\dot{Q} = \frac{\mathrm{d}\dot{\mathcal{Q}}}{\mathrm{d}T} \Delta T = \frac{2\pi^2 k_{\mathrm{B}}^4 \varrho_\ell v_\ell}{15\hbar^3 \varrho_{\mathrm{s}} v_{\mathrm{s}}^3} A T^3 \Delta T$$

#### Kapitza resistance

$$R_{
m K} = rac{A\Delta T}{\dot{Q}} = rac{15\hbar^{3}arrho_{
m s} v_{
m s}^{3}}{2\pi^{2}k_{
m B}{}^{4}arrho_{\ell} v_{\ell}} \, rac{1}{T^{3}}$$





## Cooling power

assuming 100% <sup>3</sup>He circulation one finds in equilibrium:

$$\dot{Q}_{\rm mc} + \dot{N}_3 \left[ H_3(T_{\rm ex}) - H_3(T_{\rm mc}) \right] = \dot{N}_3 \left[ H_{3,\rm d}(T_{\rm mc}) - H_3(T_{\rm mc}) \right]$$

enthalpy

$$H = U + pV$$

circulation rate

heat leak and/or available cooling power

enthalpy of <sup>3</sup>He-dilute phase

enthalpy of <sup>3</sup>He-rich phase

mixing chamber temperature

temperature after last heat exchanger

inserting the enthalpies

$$\dot{Q}_{\rm mc} = \dot{N}_3 \left[ H_{3,\rm d}(T_{\rm mc}) - H_3(T_{\rm ex}) \right]$$
$$= \dot{N}_3 \left( 95 \, T_{\rm mc}^2 - 11 \, T_{\rm ex}^2 \right) \, \left( \frac{\rm J}{\rm mol \ K^2} \right)$$





## Temperature and circulation rate dependence of the cooling power



limiting case of vanishing cooling power:  $\dot{Q}_{
m mc}=0$ 

$$95 T_{\rm mc}^2 - 11 T_{\rm ex}^2 = 0$$

$$\frac{T_{\rm ex}}{T_{\rm mc}} = 2.8$$

 this underlines the importance of the heat exchanger quality

 $\blacktriangleright$  for  $\dot{Q}\gg\dot{Q}_{
m heat\ leak}$   $\longrightarrow$   $\dot{Q}\propto T^2$  ,  $\dot{Q}\propto\dot{N}_3$ 

heat leak determines lowest temperature

circulation rate