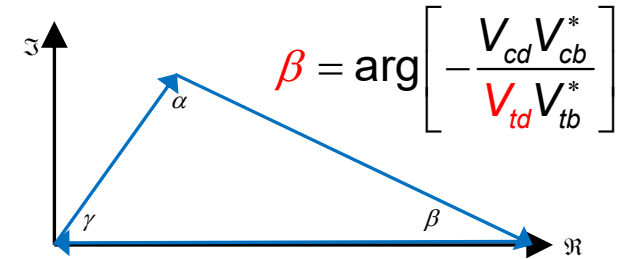
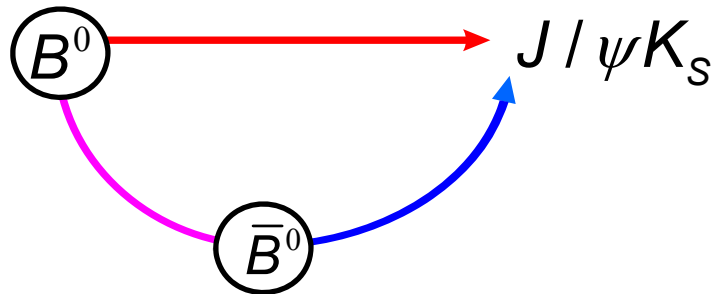


Measurement of the CKM phase β in $B^0 \rightarrow J/\psi K_S$

The decay $B^0 \rightarrow J/\psi K_S$ allows an interference between the decay with and without mixing:



With the master equation one obtains for small $\Delta\Gamma \approx 0$ and negligible (no) direct CPV (i.e. $|\bar{A}_f| = |A_f|$) and no CPV in mixing ($q/p=1$):

$$C_f = 0 \quad S_f = \frac{2\Im\lambda_f}{1+|\lambda_f|^2}$$

$$A_{CP}(t) = \frac{\Gamma(B^0 \rightarrow J/\psi K_S)(t) - \Gamma(\bar{B}^0 \rightarrow J/\psi K_S)(t)}{\Gamma(B^0 \rightarrow J/\psi K_S)(t) + \Gamma(\bar{B}^0 \rightarrow J/\psi K_S)(t)} = -\Im\lambda_f \sin(\Delta m t)$$

For the final-state $f=f_{CP} = B^0 \rightarrow J/\psi K_S$ with CP $|f_{CP}\rangle = \eta_{CP} |f_{CP}\rangle$ and $\eta_{CP} = \pm 1$ one obtains for λ_f

$$\lambda_f \equiv \frac{q}{p} \cdot \frac{\bar{A}_{f_{CP}}}{A_{f_{CP}}} = \frac{q}{p} \cdot \eta_{CP} \frac{\bar{A}_{\bar{f}_{CP}}}{A_{\bar{f}_{CP}}}$$

where we replaced $\bar{A}_{f_{CP}}$ by $\eta_{CP} \bar{A}_{\bar{f}_{CP}}$

$\mathcal{A}_{f_{CP}}$ and $\bar{\mathcal{A}}_{\bar{f}_{CP}}$ are related by **CP-conjugation**, i.e. they differ only in the

Sign of the weak phase: $\bar{\mathcal{A}}_{\bar{f}_{CP}}$ can be derived from $\mathcal{A}_{f_{CP}}$ by inverting the CKM (weak) phase.

For the final-state $J/\psi K_S$ one finds $\eta_{CP} = -1$.

Moreover for this final-state one needs to consider the following amplitudes:

$$\mathcal{A}(B^0 \rightarrow J/\psi K_S) = \mathcal{A}(B^0 \rightarrow J/\psi K^0) \cdot \mathcal{A}(K^0 \leftrightarrow K_S)$$

$$\mathcal{A}(\bar{B}^0 \rightarrow J/\psi K_S) = \mathcal{A}(\bar{B}^0 \rightarrow J/\psi \bar{K}^0) \cdot \mathcal{A}(\bar{K}^0 \leftrightarrow K_S)$$

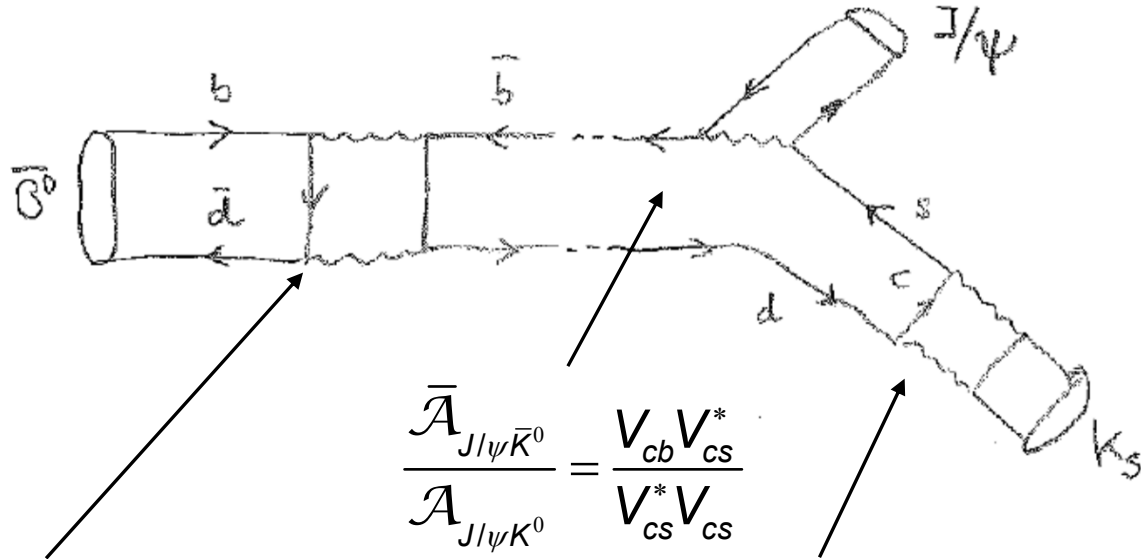
Reminder:

$$|K_S\rangle = p|K^0\rangle + q|\bar{K}^0\rangle$$

Considering the K^0 mixing one obtains:

$$\lambda_{J/\psi K_S} = \left(\frac{q}{p}\right)_{B^0} \cdot \eta_{CP} \cdot \frac{\bar{\mathcal{A}}_{J/\psi \bar{K}^0}}{\mathcal{A}_{J/\psi K^0}} \cdot \left(\frac{q}{p}\right)_{K^0}$$

The respective Feynman-diagram is:



B oscillation: $\Delta m \approx 2 |M_{12}|$

K^0 oscillation:

$$\underbrace{\left(\frac{q}{p}\right)_{B^0}}_{\text{Pure phase}} = \sqrt{\frac{M_{12}^*}{M_{12}}} = \frac{V_{td} V_{tb}^*}{V_{td}^* V_{tb}} = e^{-i2\beta}$$

assuming $|q/p|=1$
(no CPV in mixing)

for B^0 : $\Gamma_{12} \approx 0$

$$\underbrace{\left(\frac{q}{p}\right)_{K^0}}_{\text{real in leading order}} = \sqrt{\frac{M_{12}^*}{M_{12}}}_{K^0} = \frac{V_{cs} V_{cd}^*}{V_{cs}^* V_{cd}}$$

Remark:
 Non trivial phase of λ_f is result of only B^0 mixing. Contribution from V_{td}

$$V_{td} = |V_{td}| e^{-i\beta}$$

Putting everything together one obtains:

$$\lambda_{J/\psi K_S} = \eta_{CP} \left(\frac{V_{td} V_{tb}^*}{V_{td}^* V_{tb}} \right) \left(\frac{V_{cb} V_{cs}^*}{V_{cb}^* V_{cs}} \right) \left(\frac{V_{cs} V_{cd}^*}{V_{cs}^* V_{cd}} \right) = - \left(\frac{V_{td} V_{tb}^*}{V_{td}^* V_{tb}} \frac{V_{cb} V_{cd}^*}{V_{cb}^* V_{cd}} \right)$$

$$\begin{aligned} \Im(\lambda_{J/\psi K_S}) &= -\sin\left(\arg\left(\frac{V_{td} V_{tb}^*}{V_{td}^* V_{tb}} \frac{V_{cb} V_{cd}^*}{V_{cb}^* V_{cd}}\right)\right) = -\sin\left(2 \arg\left(\frac{V_{cb} V_{cd}^*}{V_{td}^* V_{tb}}\right)\right) \\ &= +\sin(2\beta) \end{aligned}$$

$$A_{CP}(t) = -\Im(\lambda_{J/\psi K_S}) = -\sin(2\beta) \sin(\Delta mt)$$

i.e. the observation of time-dependent $A_{CP} \rightarrow$ amplitude is a measure of CKM phase β (mixing phase)

The first measurements of the CKM angle β was performed by the BABAR and Belle collaborations in 2001 at the high-luminosity e^+e^- B-factories:

e^+e^- B-factories were operated at the center-of-mass energy of the $\Upsilon(4S)$ ($b\bar{b}$) resonance ($\sqrt{s}=10.56$ GeV) which decay nearly entirely to B^+B^- and B^0B^0 :

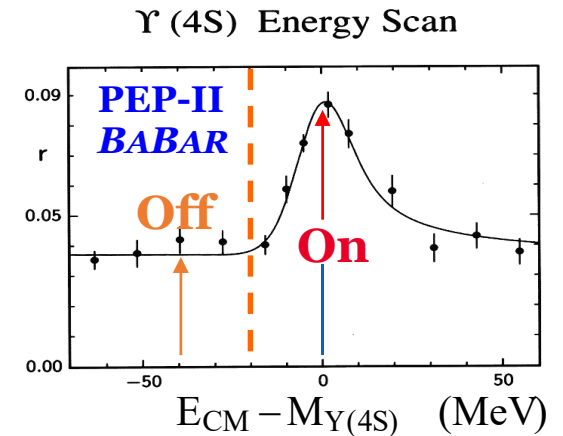
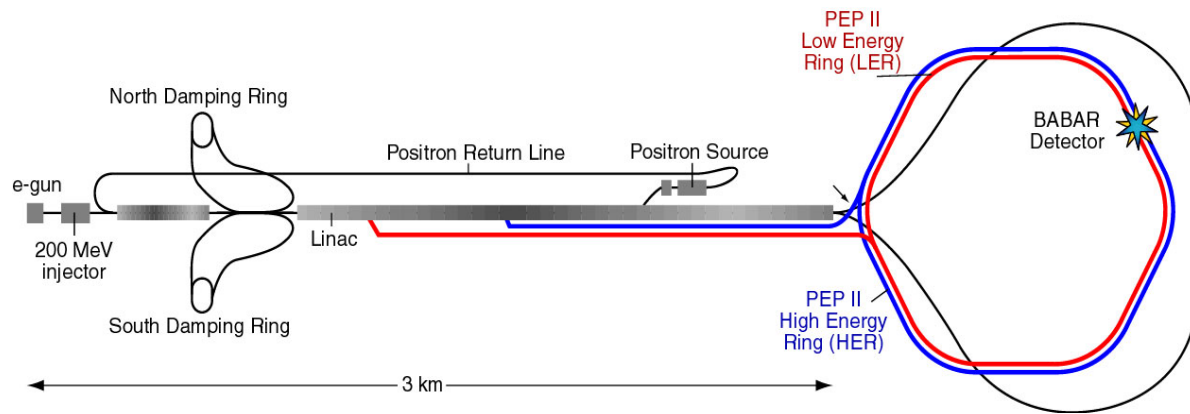
$$e^+e^- \rightarrow \Upsilon(4S) \rightarrow B\bar{B} \quad (\sqrt{s} = 10.58 \text{ GeV} \approx 2m_B)$$

The measurement of $\sin(2\beta)$ was the first observation of CP violation outside the Kaon system \rightarrow predicted by the Standard Model

Experimental problems:

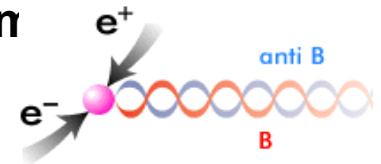
- In the $\Upsilon(4S)$ the BB -system has very small kinetic energy $\rightarrow \gamma$ -factor is very small $\rightarrow \gamma c\tau$ is very small: B's don't fly but decay at rest \rightarrow can't perform t -dependent measurement (flight-length is usually used as measure of the decay-time).
 \rightarrow Idea: produce the $\Upsilon(4S)$ with a small boost by building an asymmetric storage ring
- Flavor-Tagging: Determine the flavor of signal-B at time of production.
Idea: use the second B!

PEP-II B-factory and BABAR Experiment ($\Upsilon(4S) \rightarrow B\bar{B}$)



	Design	Reached
E [GeV] e^- / e^+	9.0 / 3.1	✓
L [$\text{cm}^{-2} \text{s}^{-1}$]	3×10^{33}	6.58×10^{33}
L_{int} [$\text{pb}^{-1}/\text{day}$]	135	391

Coherent $B^0 \bar{B}^0$ system



- 2 bosons \Rightarrow symmetric wavefunction

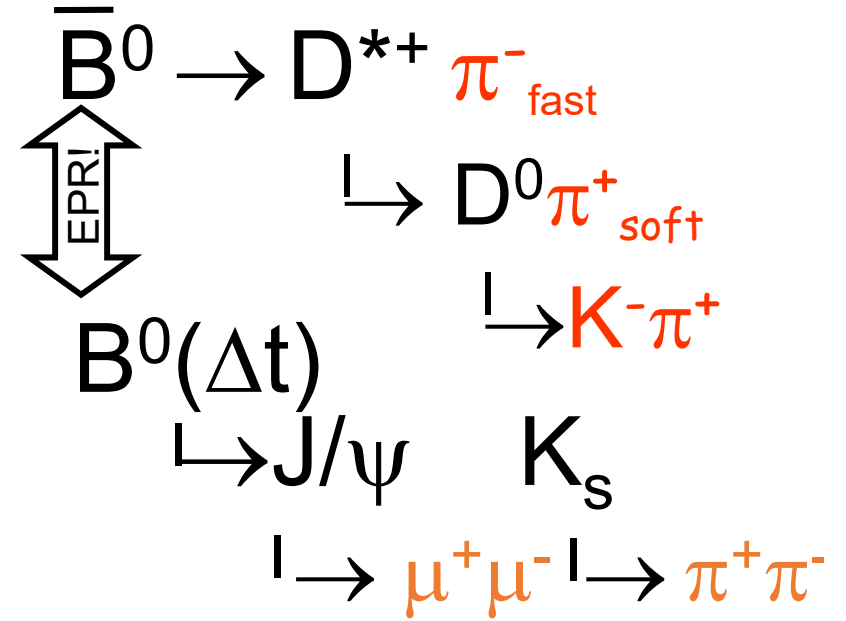
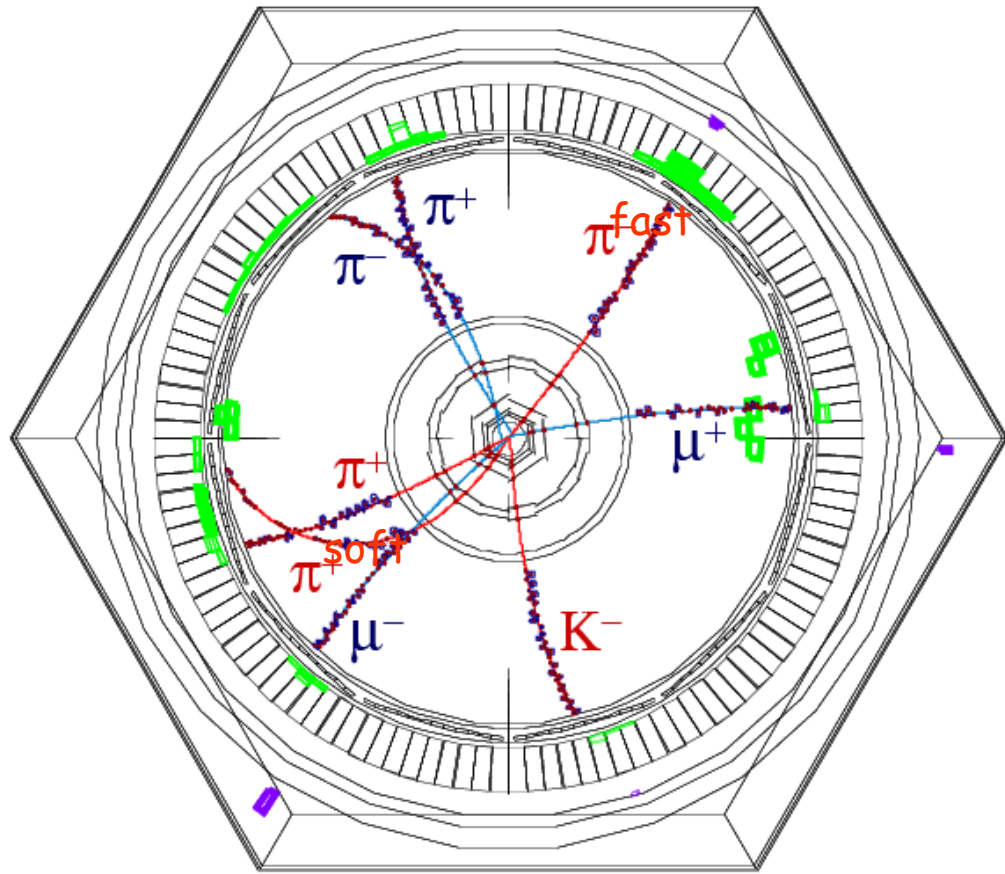
$$\Psi = \Psi_{\text{Flavour}} \cdot \Psi_{\text{space}}$$

- L=1: $\Psi_{\text{space}} \propto (-1)^L \rightarrow$ antisymmetric

- $\rightarrow \Psi_{\text{Flavour}}$ antisymmetric

$$\Psi_{\text{Flavour}} = \frac{1}{\sqrt{2}} (B^0 \bar{B}^0 - \bar{B}^0 B^0)$$

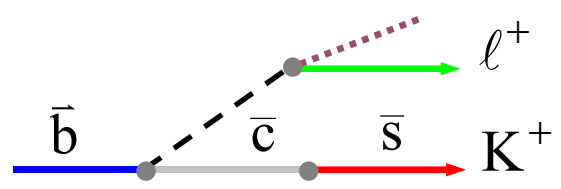
Example of a fully reconstructed $B^0 \rightarrow J/\psi K_s$ event



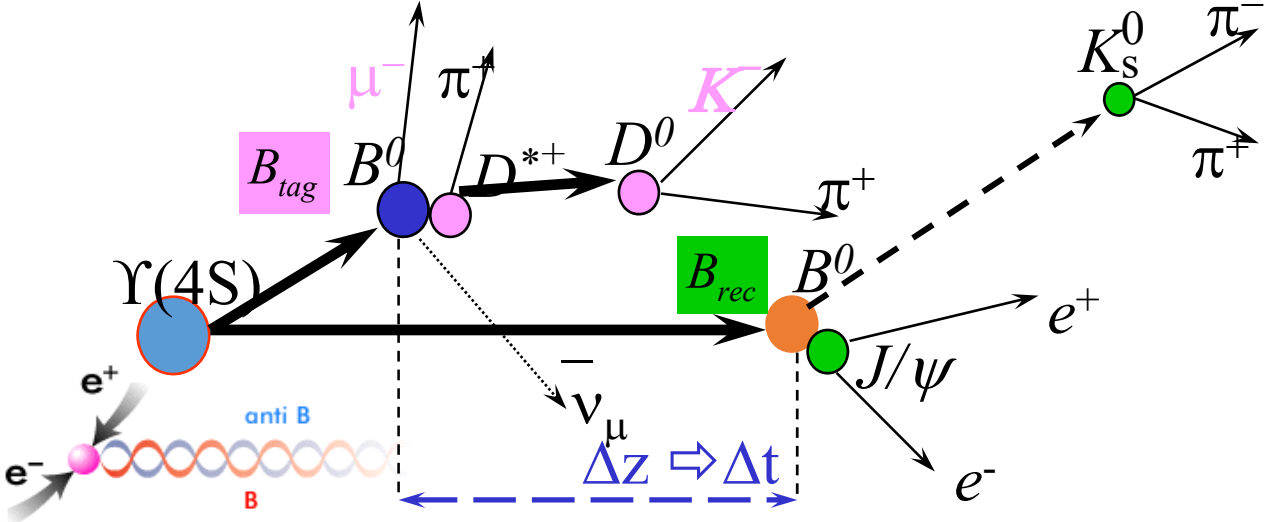
In general, use charges of identified

- leptons,
- kaons,
- soft pions

from the “the rest of the event” to tag B flavour

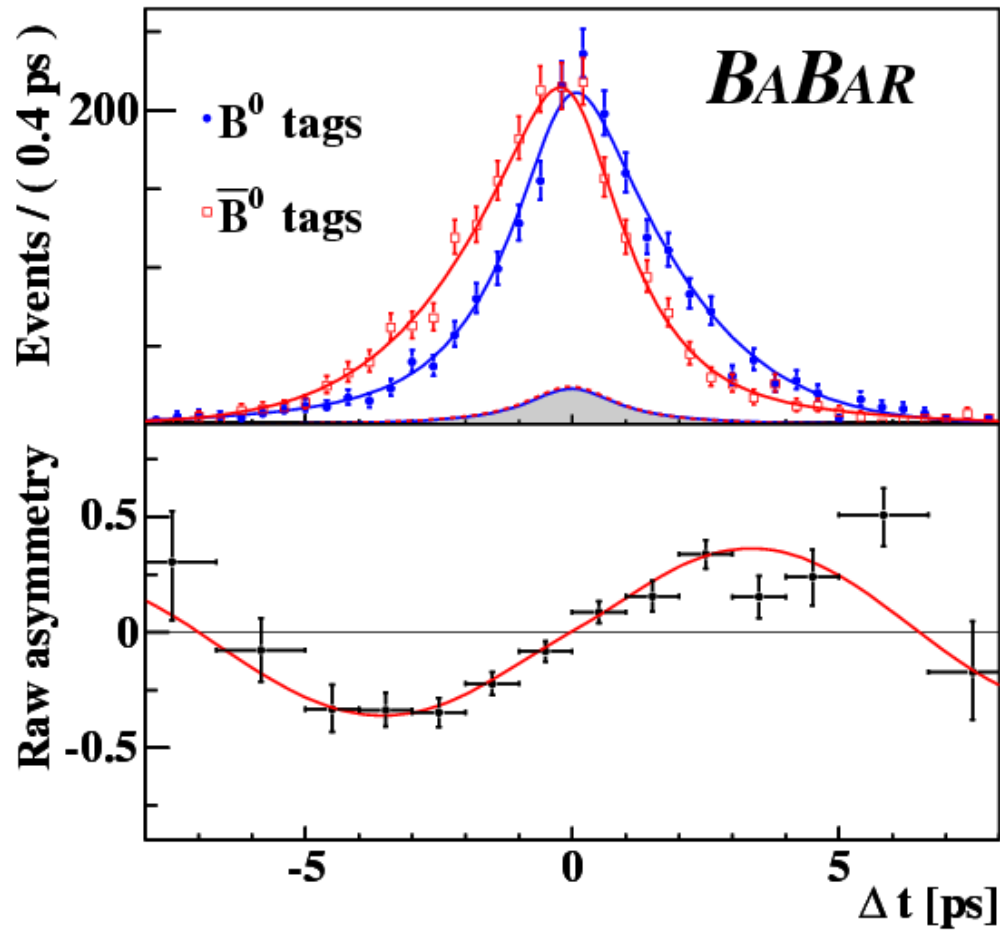


Event configuration:



The variable Δt for the coherent state allows for negative values (tagging B decays after the signal B)

$(c\bar{c}) K_S$ (CP odd) modes



BABAR:

$$\sin(2\beta) = 0.687 \pm 0.028 \pm 0.012$$

Measurement at LHCb

Flavor-tagging w/ 2nd b-hadron in the event:

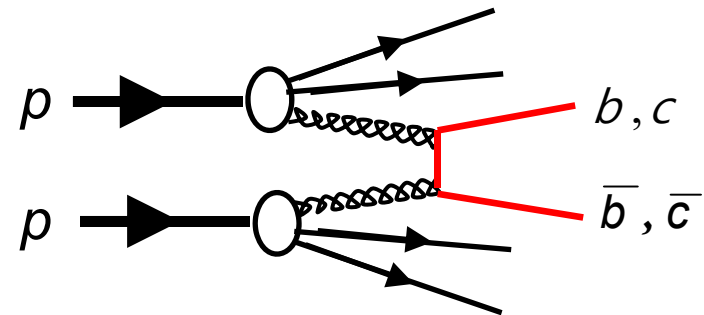
- b and anti-b-quarks are always produced in pairs. Use flavor specific decays

$$b \rightarrow c + \ell^-$$

$$\bar{b} \rightarrow \bar{c} + \ell^+$$

$$B_{tag} \rightarrow D + \ell \nu$$

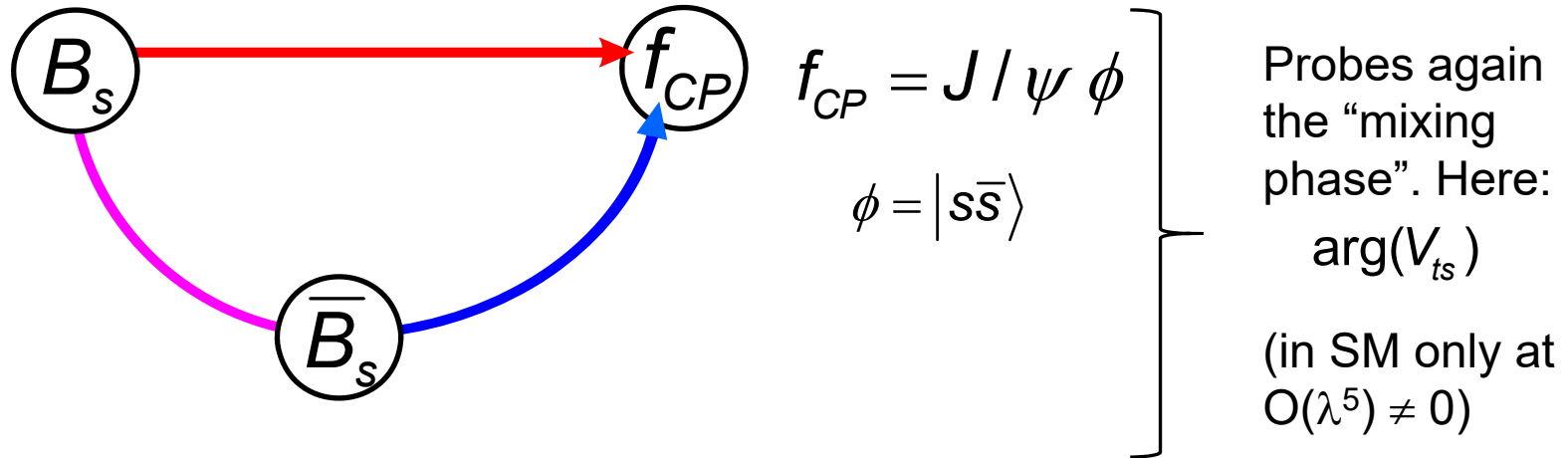
- Hadronization allows the production of two different b-hadrons. No coherence! (large mistag fraction)



$$\sin(2\beta) = 0.76 \pm 0.034$$

Measurement of the CKM phase β_s

Reminder: β_s is the phase of V_{ts} (very small)



Differences w/r to $B^0 \rightarrow J/\psi K_S$

$$1.) \quad \left(\frac{q}{p} \right)_{B_s^0} = \frac{V_{ts} V_{tb}^*}{V_{ts}^* V_{tb}} \approx e^{-i2\beta_s} \quad \Longrightarrow \quad \Im(\lambda_{J/\psi\phi}) = +\sin(-2\beta_s)$$

($2\beta_s$ is very small (0.04) \rightarrow very small CPV)

2.) ϕ is not oscillating like the K^0

3.) $\Delta\Gamma \neq 0$: $A_{CP}(t)$ is described by the more complicated Master formula

$$A_{CP}(t) = \frac{-\Im(\lambda_{J/\psi\phi}) \sin(\Delta m_s t)}{\cosh(\frac{1}{2} \Delta\Gamma_s t) + \Re(\lambda_{J/\psi\phi}) \sinh(\frac{1}{2} \Delta\Gamma_s t)}$$

4.) Additional problem: $J/\psi\phi$ is not a pure CP eigenstate but a mixture of $\eta_{CP} = \pm 1$ states.

Reason: J/ψ and ϕ are both vector mesons (Spin = 1)

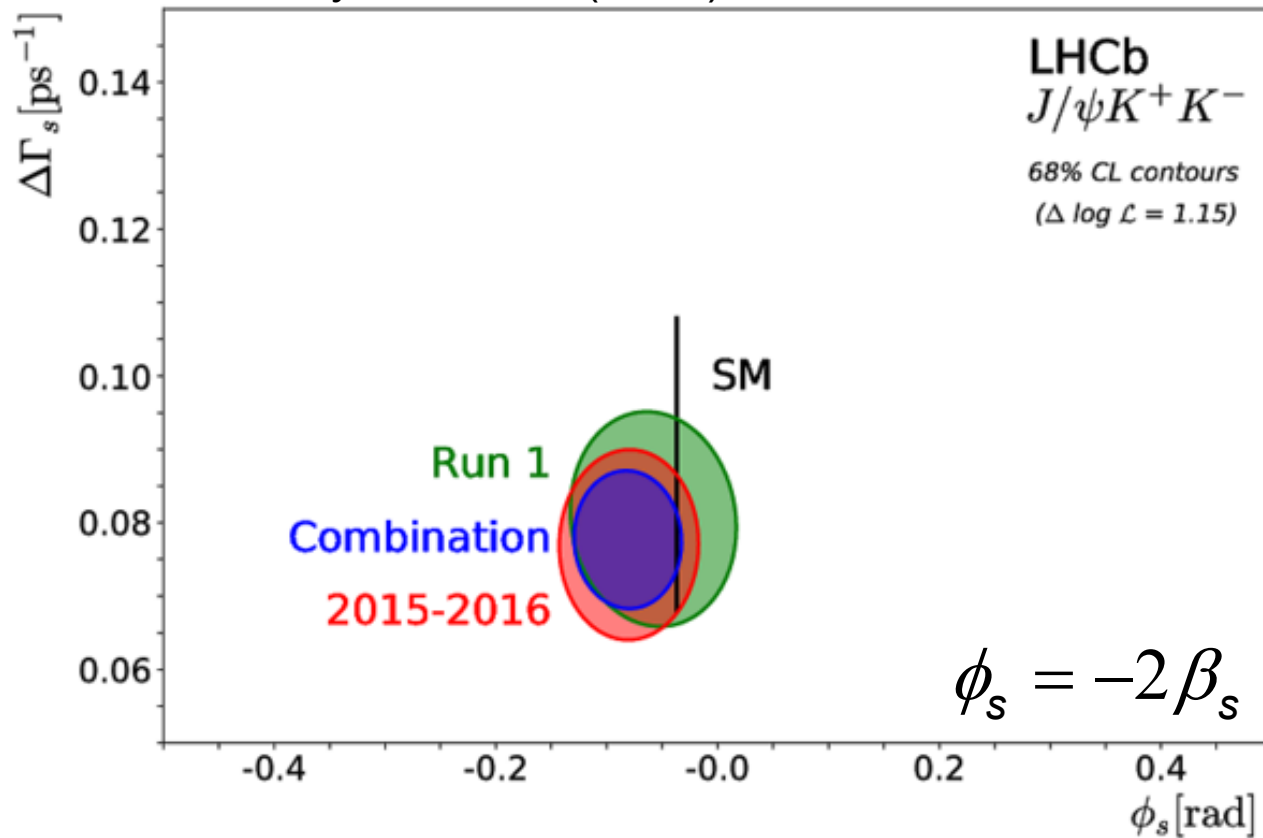
$$B_s \rightarrow J/\psi + \phi \quad \left. \begin{array}{l} J^{PC} = 0^{--} \quad J^{PC} = 1^{--} \quad J^{PC} = 1^{--} \end{array} \right\} \text{Relative angular momentum } L=0,1,2,$$

$$\eta_{CP} = \underset{+1}{\eta_{CP}(J/\psi)} \cdot \underset{+1}{\eta_{CP}(\phi)} \cdot (-1)^L \quad \begin{array}{l} L=1: \text{ CP-odd} \\ L=0,2: \text{ CP-even} \end{array}$$

Requires an angular analysis to separate the CP even from the CP odd states.

Simultaneous measurement of $\Delta\Gamma$ and $\phi_s = -2\beta_s$

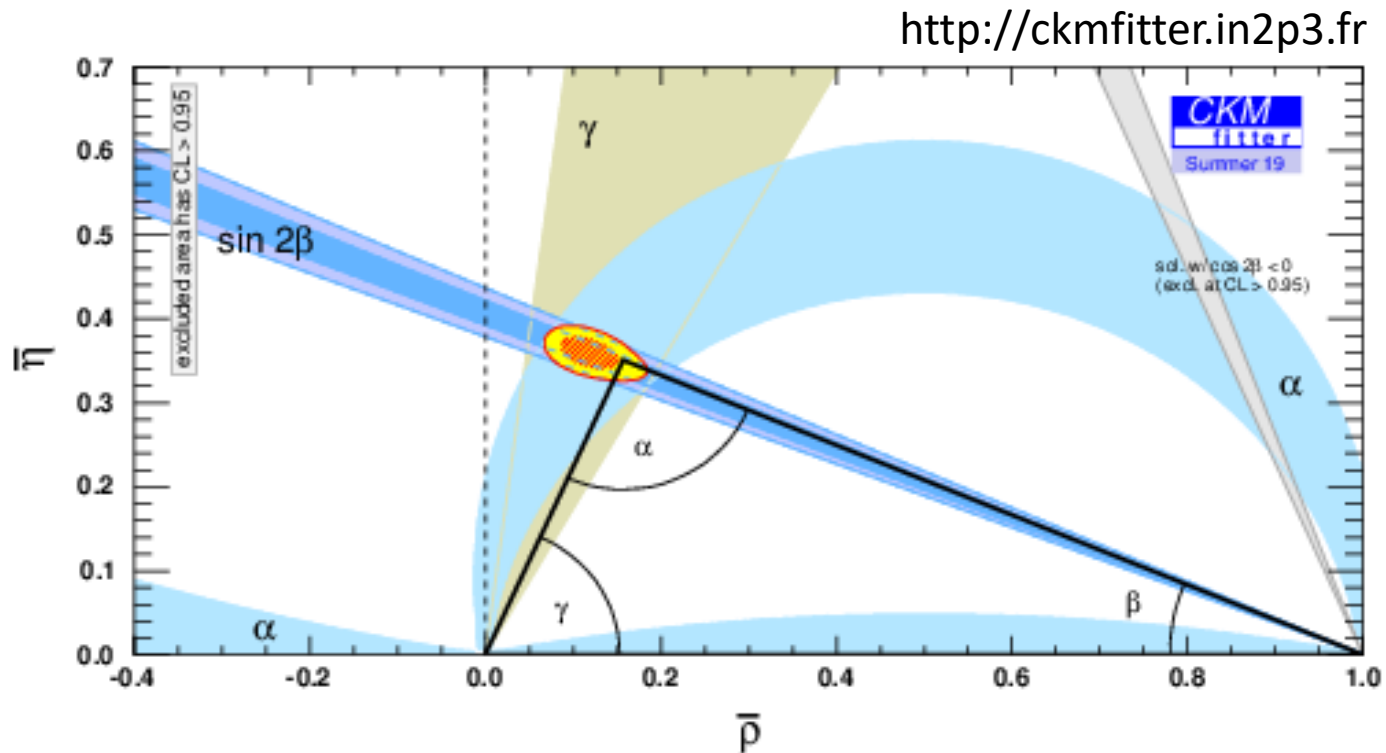
Eur. Phys. J. C. 79 (2019)



Similarly to the β determination in $B^0 \rightarrow J/\psi K_s$ the determination of the weak phase has very little uncertainty: Using the β -measurement \rightarrow predict $\phi_s = -2\beta_s$ with an excellent precision.

Status of CKM angles

Unitarity triangle can be checked with only measurements of angles (CPV):



and compared to the determination of unitarity triangle using all information: