



quantization of circulation

Josephson effects

wave function of superfluid component

$$\psi(\boldsymbol{r}) = \psi_0 e^{\mathrm{i}\varphi(\boldsymbol{r})} \qquad \text{(*)} \qquad \text{with} \quad \psi^*\psi = |\psi_0|^2 = \frac{\varrho_\mathrm{s}}{m_4}$$

Schrödinger equation

$$-\mathrm{i}\hbar\nabla\psi=\boldsymbol{p}\,\psi$$
 with $(*)$ $\boldsymbol{p}=\hbar\nabla\varphi(\boldsymbol{r})=m_4\boldsymbol{v}_\mathrm{s}$ $\boldsymbol{v}_\mathrm{s}=\frac{\hbar}{m_4}\nabla\varphi(\boldsymbol{r})$



$$oldsymbol{v}_{\mathrm{s}} = rac{\hbar}{m_4}
abla arphi(oldsymbol{r})$$

comment:

only valid at sufficiently low velocity were Q_{S} is constant

 $v_{\rm S}$ determines the phase shift of wave function

$$m{v_{
m s}} = 0 \longrightarrow {
m phase is constant}$$
 $m{v_{
m s}} = {
m const.} \longrightarrow {
m phase is changes uniformly}$

Interpretation

- phase is well-defined in entire liquid
- macroscopic wave function
- "rigid" coupling in momentum space



Proof of the concept: He-II under rotations

measurement of liquid meniscus

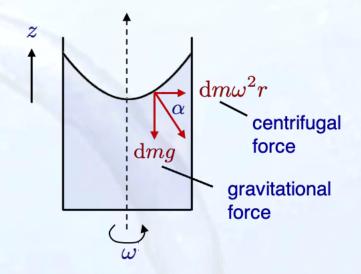
classical fluid \triangleq normalfluid component Q_n

 \longrightarrow solid body rotation $v_{\mathrm{n}} = \omega r$

distance from axis of rotation

profile of liquid surface parabola

$$\tan \alpha = \frac{\mathrm{d}z}{\mathrm{d}r} = \frac{\omega^2 r}{g} \longrightarrow z = \frac{\omega^2}{2g} r^2$$



what about the superfluid component?

two-fluid model $\operatorname{curl} oldsymbol{v}_{\mathrm{s}} = 0$

for a simply-connected region this means every loop can be contracted to a point

$$\int_{A} \underbrace{\operatorname{curl} \boldsymbol{v}_{s}}_{=0} \cdot d\boldsymbol{f} = \oint_{L} \boldsymbol{v}_{s} \cdot d\boldsymbol{l} = 0$$

Stokes

area enclosed by contour ${\cal L}$

$$\triangleright$$
 $\varrho_{\rm S}$ should not rotate (should be at rest)

if so, centrifugal force is reduced



Experimental results

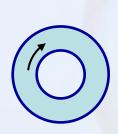
surface curvature: $\gamma = \omega^2/g$ all liquid

$$\gamma = (arrho_{
m n}/arrho)\omega^2/g$$
 only normalfluid

curvature for all liquid is observed in Osborn experiment

Why is this the case?

let's do a thought experiment with an annular-shaped container



circulation:

$$\kappa = \oint_{L} \mathbf{v}_{s} \cdot d\mathbf{l} \qquad \qquad \kappa = \frac{\hbar}{m_{4}} \Delta \varphi_{L}$$
$$\mathbf{v}_{s} = \frac{\hbar}{m_{4}} \nabla \varphi(\mathbf{r})$$

multiply-connected region

- since $\psi({m r})$ is a uniquely-defined function phase can only be changed by $2\pi n$ for full cycle
- $\Delta \varphi = 2\pi n \qquad n = 0, 1, 2, 3, \dots$

He II Surface curvature γ / cm $^{-1}$ T = 1.1 K0.1 Angular velocity v / rad s⁻¹

circulation is quantized

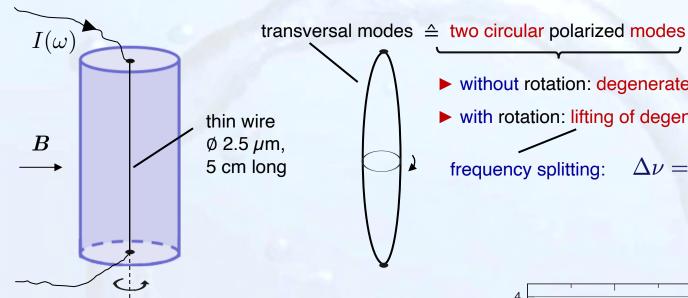
$$\kappa = \frac{h}{m_4} \, n$$





Experimental discovery of quantization of circulation

vibrating wire excited by current pules (Joe Vinen 1961)



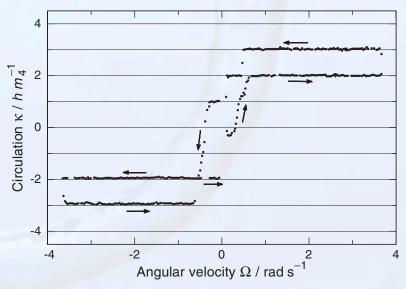
- without rotation: degenerate
- with rotation: lifting of degeneracy by Magnus force

frequency splitting:
$$\Delta \nu = \frac{\varrho_{\rm s}}{2\pi\,\mathcal{M}}\,\kappa$$

effective mass / length (wire + ½ of expelled liquid)

experimental results

- quantization with expected value
- hysteresis effects are observed
- modern measurements up to n = 4

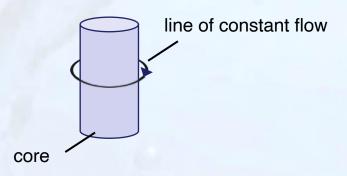






What has this to do with the rotation of bulk helium in a simply connected region?

- vortices may occur with normal fluid core
- resulting in a multiply connected region

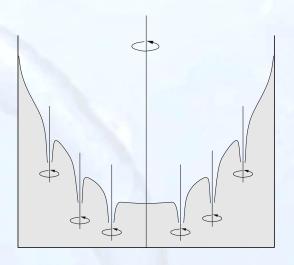


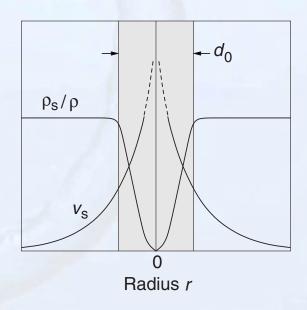


$$v_{\rm s}(r) = \frac{\kappa}{2\pi r} = \frac{1.58 \times 10^{-8}}{r} n \quad \left[\frac{\rm m}{\rm s}\right]$$

 \longrightarrow normal core: $v_{
m s}
ightarrow v_{
m c}$

$$d_0 \approx 2 - 3 \, {
m \AA} \, \, \stackrel{\triangle}{=} \, \, {
m coherence \, length} \, \, {
m healing \, length} \, \,$$









Energy of a vortex

$$E_{
m v}=\int\limits_{a_0}^{b}rac{arrho_{
m s}v_{
m s}^2}{2}\,2\pi r\;{
m d}r$$
 energy / length

 a_0 : radius of vortex core

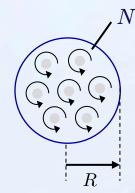
b: radius of vessel or ½ distance to next vortex

$$\kappa = v_{\rm s} \, 2\pi r \quad \longrightarrow \quad v_{\rm s}^2 = \frac{\kappa^2}{4\pi^2 r^2}$$

$$E_{\rm v}=rac{arrho_{\rm s}\kappa^2}{4\pi}\,\ln\left(rac{b}{a_0}
ight)\,\propto\,\kappa^2\,\propto\,n^2$$
 vortex formation with $n=1$ is preferred

Why is not a large vortex forming?

splitting up in many small vortices prohibits large kinetic energy in core of vortex near the axis of rotation (velocity at the edge of vessel is given)



$$N$$
 vortices

$$v_{\rm s}(R) = N \frac{h}{m^4} \frac{1}{2\pi R}$$

 $N \propto R^2$ if evenly distributed



At what velocity vortices are formed?

critical angular velocity

$$\omega_{
m c}=E_{
m v}/L_{
m v}$$

$$L_{
m v}=\int\limits_0^R arrho_{
m s} r\,v_{
m s}\,2\pi r\;{
m d}r=rac{1}{2}arrho_{
m s}\kappa R^2$$
 angular momentum

$$\longrightarrow \omega_{\rm c} = \frac{h}{2\pi m_4 R^2} \, \ln \left(\frac{R}{a_0} \right)$$

$$R = 1 \, \mathrm{cm} \longrightarrow \omega_{\mathrm{c}} \approx 10^{-3} \, \mathrm{s}^{-1}$$

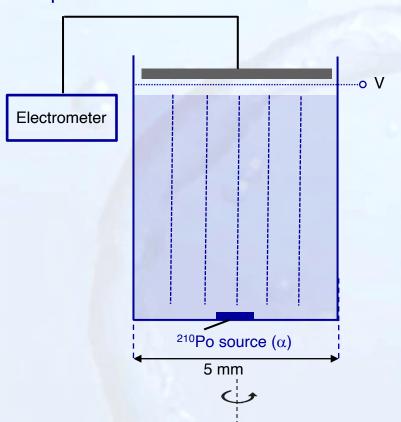
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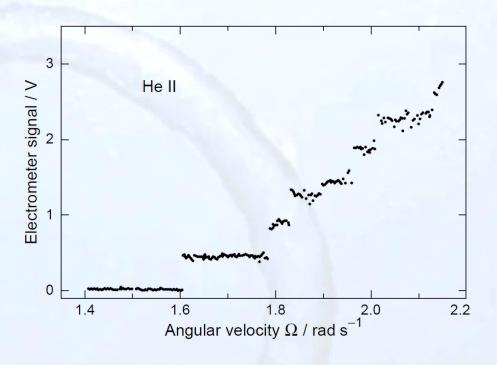
concept of critical velocity will be discussed in section 2.6

- meniscus is rotating vessels
- damping of second sound
- electrometer experiments
- exploding electron bubbles
- decorating with hydrogen ice particles





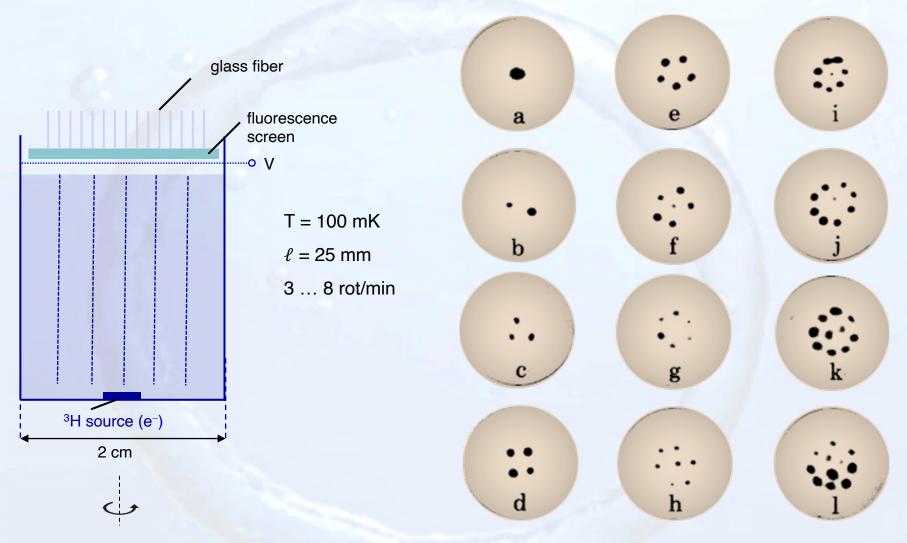




- $ightharpoonup \alpha$ source ightharpoonup helium ionized ightharpoonup electrons form bubbles
- bubbles are captured by vortex lines via Magnus force
- ▶ E field is pulling bubbles alongside of vortex line to surface
- ▶ measurement of charge → is proportional to number vortex lines
- uniform acceleration over 10 h to 10 rot/min



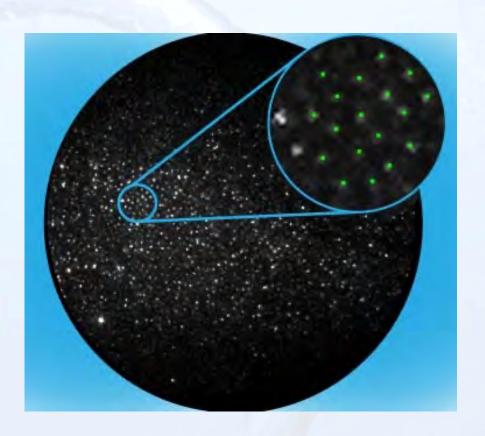








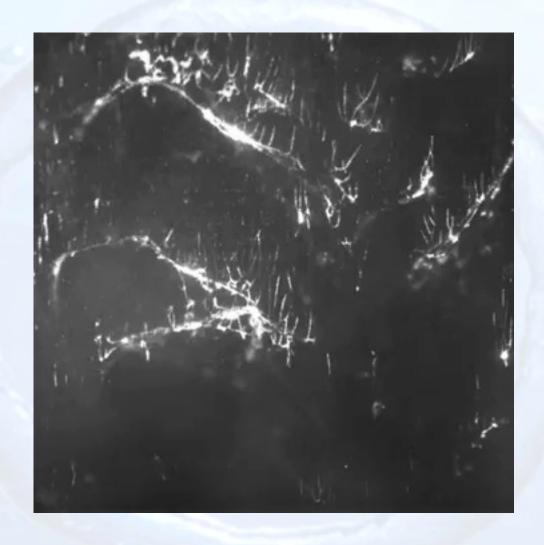




Abrikosov lattice → Type 2 superconductor



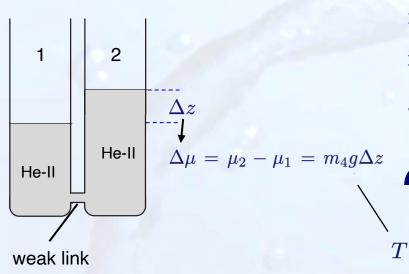








Josephson Effects



healing length $dpprox \xi=1\dots2\, ext{\AA}$ $\xi_4=rac{0.3\, ext{nm}}{(1-T/T_\lambda)^{2/3}}$ diverges for $T o T_\lambda$

Schrödinger Eq.

$$\mathrm{i}\hbar\dot{\Psi}_1 = \mu_1\Psi_1 + \mathcal{K}\Psi_2$$

$$i\hbar\dot{\Psi}_2 = \mu_2\Psi_2 + \mathcal{K}\Psi_1$$

with
$$\varPsi_1=\sqrt{\varrho_{\mathrm{s}}}\mathrm{e}^{\mathrm{i}arphi_1}$$
 and $\varPsi_2=\sqrt{\varrho_{\mathrm{s}}}\mathrm{e}^{\mathrm{i}arphi_2}$



$$\frac{\partial \varrho_{\rm s}}{\partial t} = \frac{2\mathcal{K}}{\hbar} \ \varrho_{\rm s} \sin \left(\varphi_2 - \varphi_1\right)$$

$$\frac{\partial \mathcal{L}^{2}}{\partial t} = \frac{2\mathcal{K}}{\hbar} \, \varrho_{s} \sin(\varphi_{2} - \varphi_{1})$$

$$T = 0 \quad \frac{\partial}{\partial t} (\varphi_{2} - \varphi_{1}) = -\frac{1}{\hbar} (\mu_{2} - \mu_{1}) = -\frac{1}{\hbar} m_{4} g \Delta z$$

 $\Delta \mu = 0$ phase difference constant \longrightarrow Josephson dc effect

 $\Delta \mu \neq 0$ phase difference changes \rightarrow Josephson ac effect

$$T=0$$
 \longrightarrow $\omega_{\mathrm{J}}=rac{\Delta\mu}{\hbar}=rac{m_{4}\Delta p}{\hbar\varrho}$ with $\Delta p=\varrho g\Delta z$

$$T \neq 0 \longrightarrow \omega_{\rm J} = \frac{\Delta \mu}{\hbar} = \frac{m_4 \Delta p}{\hbar \varrho} - m_4 \frac{S \Delta T}{\hbar}$$