

## Normalfluid component:

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$$\rho_{\rm n,r} = \frac{2 p_0^4}{3 \hbar^3} \sqrt{\frac{m^*}{(2\pi)^3 k_{\rm B} T}} e^{-\Delta_{\rm r}/k_{\rm B} T}$$
Rotons

 $\varrho_{\rm n} = \varrho_{\rm n,ph} + \varrho_{\rm n,r}$ 

$$\sum \varrho_{\rm n,ph} = \frac{2\pi^2 k_{\rm B}^4}{45\,\hbar^3\,v_1^5}\,T^4$$





- ▶ at low temperatures  $\rho_{\rm n} \propto T^4$  due to phonons
- rotons dominate between 0.5 K and 1.2 K
- above 1.2 K nature of excitations more complex

# 2.6 Excitation Spectrum of He-II: Landau Model



## Specific heat:

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a) low temperatures T < 0.6 K

only long wavelength phonons contribute

→ Debye model

$$C_{\rm ph} = \frac{2\pi^2 k_{\rm B}^4}{15 \varrho \hbar^3 v_1^3} \ T^3$$

measurement of thermal conductivity

Casimir regime  $\ell = d$  capillary cross section

$$\longrightarrow \Lambda = \frac{1}{3} C_{\rm ph} v \, d \propto T^3$$



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b) intermediate temperatures  $0.6 < T < 1.2 \,\mathrm{K}$ 

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free energy 
$$F_{\rm r} = -k_{\rm B}Tn_{\rm r}$$
  
 $S_{\rm r} = -\partial F_{\rm r}/\partial T$   
 $\downarrow$   
 $C_{\rm r} = T\partial S_{\rm r}/\partial T$   
 $n_{\rm r} = \frac{2p_0^2}{3\varrho\hbar^3}\sqrt{\frac{m^*k_{\rm B}T}{(2\pi)^3}} e^{-\Delta_{\rm r}/k_{\rm B}T}$   
number density of rotons

$$C_{\rm r} = \frac{2k_{\rm B}p_0^2}{3\rho\hbar^3} \sqrt{\frac{m^*k_{\rm B}T}{(2\pi)^3}} \left\{ \frac{3}{4} + \frac{\Delta_{\rm r}}{k_{\rm B}T} + \left(\frac{\Delta_{\rm r}}{k_{\rm B}T}\right)^2 \right\} \, \mathrm{e}^{-\Delta_{\rm r}/k_{\rm B}T}$$

c) high temperatures  $1.2 \, {
m K} < T < T_{\lambda}$ 

additional excitations contribute: maxons lifetime of rotons becomes very short

excitations are not well-defined





## Landau's concept of critical velocity

superconductors ---- energy gap

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superfluid He-II ----- no energy gap, but velocity gap!

Landau's Gedankenexperiment: dropping a massive sphere in He-II at T = 0

let's assume that sphere generates one excitation with energy  $\mathcal{E}$  and momentum p

> How fast must this sphere fall in He-II to generate dissipation ?

> > (2)



energy conservation

$$\frac{1}{2}\mathcal{M}v^2 = \frac{1}{2}\mathcal{M}v'^2 + \mathcal{E}$$
 (1)

momentum conservation  $\mathcal{M} \boldsymbol{v} - \boldsymbol{p} = \mathcal{M} \boldsymbol{v}'$ 

not all combinations of  $\mathcal{E}$  and p fulfill both conservation law's at the same time, even if the direction of the excitation is not fixed

phonons can be excited at arbitrary small energies

# 2.6 Excitation Spectrum of He-II: Landau Model





▶ for  $v \ge v_c$  sudden onset of dissipation, laminar  $\longrightarrow$  turbulent flow

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Testbed for the generation of excitations and the critical velocity

type of ions:

- ▶ electrons (–) : zero-point motion  $\rightarrow$  bubbles r = 19 Å
- ▶ <sup>4</sup>He<sup>+</sup>, H<sub>2</sub><sup>+</sup> (+) : attract He atoms  $\rightarrow$  snowballs  $r \approx 7$  Å
- other ions (-, +) : properties depend on wave function

#### Electrons in liquid He

electrons need energy to be emerged in helium  $\sim 1 \text{ eV}$ , which means they need more that 1 eV of kinetic energy to enter liquid He.

#### comment:

similar to work function of electrons in metals





## 2.7 Motion of lons in He-II



#### Energy of bubble

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bubble size:

 $\frac{\partial E}{\partial r} = 0 \longrightarrow r_{\min}(p=0) = 19 \text{ Å}$ 

#### size depends on pressure







exploding bubbles at negative pressure





## Acceleration of ions in constant field



0.7 K < T < 1.8 K: rotons should dominate however, difficult to observe because of other excitations / impurities







in ultra-pure He-II under pressure ions can be accelerated up to Landau velocity

- negative ions accelerated in electric field under high pressure
- drag is measured by time-of-flight method
- in He-I: drag proportional to velocity
- ▶ in He-II: drag is not observable until critical velocity is reached

pressure dependence of  $v_{
m c}$ 









## *T* < 0.3 K

