

Physics at the LHC

Hot QCD matter produced in ultra-relativistic heavy-ion collisions

Lecture 3 January 15, 2020



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Outline





Heavy-ion collisions: Little Bangs



Experimental program at the LHC, ALICE



Global characteristics





Quark-gluon plasma tomography with hard probes



Research plans for near and further future

Outline



- Historical view: experimental opportunities, first theoretical ideas
- Thermodynamics of strongly interacting matter and phase diagram
- QCD matter under extreme conditions in nature and in the lab
- Global characteristics: centrality, energy density, multiplicities
- Bulk (soft) particle production
 - Thermal model, particle yields and chemical freeze-out
 - (hyper-)(anti-)nuclei
 - Hydrodynamics, flow and correlations
 - Particle spectra, Blast-Wave fits and kinetic freeze-out
 - Small systems
- Hard probes
 - Jets, heavy quarks, heavy quarkonia
- Future



- 1. Initial conditions, pre-equilibrium, hard scattering processes
- 2. Thermalization: equilibrium is established ($t_{eq} \le 1$ fm/c)
- **3**. Expansion and cooling ($t_{QGP} < 10 15$ fm/c)
- 4. Hadronization (quarks and gluons form hadrons)
- 5. Chemical freeze-out: inelastic collisions cease, yields are defined
- Kinetic freeze-out: elastic collisions cease, spectra are frozen (t_{had} ~ 3-5 fm/c)

Measurements can only be performed at stages 5 and 6 From those, we want to deduce information on phases 2, 3, 4

≈ 10⁻²² s



- Initial conditions, pre-equilibrium, hard scattering processes 1.
- 2. Thermalization: equilibrium is established ($t_{eq} \le 1$ fm/c)
- 3. Expansion and cooling ($t_{OGP} < 10 15$ fm/c)
- Hadronization (guarks and gluons form hadrons) 4.
- Chemical freeze-out: inelastic collisions cease, yields are defined 5.
- $(t_{had} \sim 3-5 \text{ fm/c})$ 6. Kinetic freeze-out: elastic collisions cease, spectra are frozen

Evolution of system from equilibration to freeze-out can be described as motion and interaction of fluid cells, as collective system \rightarrow use of fluid dynamics



Good applicability to large systems, as in heavy-ion collisions:

- Long distances, long times, strong fields
- System described by macroscopic medium properties:
 - Equation of state P(T)
 - Shear viscosity η(T)
 - Bulk viscosity ζ(T)
 - Heat conductivity κ(T)
 - Relaxation times $\tau_{shear}(T)$, $\tau_{bulk}(T)$, etc



S. Floerchinger

Assumption: local equilibrium condition

Ingredients:

- Streaming fluid described by velocity field
- ε, p: energy and pressure in co-moving frame of fluid cell

 T^{μ}

- Energy-momentum tensor:
- Charge current: $j^{\mu} = n u^{\mu}$ e.g. baryon number $j^{\mu_{B}} = n_{B}(x)u^{\mu}(x)$

• Equations of relativistic hydrodynamics:

Equation of motion from energy-momentum conservation:

Charge conservation, e.g. baryon number:



$$u^{\mu} = \frac{dx^{\mu}}{dt}$$

$$v = (\epsilon + p)u^{\mu}u^{\nu} - pg^{\mu\nu}$$

$$\partial_{\mu} T^{\mu\nu} = 0$$

 $\partial_{\mu} j^{\mu} = 0$

Ingredients

- Equations of relativistic hydrodynamics need proper formalism (e.g. Isreal Stewart theory) to have stable and causal equations
- Equation of State (EOS) relates thermodynamic variables ε(p,...) generally taken from lattice QCD computations
- Transport coefficients: bulk and shear viscosity, their relaxation times, heat conductivity, ...
- Initial conditions: energy density distribution MC Glauber, IP Glasma, MC-KLN, TRENTo ...
- Final hadron phase (interactions)
 UrQMD, SMASH

$$\partial_{\mu} \mathsf{T}^{\mu\nu} = \mathbf{0} \qquad \partial_{\mu} \mathsf{j}^{\mu} = \mathbf{0}$$



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Hydrodynamic frameworks





Different realizations:

- Simulation codes (very CPU intensive!):
 VISHNU, MUSIC, ECHO-QGP, etc
- Analytic, mode-by-mode: elegant and fast!
 FLUIDUM (+FastReso for final phase)



Radial expansion

Central collisions, transverse plane



Depends on bulk viscosity $\zeta(T)$

Demonstrated by particle spectra



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Thermalized source + radial flow boost



Transverse momentum spectra described as sum of two components:

• Thermal distribution:

$$\frac{1}{m_{T}} \frac{dN}{dm_{T}} \propto m_{T} K_{1} \left(\frac{m_{T}}{T} \right) \stackrel{m_{T} \gg T}{\rightarrow} \sqrt{m_{T}} e^{-m_{T}/T}$$

- Additional component due to common transverse velocity from radial expansion Boltzmann-Gibbs Blast-Wave: simplified hydro model, with parameters:
 - T_{kin}: kinetic freeze-out temperature
 - $<\beta_{T}>$: transverse radial flow velocity





E. Schnedermann et al., Phys. Rev. C48 (1993) 2462

Identified particle spectra at the LHC



Perform a simultaneous fit of pion, kaon, proton spectra in given centrality range, in the low (hydrodynamic-dominated) region







- < β_T >: transverse radial flow velocity increasing with centrality Order of 0.6 – 0.66 at the LHC
- Pb-Pb and Xe-Xe consistent

Blast-wave fits: all collision systems





A. Andronic, arXiv:1407.5003

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Hydro-models

iEbyE + VISHNU + Trento/AMPT: viscous hydro with different initial conditions

[arXiv:1703.10792v1, PRC 92, (2015) 014903, PRC 92 (2015) 011901(R)]

McGill: MUSIC viscous hydro with IP-Glasma initial conditions [PRC 95 (2017) 064913]

EPOS-LHC: core (hydro) + corona [*PRC 92 (2015) 034906*]



v_n flow coefficients

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Non-central collisions





Initial spatial asymmetry b Pressure gradients larger in one direction \rightarrow larger fluid velocity in that direction

 \rightarrow more particles fly in that direction



Particle distribution:

$$\frac{\mathrm{dN}}{\mathrm{d}\phi} = \frac{\mathrm{N}}{2\pi} [1 + 2\sum_{\mathrm{n}} \mathrm{v}_{\mathrm{n}} \cos(\mathrm{n}(\phi - \psi_{\mathrm{R}}^{\mathrm{n}}))]$$

In case of homogeneous initial energy density, for symmetry reasons $(\Phi \rightarrow \Phi + \pi)$: $v_1 = v_3 = v_5 = ... = 0$



$$\frac{\mathrm{dN}}{\mathrm{d}\phi} = \frac{\mathrm{N}}{2\pi} [1 + 2\sum_{\mathrm{n}} \mathrm{v}_{\mathrm{n}} \cos(\mathrm{n}(\phi - \psi_{\mathrm{R}}^{\mathrm{n}}))]$$

In non-central collisions, the second component (coefficient v_2 , elliptic flow) is the dominant one



Depending on the collision energy, particles prefer to be "in plane" or "out of plane"





A. Andronic, arXiv:1407.5003

Very low energies: in-plane, rotation like emission. Low energy density and long reaction times. SIS region: competition between increasing speed of expansion and decreasing passage time of the spectators (hindering expansion in reaction plane) → squeeze-out High energies: unhindered collective expansion of the initially-anisotropic fireball



Comparison between RHIC (200 GeV) and the LHC (2.76 TeV)

- p_{T} -integrated v_{2} at the LHC ~30% higher than at RHIC
- as a function of the charged-particle p_{τ} : practically identical



PRL 105 (2010) 252302





JHEP 06 (2015) 190

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Identified particles: mass dependence and baryon/meson grouping

Mass ordering of v_2 (combined effect with radial flow)



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Example from 2013





Figure from A. Andronic, arXiv:1407.5003 Hydro: VISHNU, from arXiv:1311.0157 MC-KLN, VISHNU with constant η/s, UrQMD

2-particle correlations



Most central events (0-1%): 2-particle azimuthal correlation

sum of $v_2^{}$, $v_3^{}$, $v_4^{}$, $v_5^{}$ harmonics

 v_3 , v_5 non-zero!



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Lumpy energy density profile of initial state



Fluctuations of the complex structure of the energy density profile of the initial state produce non-zero contribution to all harmonics









The energy dependence of anisotropic flow provides additional discriminating power over initial-state models and temperature dependence of transport coefficients.

JHEP 07 (2018) 103





 $v_{2}^{2}, v_{2}^{4}, v_{2}^{6}, v_{2}^{8}$

The energy dependence of anisotropic flow provides additional discriminating power over initial-state models and temperature dependence of transport coefficients.

JHEP 07 (2018) 103

Comparison to hydro models

Results: 2018, ALICE and CMS

Charged particles

 $v_{2}^{2}, ..., v_{6}^{2}$ with 2 particle cumulants

Hydro models:

- MUSIC with IP-Glasma initial conditions, UrQMD
- iEbE-VISHNU with initial conditions from either AMPT or Trento, UrQMD
- η/s constant or parametric function



JHEP 07 (2018) 103



Hot and Dense QCD White Paper (2012)

- Tremendous evolution over the last years
- Hydro implementation more and more refined ...
- ... however still depending on much modeling!!
- High-precision data challenge the hydro computations







Multi-parameter fit with Bayesian approach

 $20.2^{+2.2}_{-2.2}$

28.0

18.5 N



0-5% Pb-Pb, $\sqrt{s_{_{\rm NN}}}$ = 2.76 TeV

Local contribution within ISOQUANT (SFB1225)

Transport coefficients: compute their temperature and density dependence with first principle QCD

- Functional renormalization group approach
- Lattice QCD

IP-Glasma IC, MUSIC, UrQMD

Dubla, S.M., Pawlowski, Schenke, Shen, Stachel, Nucl. Phys. A979 (2018) 251

30-40% Pb-Pb, $\sqrt{s_{NN}} = 2.76 \text{ TeV}$

- Fluid dynamics of heavy ion collisions with Mode expansion (FluiduM) Analytic solution of hydrodynamic equations
- Resonance decays (FastReso)

Speed and flexibility: opens new doors and opportunities!

Floerchinger, Grossi, Lion arXiv: 1811.01870

Devetak, Dubla, Floerchinger, Grossi, Masciocchi, Mazeliauskas, Selyuzhenkov, arXiv:1909.10485

Hadrochemistry vs system size with a "parenthesis" on "small systems"

proton - proton -----

Reference system, normalization e.g. denominator in the nuclear modification factor R_{AA}

Control experiment

mostly to address cold nuclear matter, initial-state effects (shadowing, gluon saturation, nPDF)

New at the LHC: p-Pb and pp collision data (at high multiplicities) show features which are suggestive of the collective behavior known from Pb-Pb

→ examine evolution vs system size ↔ multiplicity of produced particles Now: look at hadron yields vs charged-particle multiplicity

Smooth evolution of particle production yields from small to large systems, as a function of charged particle multiplicity at mid-rapidity

pp \sqrt{s} = 7 TeV p-Pb $\sqrt{s_{NN}}$ = 5.02 TeV Pb-Pb $\sqrt{s_{NN}}$ = 5.02 TeV

Common hadron production mechanism in all systems, regardless of type and \sqrt{s} ?

Smooth evolution of particle production yields from small to large systems, as a function of charged particle multiplicity at mid-rapidity

pp $\sqrt{s} = 7$ TeV **pp** $\sqrt{s} = 13$ TeV p-Pb $\sqrt{s_{NN}} = 5.02$ TeV Pb-Pb $\sqrt{s_{NN}} = 5.02$ TeV Xe-Xe $\sqrt{s_{NN}} = 5.44$ TeV

Common hadron production mechanism in all systems, regardless of type and \sqrt{s} ?

Focus on strageness vs system size

Strong increase of strangeness production with growing charged particle multiplicity at mid-rapidity, until it reaches a saturation

→ gran-canonical plateau

Increasing slope from s (K⁰_s, Λ^0) to ss (Ξ) to sss (Ω)

Clear difficulty of models to describe the data, increasing with S

Same mass ordering in the elliptic flow measured:

Phys. Lett. B 726 (2013) 164–177

The "baryon anomaly" in Pb-Pb: p_{τ} -dependent distribution of baryons and mesons abundances in different centrality classes:

 Λ/K_{s} ratio

Effect of radial flow, which makes the spectra harder, with a strength proportional to the mass of the particle

arXiv:1307.5530

The "baryon anomaly" in Pb-Pb: p_{τ} -dependent distribution of baryons and mesons abundances in different centrality classes:

Λ/K_{s} ratio

Effect of radial flow, which makes the spectra harder, with a strength proportional to the mass of the particle

Also in p-Pb?

Summary

Anisotropic flow is determined by the response of the system to its initial spatial anisotropies. Initial state spatial anisotropies come in turn from both the geometry of the collision and fluctuations in the wave function of the incident nuclei. The significant magnitude of anisotropic flow is interpreted as evidence of the formation of a strongly-coupled system, which can be described as s fluid with very low shear viscosity to entropy density ratio.

- Standard model of heavy-ion collisions
- Provides an overall good description of many features of experimental data
- Relatively many free parameters: important to follow a much more rigorous approach to fix some (e.g. with principle QCD input) and learn about/constrain the others
- Pursue the investigation of small systems to ultimately determine the true driving physics behind the common phenomena

Spares

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